

Cite this: DOI: 10.1039/xxxxxxxxxx

Real-space characterization of hydroxyphenyl porphyrin derivatives designed for single-molecule devices

Akitoshi Shiotari,^{a‡} Yusuke Ozaki,^a Shoichi Naruse,^a Hiroshi Okuyama,^{*a} Shinichiro Hatta,^a Tetsuya Aruga,^a Takashi Tamaki,^b and Takuji Ogawa^bReceived Date
Accepted Date

DOI: 10.1039/xxxxxxxxxx

www.rsc.org/journalname

The porphyrin derivatives are potential candidates as constituents of functional molecular devices because their electronic levels can rationally be manipulated by chemical modification. In this work, we deposit a porphyrin molecule with a hydroxyphenyl side group on Au(111), which is designed and synthesized as a basic unit for functional single molecule devices, and observe the bonding structure and electronic states with a scanning tunneling microscope (STM). The molecule changes the configuration from monomer, cluster to monolayer as the coverage increases, ruled by the H-bonding interaction through the hydroxyphenyl group and steric repulsion by isoamiloxy groups. The highest occupied molecular orbital (HOMO) and the lowest unoccupied molecular orbital (LUMO) localized in the porphyrin macrocycle are observed at -1.1 and $+1.1$ eV, respectively, with respect to the Fermi level. We also deposit a *para*-phenylene-bridged porphyrin array on the surface by electrospray method, and observe the local density of states along the array.

1 Introduction

Porphyrins and the related macrocycles have rich optical and redox properties, which can be modulated by appropriate metalation. The flexible tunability of their properties makes them potential for molecular electronics applications^{1–5}. Recently, some of the authors reported a methodology for the synthesis of porphyrin arrays with programmed sequence⁶. Because the individual porphyrin macrocycles possess localized valence state of distinct energy levels, it is possible to engineer the distribution of the frontier orbitals along the array, and thereby to rationally design the transport property through the molecule. The electronic states of the arrayed molecules were studied by using optical and electrochemical methods, which revealed the highest occupied

molecular orbital (HOMO) and the lowest unoccupied molecular orbital (LUMO) levels of each component and their shift upon the formation of the arrays. However, these versatile techniques are not able to probe the electronic states in a site-resolved way; it is essential to investigate the spatial distribution of the frontier orbitals along the array for the control of the transport property through the molecules.

A scanning tunneling microscope (STM) is useful to investigate the electronic states and its spatial distribution for the molecules on the surface. The adsorption and self-assembly of porphyrin derivatives on surfaces have been intensively studied, mainly for the understanding of molecule-electrode interaction, which is essential for their application to molecular electronics^{7–13}. Not only the giant molecules^{14–16}, but also the functional molecules, such as molecular switches^{17,18} and molecular magnet¹⁹, were imaged and characterized at a single-molecule level. In this work, using STM, we image hydroxyphenyl porphyrin unit and its array which are synthesized as the basis of molecular rectifier, and characterize the electronic state associated with the transport property through the molecule.

^a Department of Chemistry, Graduate School of Science, Kyoto University, Kyoto 606-8502, Japan. Fax: +81 75 753 4000; Tel: +81 75 753 3977; E-mail: hokuyama@kuchem.kyoto-u.ac.jp

^b Department of Chemistry, Graduate School of Science, Osaka University, 1-1 Machikaneyama-cho, Toyonaka, Osaka 560-0043, Japan.

[‡] Present address: Department of Advanced Materials Science, The University of Tokyo, Kashiwa 277-8561, Japan.

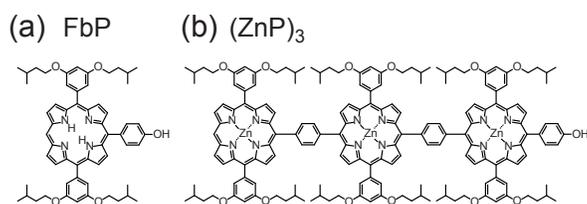


Fig. 1 Molecular structures of (a) freebase porphyrin monomer (FbP) and (b) phenylene-linked Zn porphyrin array [(ZnP)₃].

2 Experimental

The STM experiments were carried out in an ultrahigh-vacuum chamber (USM-1200, Unisoku). An electrochemically etched tungsten tip was used as an STM probe. The STM images were acquired in constant current mode at 6 K. Single-crystalline Au(111) was cleaned by repeated cycles of Ar⁺ sputtering and annealing.

10-(4-Hydroxyphenyl)-5,15-bis(3,5-diisopentoxyphenyl) freebase porphyrin [FbP; Fig. 1(a)] and *para*-phenylene-bridged porphyrin array of Zn porphyrin [(ZnP)₃; Fig. 1(b)] were synthesized according to the previous work⁶. The FbP molecules were thermally sublimated from a stainless crucible at ~540 K and deposited to the surface kept at ~90 K. We also used an electro-spray ionization method (ADCAP vacuum technology, Japan) to deposit nonvolatile (ZnP)₃ onto the surface²⁰. The molecule was dissolved in 10:1 methanol/toluene (by volume) mixture to give a concentration of 49 μM. The solution in the atmosphere was introduced into a differential pumping chamber at a typical flow rate of 2 μL/min through a needle tube with an applied voltage of 1.4 kV. The Au(111) surface was exposed to the ion flux (typical ion current of 10–30 pA) for 5 min at room temperature.

The *dI/dV* curves were obtained by using a lock-in amplifier with a modulation voltage of 8 mV_{rms} at 590 Hz. Each spectrum is displayed after subtraction of *dI/dV* obtained over the clean surface obtained at the constant tip height.

3 Results and Discussion

3.1 Porphyrin monomer

Figure 2 shows STM images of FbP molecules on Au(111) as a function of coverage. The molecules were adsorbed on the surface at ~90 K. At low coverage [Fig. 2(a)], the molecules are bonded to the elbow sites of the herringbone structure²¹, in a similar way to other porphyrin derivatives on the surface^{22,23}. The molecule surrounded by the dotted ellipse is assigned to an isolated molecule as described below, and their clusters are also observed at the elbow sites. In case the surface was exposed to the molecules at room temperature, they were trapped at the step edges on the surface (not shown). The molecules diffused across the surface and migrated to more favorable step sites at room temperature. Therefore it is essential to maintain the surface at

~90 K to isolate the molecules on the terrace of the surface. As the coverage increases, the molecules form islands in the face-centered-cubic region of the herringbone structure [Fig. 2(b)], and eventually forms a monolayer over the surface [Fig. 2(c)].

Figure 3(a) shows close inspection into an isolated molecule. The image consists of a ring surrounded by five spots indicated by the colored dots. By comparison with the molecular structure [Fig. 1(a)], we assign the center ring to a porphyrin macrocycle, and the spots indicated by the yellow and green dots to hydroxyphenyl (HO-C₆H₄-) and isoamyloxy [(CH₃)₂CHCH₂CH₂O-] groups, respectively. During successive scans, the appearance and position of the four isoamyloxy moieties (green dots) change independently, as shown in Figs. 3(a)–(c). The fluctuation is ascribed to the rotation of the isoamyloxy moieties around the sp³ C–C bonds, which is induced by the interaction with the STM tip; it depends on the tunnel condition (bias and current). We could not suppress the motion under the tunnel condition employed (*V* = 0.03–2 V and *I* = 10–500 pA).

Figures 3(d) and (e) show STM images of the molecule obtained at *V* = +1.5 and –1.5 V, respectively. Depending on the bias voltage, the appearance of the macrocycle drastically changes. It appears as rather uniform ring at low bias [Figs. 3(a)–(c)], but the upper part of the ring becomes more protruded at *V* = 1.5 V [Fig. 3(d)]. Note that the upper part of the image corresponds to the site to which hydroxyphenyl group is bonded. When the bias polarity is reversed (*V* = –1.5 V), the bright feature is localized at the specific position of the ring [the cross in Fig. 3(e)]. The dependence of the image on the bias voltage is ascribed to energy distribution of the electronic states of the molecule. Figure 3(g) shows a normalized *dI/dV* spectrum [i.e., ($\frac{dI}{dV}$) / ($\frac{I}{V}$) curve] obtained over the cross in Fig. 3(e). The density-functional theory (DFT) calculation for the free FbP molecule showed that LUMO and the second lowest unoccupied molecular orbital (LUMO+1) are distributed within 0.04 eV⁶ and thus may contribute to a peak at *V* = +1.1 V. HOMO and the second highest occupied molecular orbital (HOMO–1) are separated from each other by 0.4 eV for the free molecule⁶ and thus contribute to peaks at *V* = –1.1 and –1.6 V, respectively. Therefore, the STM images obtained at *V* = +1.5 V [Fig. 3(d)] and *V* = –1.5 V [Fig. 3(e)] reflect the shape of LUMO–LUMO+1 and HOMO–HOMO–1 of the molecule, respectively. The HOMO–LUMO gap of ~2.2 eV is narrower than the calculated value for the free molecule (2.7 eV)⁶. The reduction of the gap is possibly ascribed to surface polarization effects^{24–26}. The asymmetry of the image of porphyrin core with respect to the dashed line [Fig. 3(e)] reflects the presence/absence of the H atoms in the cavity of macrocycle [Fig. 3(f)]. It is also related with the deformation of the molecule; the porphyrin macrocycle undergoes saddle-shaped nonplanar deformation with pairs of opposite pyrrole rings tilted upwards or downwards^{17,22–24,27–30}. This deformation is associated with the rotated side groups by

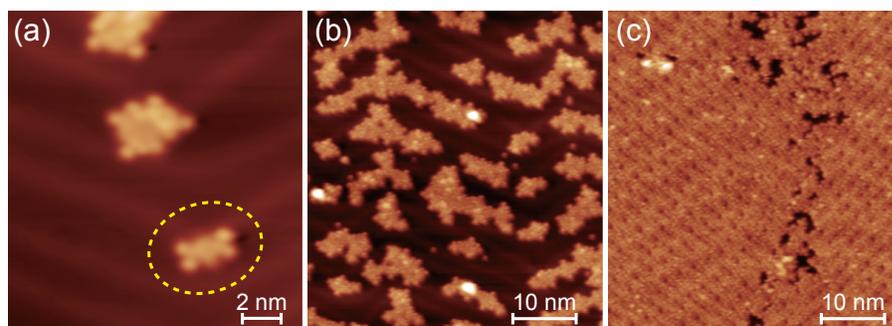


Fig. 2 (a)–(c) Typical STM images of freebase porphyrin monomers (FbP) on Au(111) as a function of the coverage. The molecules were adsorbed on the surface at ~ 90 K, and the images were obtained at 6 K in the constant-current mode with a sample bias $V = +0.4$ V and tunnel current $I = 50$ pA. An isolated molecule (surrounded by a dotted ellipse) is observed at low coverage in (a), and then, the clusters and monolayer develop as the coverage increases as shown in (b) and (c), respectively.

steric repulsion between phenyl and porphyrin rings. Note that the deformation of the macrocycle can also modify the HOMO–LUMO gap width.

Figure 4(a) shows the STM image of dimerized FbP molecules. The molecules preferentially form clusters and islands, as shown in Fig. 2(b), suggesting the presence of attractive interaction between the molecules. In Fig. 4(a), the molecules in the dimer are interacted with each other via hydroxyphenyl groups (head-to-head configuration), as illustrated in Fig. 4(b). In analogy to similar alkoxyphenyl-substituted cyanophenyl porphyrin^{22,28}, we attribute the interaction to H-bonding between the hydroxyl group and β -hydrogen of the adjacent porphyrin ring. At higher coverage the molecules form a monolayer, as shown in Fig. 4(d). On the terrace of the surface, the molecules are arranged in a quasi-rectangular arrangement with the hydroxyphenyl groups pointing random in opposite direction. This is consistent with the β -H \cdots OH interaction; the head-to-tail configuration (one H-bond) is nearly degenerate with simultaneous head-to-head (two H-bonds) and tail-to-tail (no H-bond) configurations. Therefore, we suggest that the structure of the monolayer is mainly determined by the interaction between the side isoamyloxy groups; the molecules are densely packed by avoiding steric repulsion between them.

3.2 Porphyrin array

A variety of metallated porphyrin units can be linked via phenylene in programmed sequence, enabling one to tailor the transport property through the molecule⁶. Here, as a first step toward the characterization of programmed arrays, we deposit a homologous porphyrin array [(ZnP)₃] on the surface and observe the electronic state at a level of individual molecules. Figure 5(a) shows typical STM images of (ZnP)₃ deposited by electro-spray method on Au(111). Isolated molecules and clusters are observed. Figure 5(b) shows a magnified image of an isolated

(ZnP)₃ molecule. The image is consistent with the molecular structure [Fig. 1(b)], as illustrated in Fig. 5(c). The individual macrocycles with four isoamyloxy groups are shown by the red, orange, and yellow envelopes, corresponding to the head, center and tail macrocycles, respectively. The residual protrusions (outside of the envelopes) were not reproducible, and thus, are ascribed to impurities attached to the molecule coadsorbed during the electro-spray deposition (possibly solvent molecules).

The electronic states of (ZnP)₃ are investigated with STM as shown in Fig. 5(d). The red, green, and blue curves shows normalized dI/dV curves obtained over the position marked by the crosses of corresponding colors in Fig. 5(e). The dI/dV shows the local density of states of the frontier orbitals at each site. As for the unoccupied states, a peak is observed at 1.3 V. In addition, low-voltage shoulder is prominent for the head molecule (red). The DFT calculation for the free (ZnP)₃ molecule showed that five orbitals (LUMO–LUMO+5) are distributed within 0.1 eV⁶ and thus may contribute to the peak. As for the occupied states, the dI/dV of the head and tail molecules shows a broad feature around -1 V, while that of the center molecule shows characteristic split structure. The calculation showed that HOMO–HOMO+5 are distributed within 0.3 eV⁶ and thus contribute to these structures. The site-dependent dI/dV is ascribed to the distribution of so many orbitals along the array. Note that these orbitals may be perturbed by the substrate as well as by the deformation (described later). Therefore DFT calculations for adsorbed (ZnP)₃ would help the interpretation of the dI/dV . Thus, the STM is useful for characterizing the local density of states along the molecule, which is associated with the transport property through the molecule.

The image obtained at $V = +1.5$ V [Fig. 5(e)] is contributed from several unoccupied orbitals, where the macrocycles show complex internal structure as depicted in Fig. 5(f). The occupied states are reflected in the image obtained at $V = -1.5$ V [Fig. 5(g)]. No sharp difference was visible between three macro-

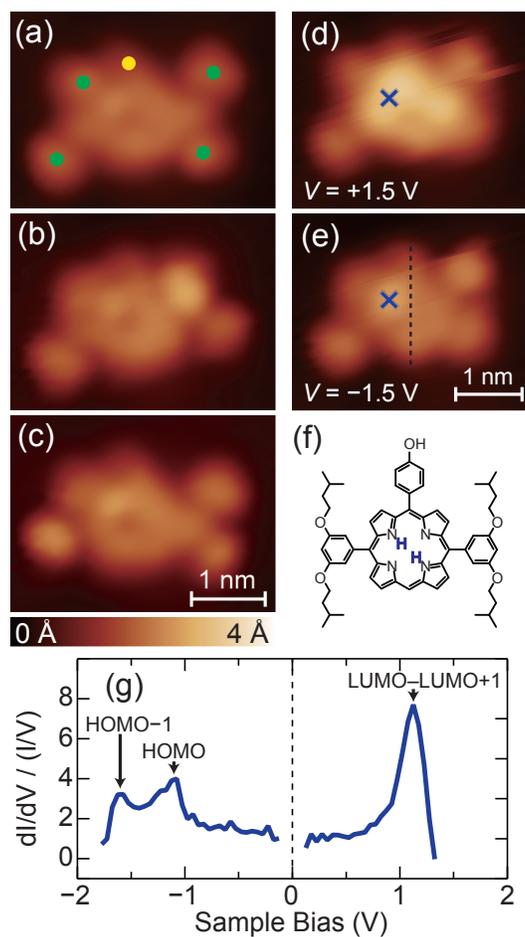


Fig. 3 (a)–(c) Sequential STM images of an isolated FbP molecule ($V = +0.4$ V and $I = 50$ pA). The yellow and green dots in (a) indicate the hydroxyphenyl and isoamyloxy groups, respectively. The appearance and position of the dots change between sequential imaging, which is caused by the inevitable perturbation by the tip. (d),(e) STM images of FbP acquired with $V =$ (d) $+1.5$ V and (e) -1.5 V at $I = 50$ pA. The porphyrin core appears asymmetric with respect to the molecular axis (dashed line). (f) Schematic illustration of the molecular structure of an FbP molecule. (g) A normalized dI/dV spectrum of the FbP molecule recorded over the cross in (e). The black arrows represent resonance peaks of the molecular orbitals.

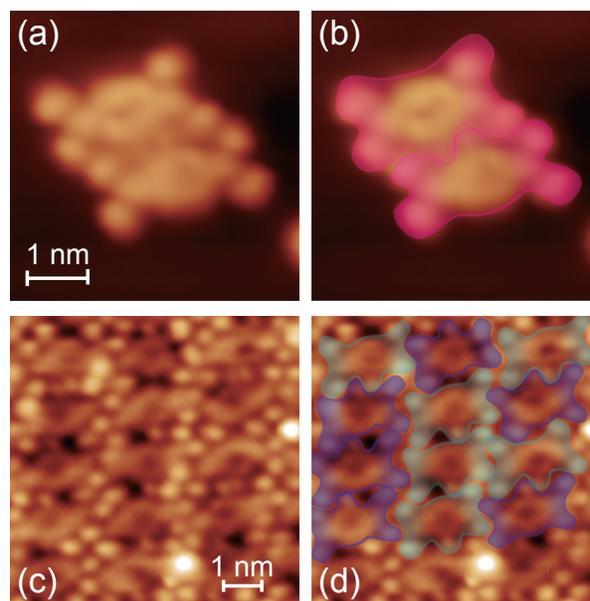


Fig. 4 (a) The STM images of dimerized FbP molecules ($V = +0.4$ V and $I = 50$ pA). (b) The same as (a) superimposed by the envelopes representing the contour of the molecules (head-to-head configuration). (c) The STM images of the monolayer FbP ($V = +0.4$ V and $I = 50$ pA). (d) The same as (b) superimposed by the envelopes representing the contour of the molecules. Blue and cyan envelopes represent the molecules which are rotated by 180° with each other.

cycles in the range of $V = +1.1$ – $+1.5$ V (-1.0 – -1.5 V), because of (i) similar local density of states [Fig. 5(d)] and (ii) delocalized nature of the orbital along the array.

The illustration of the occupied-state image [Fig. 5(h)] shows that each porphyrin macrocycle is no longer symmetric with respect to the molecular axis. The Zn atoms are positioned at the center of the cavities, and thus, the asymmetry is due to the deformation of the macrocycles. Because of the steric repulsion between the isoamyloxy moieties of linked macrocycles, the phenyl ring may twist and the isoamyloxy groups tilt to release the repulsion. This causes the saddle-shaped deformation of the macrocycle as described above. The two salient protrusions in the molecular image [Fig. 5(b)] is thus ascribed to the isoamyloxy groups tilted upward [see the orange envelope in Fig. 5(c)].

4 Conclusion

We image a freebase porphyrin monomer (FbP) and zinc porphyrin array $[(ZnP)_3]$ on Au(111), and characterize their electronic states with STM. The nonvolatile porphyrin array is successfully deposited on the surface using an electrospray deposition. The STM reveals the energy level and spatial distribution of the frontier orbitals localized on the individual porphyrin macrocycles in the molecules. The real-space imaging is also useful for studying their aggregation on the surface. Combined with

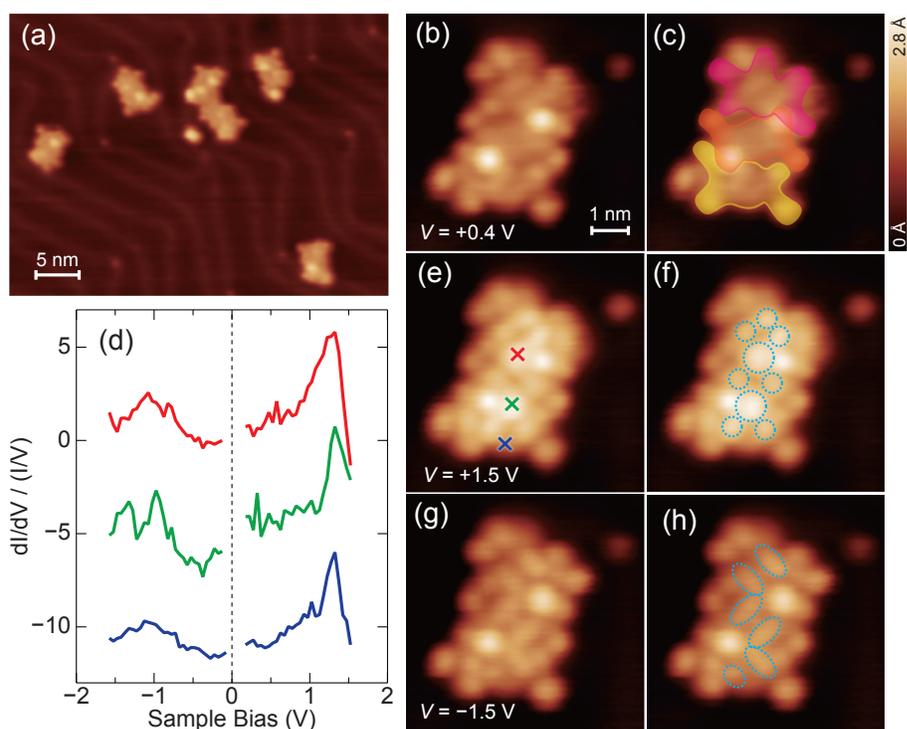


Fig. 5 (a) Typical STM images of zinc porphyrin array $[(\text{ZnP})_3]$ on Au(111) ($V = +0.4$ V and $I = 50$ pA). (b) A close-up view of a porphyrin array. (c) The same as (b) superimposed by schematic illustration of each component. (d) Normalized dI/dV curves of the array (offset for clarity). The red, green, and blue curves are recorded over the position marked by the crosses of corresponding colors in (e). The STM images of a $(\text{ZnP})_3$ molecule acquired with (e) $V = 1.5$ V and (g) $V = -1.5$ V at $I = 50$ pA. The bias-dependent STM images in (e) and (g) show unoccupied orbitals and occupied orbitals, respectively. The images (f) and (h) are the same as (e) and (g), respectively, but superimposed by the ellipses illustrating the shape of the orbitals distributed in each macrocycle.

nondestructive deposition method, the STM can probe the local electronic states along the programmed porphyrin arrays, which is associated with the transport property through the molecule.

This study was supported in part by Grant-in-Aid for Scientific Research on Innovative Areas “Molecular Architectonics: Orchestration of Single Molecules for Novel Functions.” A.S. acknowledges the support of the Japan Society for the Promotion of Science. The authors acknowledge T. Yokoyama for technical support of electrospray deposition.

References

- 1 M. Jurow, A. E. Schuckman, J. D. Batteas and C. M. Drain, *Coord. Chem. Rev.*, 2010, **254**, 2297–2310.
- 2 H. Xu, R. Chen, Q. Sun, W. Lai, Q. Su, W. Huang and X. Liu, *Chem. Soc. Rev.*, 2014, **43**, 3259–3302.
- 3 K. Tagami, M. Tsukuda, T. Matsumoto and T. Kawai, *Phys. Rev. B*, 2003, **67**, 245324.
- 4 J. A. A. W. Elemans, R. van Hameren, R. J. M. Norte and A. E. Rowan, *Adv. Mater.*, 2006, **18**, 1251–1266.
- 5 G. Sedghi, V. M. García-Suárez, L. J. Esdaile, H. L. Anderson, C. J. Lambert, S. Martin, D. Bethell, S. J. Higgins, M. Elliott, N. Bennett, J. E. Macdonald and R. J. Nichols, *Nat. Nanotech.*, 2011, **6**, 517–523.
- 6 T. Tamaki, T. Nosaka and T. Ogawa, *J. Org. Chem.*, 2014, **79**, 11029–11038.
- 7 J. Otuski, *Coord. Chem. Rev.*, 2010, **254**, 2311–2341.
- 8 S. Mohnani and D. Bonifazi, *Coord. Chem. Rev.*, 2010, **254**, 2342–2362.
- 9 S. Yoshimoto and N. Kobayashi, *Struct. Bond.*, 2010, **135**, 137–168.
- 10 J. Xiao, S. Ditze, M. Chen, F. Buchner, M. Stark, M. Drost, H.-P. Steinrück, J.M. Gottfried and H. Marbach, *J. Phys. Chem. C*, 2012, **116**, 12275–12282.
- 11 T. Niu and A. Li, *J. Phys. Chem. Lett.*, 2013, **4**, 4095–4102.
- 12 S. Haq, F. Hanke, J. Sharp, M. Persson, D.B. Amabilino and R. Raval, *ACS Nano*, 2014, **9** 8856–8870.
- 13 W. Auwärter, D. Écija, F. Klappenberger and J. V. Barth, *Nat. Chem.*, 2015, **7**, 105–120.
- 14 A. Saywell, J. K. Sprafke, L. J. Esdaile, A. J. Britton, A. Rienzo, H. L. Anderson, J. N. O’Shea and P. H. Beton, *Angew. Chem. Int. Ed.*, 2010, **49**, 9136–9139.
- 15 M.C. O’Sullivan, J. K. Sprafke, D. V. Kondratuk, C. Rinfray, T. D. W. Claridge, A. Saywell, M. O. Blunt, J. N. O’Shea, P. H. Beton, M. Malfois and H. L. Anderson, *Nature*, 2011, **469**, 72–75.
- 16 D. V. Kondratuk, L. M. A. Perdigo, M. C. O’Sullivan, S. Svatek, G. Smith, J. N. O’Shea, P. H. Beton and H. L. Anderson, *Angew. Chem. Int. Ed.*, 2012, **51**, 6696–6699.
- 17 W. Auwärter, K. Seufert, F. Bischoff, D. Ecija, S. Vijayaraghavan, S. Joshi, F. Klappenberger, N. Samudrala and J. V. Barth, *Nat. Nanotech.*, 2012, **7**, 41–46.
- 18 S. Ditze, M. Stark, F. Buchner, A. Aichert, N. Jux, N. Luckas, A. Görling, W. Hieringer, J. Hornegger, H.-P. Steinrück and H. Marbach, *J. Am. Chem. Soc.*, 2014, **136**, 1609–1616.
- 19 W. Wang, R. Pang, G. Kuang, X. Shi, X. Shang, P.N. Liu and N. Lin, *Phys. Rev. B*, 2015, **91** 045440.
- 20 T. Yokoyama, Y. Kogure, M. Kawasaki, S. Tanaka and K. Aoshima, *J. Phys. Chem. C*, 2013, **117**, 18484–18487.
- 21 J.V. Barth, H. Brune, G. Ertl and R.J. Behm, *Phys. Rev. B*, 1990, **42**, 9307–9318.
- 22 T. Yokoyama, S. Yokoyama, T. Kamikado, Y. Okuno and S. Mashiko, *Nature*, 2001, **413**, 619–621.
- 23 T. Yokoyama, S. Yokoyama, T. Kamikado and S. Mashiko, *J. Chem. Phys.*, 2001, **115**, 3814–3818.
- 24 A. Weber-Bargioni, W. Auwärter, F. Klappenberger, J. Reichert, S. Lefrançois, T. Strunskus, C. Wöll, A. Schiffrin, Y. Pennec and J. V. Barth, *Chem. Phys. Chem.*, 2008, **9**, 89–94.
- 25 J. B. Neaton, M. S. Hybertsen and S. G. Louie, *Phys. Rev. Lett.*, 2006, **97**, 216405.
- 26 G. Heimel and J.-L. Bredas, *Nat. Nanotech.*, 2013, **8**, 230–231.
- 27 F. Buchner, K. Comanici, N. Jux, H.-P. Steinrück and H. Marbach, *J. Phys. Chem. C*, 2007, **111**, 13531–13538.
- 28 N. Wintjes, J. Hornung, J. Lobo-Checa, T. Voigt, T. Samuely, C. Thilgen, M. Stöhr, F. Diederich and T. A. Jung, *Chem. Eur. J.*, 2008, **14**, 5794–5802.
- 29 W. Auwärter, K. Seufert, F. Klappenberger, J. Reichert, A. Weber-Bargioni, A. Verdini, D. Cvetko, M. Dell’Angela, L. Floreano, A. Cossaro, G. Bavdek, A. Morgante, A. P. Seitsonen and J. V. Barth, *Phys. Rev. B*, 2010, **81**, 245403.
- 30 T. Yokoyama and F. Nishiyama, *J. Phys. Chem. Lett.*, 2014, **5**, 1324–1328.