Magnesium gating of cardiac gap junction channels

Hiroyuki Matsuda¹, Yasutaka Kurata², Chiaki Oka¹, Satoshi Matsuoka¹, and Akinori Noma¹

¹Department of Physiology and Biophysics, Graduate School of Medicine, Kyoto University, Yoshida-konoe, Sakyo-ku, Kyoto 606·8501, Japan
²Department of Physiology, Kanazawa Medical University, Ishikawa 920·0293, Japan

Present address:
Hiroyuki Matsuda; Nano Medicine Merger Education Unit, Graduate School of Medicine, Kyoto University, Yoshida-konoe, Sakyo-ku, Kyoto 606·8501, Japan.
Satoshi Matsuoka; Center for Innovation in Immunoregulative Technology and Therapeutics, Graduate School of Medicine, Kyoto University, Yoshida-konoe, Sakyo-ku, Kyoto 606·8501, Japan.
Akinori Noma; College of Life Sciences, Ritsumeikan University, 1·1·1 Noji Higashi, Shiga 525·8577, Japan.

Corresponding Author: Satoshi Matsuoka:
Center for Innovation in Immunoregulative Technology and Therapeutics, Graduate School of Medicine, Kyoto University, Yoshida-konoe, Sakyo-ku, Kyoto 606·8501, Japan.
Tel.: +81·75·753·9507
Fax: +81·75·753·9500
E-mail: smatsuoka@ak.med.kyoto-u.ac.jp
Abstract

We aimed to study kinetics of modulation by intracellular Mg$^{2+}$ of cardiac gap junction (Mg$^{2+}$ gate). Paired myocytes of guinea-pig ventricle were superfused with solutions containing various concentrations of Mg$^{2+}$. In order to rapidly apply Mg$^{2+}$ to one aspect of the gap junction, the non-junctional membrane of one of the pair was perforated at nearly the connecting site by pulses of nitrogen laser beam. The gap junction conductance ($G_j$) was measured by clamping the membrane potential of the other cell using two-electrode voltage clamp method. The laser perforation immediately increased $G_j$, followed by slow $G_j$ change with time constant of 3.5 sec at 10 mM Mg$^{2+}$. Mg$^{2+}$ more than 1.0 mM attenuated dose-dependently the gap junction conductance and lower Mg$^{2+}$ (0.6 mM) increased $G_j$ with a Hill coefficient of 3.4 and a half-maximum effective concentration of 0.6 mM. The time course of $G_j$ changes was fitted by single exponential function, and the relationship between the reciprocal of time constant and Mg$^{2+}$ concentration was almost linear. Based on the experimental data, a mathematical model of Mg$^{2+}$ gate with one open state and three closed states well reproduced experimental results. One-dimensional cable model of thirty ventricular myocytes connected to the Mg$^{2+}$ gate model suggested a pivotal role of the Mg$^{2+}$ gate of gap junction under pathological conditions.

Key Words:

Gap junction channel, chemical gating, numerical ventricular cell model, arrhythmia, Mg$^{2+}$
1 Introduction

Gap junction channels (GJCs) provide a pathway of low electrical resistance between the adjacent cells and play a pivotal role in synchronized contraction of the heart. Conductance of GJCs is regulated by the trans-junctional potential (voltage gate), phosphorylation as well as several intracellular molecules (chemical gate), such as $\text{H}^+$ and $\text{Ca}^{2+}$ (Bukauskas et al. 2004; Peracchia et al. 2004; Desplantez et al. 2007). Under some circumstances, such as ischemia, GJCs close and cellular communication is uncoupled electrically and metabolically, separating damaged cells from intact cells. Thus, GJCs work as fire shutters.

Mg$^{2+}$ is one of the intracellular molecules that modulate GJCs. Free intracellular Mg$^{2+}$ concentration ($[\text{Mg}^{2+}]_i$) is maintained within a narrow range of 0.5 - 1 mM (Grubbs 2002). Under ischemic conditions, $[\text{Mg}^{2+}]_i$ increases because of decrease in ATP (Carmeliet 1999; Murphy et al. 1989) that is the major buffer of Mg$^{2+}$ (Michailova et al. 2001; Grubbs 2002). Intracellular Mg$^{2+}$ regulates several ion channels in cardiac myocyte (Gusev et al. 2008; Wang et al. 2006; Vemana et al. 2008; Ishihara et al. 2009; Shi et al. 2002). Therefore, Mg$^{2+}$ may be one of key factors in ischemic arrhythmia like $\text{H}^+$ and $\text{Ca}^{2+}$ (Swietach et al. 2007). Mg$^{2+}$ gating of GJCs was first reported in salivary gland cells by Oliveira-Castro and Loewenstein (1971). Noma and Tsuboi (1987) confirmed the Mg$^{2+}$ effect and steady state Mg$^{2+}$ dependence of GJCs conductance was studied in adult guinea-pig ventricular myocytes using ‘open cell’ method. However, no significant effect of Mg$^{2+}$ was observed on the GJCs conductance of cultured neonatal ventricular myocytes of rat (Firek et al. 1995). No further systemic studies on the Mg$^{2+}$ gating of GJCs have been carried out.

In this study, we used ‘laser beam perforation’ method for rapidly applying $[\text{Mg}^{2+}]_i$ to GJCs and two-electrode voltage clamp method for accurately controlling membrane voltage. A numerical model was constructed based on the gating kinetics.

2 Material and Methods

2.1. Isolation of guinea-pig ventricular myocytes

Paired ventricular cells were prepared according to previous studies (Kameyama. 1983; Mitra et al. 1985). Briefly, female guinea-pigs (4-weeks-old) were
anaesthetized by an intraperitoneal injection of sodium pentobarbital (50mg/kg). The heart was quickly excised after the open of thoracic cavity, and suspended on a Langendorff perfusion apparatus. The flow rate was approximately 10 mL min⁻¹. The isolated heart was perfused initially with a control Tyrode solution for 10 to 15 minutes, secondary with a Ca²⁺-free Tyrode solution for 5 minutes, and thirdly with a Ca²⁺-free Tyrode solution containing collagenase (50 U·ml⁻¹, Warthingtone type I) for 20 minutes. Finally, the heart was perfused with a Kraftbrühe (KB) medium for 10 minutes. These procedures were performed at 37°C. Myocytes were gently dispersed by mechanical agitation. Isolated cells including paired cells were kept in the KB medium at 4°C for at least 1 hour before use in experiments.

2.2. Whole-cell voltage clamp using double giga-sealed patch electrodes

Isolated cells were transferred to a chamber mounted on an inverted microscope (Axiovert135; Zeiss, Germany). Paired cells with approximately same size were used for experiments. To accurately control membrane potential, the two-electrode voltage clamp (TEVC) method was applied (Axoclamp 2B amplifier; Axon instruments, USA) with the essentially same protocol as Noma and Tsuboi (1987). Briefly, two patch pipettes filled with an internal solution were prepared (electrical resistance = 2~5 MΩ) on separate micro manipulators. After the two pipettes were attached on the cell surface of one of the paired cells (cell 1), giga ohm seal and whole-cell configuration were established by applying negative pressure of 20-50 cm H₂O. A schematic diagram of voltage clamp circuit was shown in Fig.1A.

2.3. Experimental protocol

After replacing the control Tyrode solution with a Cs⁺-rich intracellular solution containing a given concentration of Mg²⁺, the membrane of the paired cell (cell 2) close to cell 1 was instantaneously perforated by applying pulses of nitrogen laser (VSL-337ND; Laser Science, USA). The distance between the cell connection site and the holes were within about 15 μm. Usually about 10 holes were made within 1 sec. At the end of experiments, the cells were superfused with a solution containing 0.1 mM Mg²⁺ to normalize input conductance. A series of test potential to ± 5 mV from a holding potential of 0 mV was applied every 500 ms (Fig. 1B). This voltage protocol is unlikely to affect the voltage gating of GJCs because conductance of homomeric-homotypic Cx43 channels, which is a
dominant isoform of connexins in ventricular myocardium, is almost constant at voltage from -30 to 30 mV (González et al. 2007). The equivalent electrical circuits before and after the laser perforation are illustrated in Fig. 1C. The GJCs conductance ($G_j$) was obtained on the following assumptions as previous studies (Noma and Tsuboi 1987; Sugiura et al. 1990; Toyama et al. 1994).

Before the laser perforation, input conductance ($G_i$) is expressed as follows

\[ G_i = G_{m1} + G_j \cdot G_{m2} / (G_j + G_{m2}), \]

where $G_{m1}$ and $G_{m2}$ are non-junctional membrane conductance of cell 1 and cell 2, respectively. If $G_{m1}$ is equal to $G_{m2}$ and $G_i$ is much greater than $G_{m1}$ and $G_{m2}$ ($G_j \gg G_m$), $G_i$ will be simplified.

\[ G_i = 2 \cdot G_m, \quad G_m = G_{m1} = G_{m2} \]

(1)

After the perforation, $G_{m2}$ could be ignored.

\[ G_i = G_{m1} + G_j / (G_j + G_s), \]

where $G_s$ is the conductance through the holes made in cell 2. If $G_s$ is much larger than $G_j$, $G_i$ will be obtained with following equation

\[ G_i = G_m + G_j, \quad G_m = G_{m1} \]

(2)

$G_j$ can be calculated with equations (1) and (2).

In this study, non-junctional membrane conductance was suppressed by a K$^+$ free, Ca$^{2+}$ free and low Cl$^-$ solution. Contribution of the H$^+$ and Ca$^{2+}$ gatings of GJC might be negligible because the Ca$^{2+}$ and H$^+$ was strongly buffered with 5 mM EGTA and HEPES (see below).

2.4. Solutions

The composition of Tyrode solution was (in mM): 140 NaCl, 5.4 KCl, 0.5 MgCl$_2$, 1.8 CaCl$_2$, 0.33 NaH$_2$PO$_4$, 5.5 glucose and 5 Hepes; pH was adjusted to 7.4 with NaOH. KB solution contained (in mM): 90 KOH, 30 KCl, 10 KH$_2$PO$_4$, 70 L-glutamic acid, 20 taurine, 10 HEPES, 10 D-glucose, 0.5 EGTA, 1 MgCl$_2$; pH was adjusted to 7.4 with KOH. The internal solution in the patch electrode contained (in mM) 110 CsOH, 20 TEACl, 5 Na$_2$ATP, 5 HEPES, 5 EGTA, 5 Na$_2$CRP, 110 Aspartic acid, 3 MgCl$_2$ (pMg = 4); pH was adjusted to 7.4 with CsOH. The composition of the test solution was the same as the internal solution, except for MgCl$_2$. Free Mg$^{2+}$ concentration (0.1-10 mM) was calculated according to a software ‘WinMaxC’ (http://www.stanford.edu/~cpatton/maxc.html).
2.5. Model construction

A mathematical model of gap junction was constructed using Delphi 2009 (Embarcadero technologies, USA) and differential equations were solved with fourth-order Runge–Kutta method.

Propagation of action potential was examined with a one-dimensional (1-D) cable model, which consisted of thirty ventricular cell models connected to gap junctions. The guinea-pig ventricular cell model published by kuzumoto et al. (2008) was used, where cell volume was fixed to a constant value for simplicity. Open probability ($P_{\text{Open},i}$) of a gap junction connecting cell (i) and cell (i+1) is expressed as

$$P_{\text{Open},i} = P_i \cdot P_{i+1}$$

$$G_{j,i} = N \cdot g_j \cdot P_{\text{Open},i}$$

where $P_i$ and $P_{i+1}$ are the open probabilities of the hemichannels facing $[\text{Mg}^{2+}]_i$ of cell (i) and cell (i+1), respectively. $N$ is the number of channels, $g_j$ the maximal conductance of the unitary channel. Since the half-life of Cx43 was hours order in adult heart (Beardslee et al. 1998), and cation such as H$^+$ or Ca$^{2+}$ do not affect maximal conductance (Brink et al. 1989; Veenstra et al. 1988), we assumed that $N$ and $g_j$ were constant and set $N \cdot g_j$ to 9000 nS to yield normal conductance velocity (450 – 700 mm/sec, Eloff et al. 2003; Harada et al. 2008; Rossi et al. 2008; Akar et al. 2007) at physiological $[\text{Mg}^{2+}]_i$.

The membrane potential of cell (i) ($V_{m,i}$) in a 1-D cable model was calculated as follows.

$$\frac{dV_{m,i}}{dt} = - \left( I_{\text{tot},i} + I_{j,i} - I_{j,i-1} \right) / C_m$$

$$I_{j,i} = G_{j,i} \cdot (V_{m,i+1} - V_{m,i})$$

Here $C_m$ and $I_{\text{tot},i}$ are membrane capacitance and total membrane current of cell (i), respectively. $I_{j,i}$ and $G_{j,i}$ are the current and conductance of a gap junction connecting cell (i) and cell (i+1), respectively. For simplicity, the electrical resistances of the intra- and extracellular medium were neglected. The model was excited at 2.5 Hz by applying to cell (i) a stimulus pulse, which was 2.0 ms in duration and 20 nA in amplitude. Time step was 0.005 msec.

The conduction velocity (CV) between cell (i) and cell (i+1) was calculated by a following equation

$$CV = L / \left( t_{0,i+1} - t_{0,i} \right)$$

Here the cell length (L) is set at 100 μm, and $t_{0,i}$ represents the time of AP rising over 0 mV in cell (i). In simulation, only $[\text{Mg}^{2+}]_i$, which affects gap junction was changed when changing the concentration, to clarify the effect of $[\text{Mg}^{2+}]_i$ on gap junction.
2.6. Data analysis and equipments

Model equations were fitted to data by a non-linear, least squares method using Sigmaplot 10.0 (Systat Software Inc, CA, USA).

3 Results

3.1. Mg\textsuperscript{2+} dependence of GJC conductance

Fig. 2 shows a typical experimental result. Representative current traces and time course of \(G_j\) are shown at A and B, respectively. A pair of myocytes was superfused with a solution containing 10 mM Mg\textsuperscript{2+}. Laser perforation induced a rapid jump-up of \(G_j\) (b), followed by slower decrease (c). The \(G_j\) immediately after the laser perforation reflects \(G_j\) at physiological condition. The following decrease in \(G_j\) was probably related to closing of the hemichannels of cell 2 by 10 mM Mg\textsuperscript{2+}, because changing the perfusion solution to that containing 0.1 mM Mg\textsuperscript{2+} (slow perfusion) increased \(G_j\) by about 500 times. Thus, \(G_j\) was clearly dependent on [Mg\textsuperscript{2+}].

3.2. Time course of \(G_j\) modulation by Mg\textsuperscript{2+}

In the majority of experiments, no significant delay was observed in the \(G_j\) change after the laser perforation. However, if the laser holes were made away from the site of cell connection, significant delay was observed. In Fig.3, laser beam were applied at 11, 18 and 48 \(\mu\)m from the junction. At 48 \(\mu\)m, delay of about 5 sec was observed and decay speed was markedly slower than the ones at 11 and 18 \(\mu\)m. The delay was presumably caused by the intracellular diffusion of Mg\textsuperscript{2+}. The \(G_j\) decay without delay suggested that Mg\textsuperscript{2+} gate does not include multiple open states.

The \(G_j\) changes after the laser perforation could be fitted by single exponential function as shown in Fig.4. When a solution containing lower [Mg\textsuperscript{2+}], was used (0.6 mM), \(G_j\) exponentially increased after the laser perforation. This finding is consistent with data in Fig.2 and further confirms a view that \(G_j\) depends on [Mg\textsuperscript{2+}].

3.3. Dependencies of \(\tau^{-1}\) and steady-state \(G_j\) on Mg\textsuperscript{2+}

Steady state dependence of \(G_j\) on [Mg\textsuperscript{2+}], was fitted to a Hill equation with a Hill coefficient of 3.4 and a half-maximum effective concentration (EC\textsubscript{50}) of 0.6 mM (Fig.4C). The value of EC\textsubscript{50} was smaller than the one reported by Noma and Tsuboi (1987), but the Hill coefficient was similar. High cooperativity of Mg\textsuperscript{2+} suggested that Mg\textsuperscript{2+} more than 3 binds to the hemichannel. Fig.4D shows a relationship
between \([\text{Mg}^{2+}]\), and reciprocal time constant of the \(G_j\) change \((\tau^{-1})\). A linear relationship was found.

### 3.4. Construction of \(\text{Mg}^{2+}\) gate model

We constructed a \(\text{Mg}^{2+}\) gating model of gap junction as follows. If the \(\text{Mg}^{2+}\) gate of the hemichannel is followed by a simple two-state model including one open (O) and one closed state (C, \(C = 1 – O\)) (Fig. 5A), time derivative \((dO/dt)\), steady state open probability \((P_{open,ss})\) and rate constant \((\tau^{-1})\) are given as follows:

\[
\frac{dO}{dt} = -\alpha \cdot [\text{Mg}^{2+}]^n \cdot O + \beta \cdot C = - \left(\alpha \cdot [\text{Mg}^{2+}]^n + \beta\right) \cdot O + \beta \\
P_{open,ss} = \frac{\beta}{\left(\alpha \cdot [\text{Mg}^{2+}]^n + \beta\right)} \\
\tau^{-1} = \alpha \cdot [\text{Mg}^{2+}]^n + \beta,
\]

where \(\alpha\), \(\beta\) and \(n\) represent forward rate constant, backward one and a number of \(\text{Mg}^{2+}\) molecule, respectively. The Hill coefficient obtained in Fig 4C indicated that \(n\) is more than 3. However, since \(\tau^{-1}\) depends on \([\text{Mg}^{2+}]^n\) in this model, the linear relationship between \(\tau^{-1}\) and \([\text{Mg}^{2+}]\) obtained in Fig. 4D could not be reproduced. If we assume that 3 \(\text{Mg}^{2+}\) bind to the hemichannel, four states may exist. Since one single exponential function fit well to the time course of \(\text{Mg}^{2+}\) gate, transition between open and close states may be slow (rate limiting) and other steps could be assumed as instantaneous reaction. The possible combination is listed in Fig. 5B.

Differential equation of the open probability can be described as follows.

\[
\frac{dO_r}{dt} = -Q \cdot O_r + R,
\]

where \(O_r\) is a summation of all open states, \(Q\) rate constant \((\approx \tau^{-1})\), R a function of \(\alpha_1, \alpha_2, \alpha_3, \beta_1, \beta_2, \beta_3\), and \([\text{Mg}^{2+}]\).

If 1\(^{st}\) step is rate limiting (Fig. 5B1),

\[
Q = \alpha_1 [\text{Mg}^{2+}] + \beta_1/(1+(\alpha_2/\beta_2)[\text{Mg}^{2+}])^2
\]

If 2\(^{nd}\) step is rate limiting (Fig. 5B2),

\[
Q = \alpha_2 [\text{Mg}^{2+}] /((\alpha_1/\beta_1) [\text{Mg}^{2+}])^{1+1} + \beta_2/(1+(\alpha_3/\beta_3)[\text{Mg}^{2+}])
\]

If 3\(^{rd}\) step is rate limiting (Fig. 5B3),

\[
Q = \alpha_3 [\text{Mg}^{2+}] /((\alpha_1/\beta_1 + (\alpha_2/\beta_2) [\text{Mg}^{2+}]^{1+1} + \beta_3/\alpha_1 \alpha_2) [\text{Mg}]^{2}) + \beta_3
\]

Only equation (3) has a term which depends linearly on \([\text{Mg}^{2+}]\). The linear relationship between \(\tau^{-1}\) and \([\text{Mg}^{2+}]\) could not be reproduced with other models. Furthermore, The \(G_j\) decay without delay (Fig. 3) also supports the model with one open state. Therefore, we adopted the model of Fig. 5B1.

To further simplify this model, we assumed forward \((k_f)\) and backward \((k_b)\) rate constants of single \(\text{Mg}^{2+}\) binding within closed state. Thus, \(K_d = k_f/k_b\), \(\alpha_2 = 2k_f\), \(\beta_2 = k_b\), \(\alpha_3 = k_f, \beta_3 = 2k_b\). The differential equation, rate constant and steady state open probability of model Fig. 5B1 are described as follows.
\[
d\frac{[O]}{dt} = -\{\alpha_1 [\text{Mg}^{2+}] + \beta_1 /(1+ K_d [\text{Mg}^{2+}])^2\} [O] + \beta_1 /((\alpha_1 [\text{Mg}^{2+}] (1+ K_d [\text{Mg}^{2+}]))^2 + \beta_1)
\]
\[
Q = \alpha_1 [\text{Mg}^{2+}] + \beta_1 /(1+ K_d [\text{Mg}^{2+}])^2
\]
\[
P_{\text{open ss}} = \beta_1 /((\alpha_1 [\text{Mg}^{2+}] (1+ K_d [\text{Mg}^{2+}]))^2 + \beta_1)
\]

Parameters, \( \alpha_1, \beta_1 \), and \( K_d \) were determined by fitting the experimental data to equations (4) and (5) (Fig 6A and B), where \( \alpha_1 = 0.028 \text{ S}^{-1}\text{mM}^{-1} \), \( \beta_1 = 0.227 \text{ S}^{-1} \), and \( K_d = 0.2372 \text{ mM}^{-2} \), respectively. To simulate the experiments, initial [Mg\(^{2+}\)]\(_i\) just after the laser perforation was estimated with the relationship between [Mg\(^{2+}\)]\(_i\) and \( G_j \) in A; 1.0 mM in C and 0.95 mM in D. Mg\(^{2+}\) was changed from the initial values to 10 mM (C) and 0.6 mM (D) in simulation. The model well reproduced the time course of Mg\(^{2+}\) gating.

Propagation of action potential at various [Mg\(^{2+}\)]\(_i\) was examined with a model of tandemly-arrayed ventricular cells (Fig. 7A). Conductance velocity was sharply dependent on [Mg\(^{2+}\)]\(_i\). At the physiological range of [Mg\(^{2+}\)]\(_i\) in cardiac cells (0.5 - 1 mM), conductance velocity was 645 – 224 mm/sec (Fig. 7B). These values are close to experimental data (450 mm/sec (Akar et al. 2007), 472 mm/sec (Eloff et al. 2003), 600 mm/sec (Harada et al. 2008) and 690 mm/sec (Rossi et al. 2008)). Action potentials of the thirty cells at 0.1, 0.7, and 2.0 mM [Mg\(^{2+}\)]\(_i\) are demonstrated in Fig. 7C. Marked conduction delay were observed at [Mg\(^{2+}\)]\(_i\) more than 1.4 mM. Action potential did not propagate to second cell at [Mg\(^{2+}\)]\(_i\) more than 2.8 mM.

The time course of action potential modulation by Mg\(^{2+}\) was studied in Fig. 8, Control [Mg\(^{2+}\)], of all the cells was 0.7 mM. At time of 4.0 sec, [Mg\(^{2+}\)]\(_i\) of cell 11 to 30 was increased to 2.6 mM. The open probability of cell 11 to 30 decreased over several tens of seconds (Fig. 8A). Action potential propagation remarkably slowed at cell 11 to 30 and two-to one block appeared (Fig. 8B).

4 Discussion

4.1. Measurements of GJC conductance with laser perforation technique

To study kinetics of chemical gate of GJC in native ventricular cells, fast application of agents to GJC is prerequisite. Although several methods have been developed (Harris et al. 2001), all the methods have limitations for applying to intact ventricular cells. Noma and Tsuboi (1987) developed ‘open-cell’ method, in which the non-junctional membrane of one of the paired cells was mechanically ruptured with a glass pipette to allow agents to reach one aspect of the junction.
However, diffusion delay was marked (> several tens of seconds), so that accurate rate constants of chemical gate could not be obtained. In this study, we used laser beam to make holes near the GJCs of paired cells. The input conductance immediately increased upon applying laser beam (Fig.2) and no marked diffusion delay was observed (Fig.3). Therefore, this new method is suitable for studying kinetics of the chemical gate in native ventricular cells.

4.2. What is the mechanism of Mg$^{2+}$ action on GJCs?

Direct gating of connexin channels by ligands is of key importance in intercellular signaling. However, it is not clear whether the cation gate is direct or indirect reaction to GJCs (Dhein. 2004). Although it is suggested that the Ca$^{2+}$ effect on GJCs is mediated by an intermediate component such as calmodulin (Lurtz et al. 2007), there is no information about Mg$^{2+}$ gating mechanism of gap junction channel. Mubagwa et al. (2007) discussed the mechanistic patterns of intracellular Mg$^{2+}$ action on several cation channels; namely direct binding, allosteric effects, modulation via action on enzymes or G proteins, interaction with surface charges, interaction with membrane phospholipids, and others. In this paper we assumed that Mg$^{2+}$ directly binds to GJCs. However, we could not rule out the possibility that the effect was mediated by other molecules. Sugiura et al. (1990) reported that gap junction conductance was increased by ATP, presumably by Mg-ATP, with Hill coefficient of 2.6 and EC$_{50}$ of 0.68 mM. However, in our experiment, total ATP concentration was constant (5 mM) and calculated Mg-ATP concentration was almost similar level among the test solutions: 2.86 mM at the 0.1 mM Mg$^{2+}$ solution and 4.95 mM at the 10 mM Mg$^{2+}$ solution. Therefore, the Mg$^{2+}$ effect observed in this study is unlikely to be mediated by ATP or Mg-ATP.

4.3. Physiological significance of the Mg$^{2+}$ gating

Referring to cardiovascular disease, [Mg$^{2+}$], increases in acute myocardial ischemia and decreases in heart failure. Murphy et al. (1989) demonstrated that [Mg$^{2+}$], increases from 0.85 to 2.1 mM during 10 to 15 min of ischemia. Our model suggested significant effects of [Mg$^{2+}$], rise from 0.85 to 2.1 mM on gap junction function. Namely, P$_{open}$ of gap junction decreased to 1.5 % of control (0.097 to 0.0015) and conduction velocity to 6.9 % of control (281.7 to 19.5 mm/s). Furthermore, two-to-one block was induced at 2.6 mM [Mg$^{2+}$]. Therefore, the increase in [Mg$^{2+}$], observed in cardiac ischemia may significantly affect ventricular conductance. Haigney et al. (1998) showed that [Mg$^{2+}$], decreased from 1.06 to 0.49
mM in the dog failing heart caused by rapid pacing. A moderate effect of the fall in [Mg$^{2+}$]$_i$ on gap junction was suggested in our simulation. The change in [Mg$^{2+}$]$_i$ increases the open probability by 7.5 times (0.053 to 0.39) and conduction velocity by 2.9 times (224.7 to 645 mm/s). In real heart, however, H$^+$ and Ca$^{2+}$ may also modulate gap junction function (Bukauskas et al. 2004; Peracchia et al. 2004; Desplantez et al. 2007). Closing of gap junction under the ischemic condition would be further accelerated by the rise of H$^+$ and Ca$^{2+}$.

Recently, complicated patterns of reentrant arrhythmia have been demonstrated in computer simulations, where numerical models of myocyte are connected by gap junctions as resistance factor. In most simulation studies, however, the gap junction conductance was regarded as constant. We previously constructed a mathematical model of the Ca$^{2+}$ gating of GJCs (Oka et al. 2006). Conduction velocity was clearly dependent on cytoplasmic Ca$^{2+}$ concentration and remarkable decrease in conduction velocity was induced under a Ca$^{2+}$-overloaded condition. Thus, chemical gating by H$^+$ and Ca$^{2+}$ and Mg$^{2+}$ has to be considered in the model simulation to get deeper insight into the mechanisms underlying arrhythmia under pathological conditions.

### 4.4. Clinical use of Mg$^{2+}$

Intravenous administration of Mg$^{2+}$ has been reported to be useful in preventing some kinds of arrhythmias such as rapid atrial fibrillation and torsades de pointes (TdP) (Onalan et al. 2007; Hoshino et al. 2006). The main antiarrhythmic mechanism of Mg$^{2+}$ has been considered to be a Ca$^{2+}$ antagonistic effect (Gomez 1998). The effective antiarrhythmic concentration of free Mg$^{2+}$ is 0.5 ~ 1.0 mM: 1.1 ~ 1.2 mM (Hoshino et al. 2006) and 0.54 mM (Kasaoka et al. 1996)). On the other hand, Mg$^{2+}$ at high concentration is known to cause several side effects such as bradycardia or atrioventricular block. Symptoms of Mg$^{2+}$ toxicity were reported to develop at its serum concentration more than 1.7 ~ 2.6 mM (Musso 2009). Considering that free Mg$^{2+}$ is about half of total serum Mg$^{2+}$ concentration and that intracellular free Mg$^{2+}$ is about the same as extracellular free Mg$^{2+}$ concentration, high dose application of Mg$^{2+}$ may cause blockade of GJCs in clinical situation.
5 Conclusions

We studied kinetics of Mg\(^{2+}\) gate of gap junction in intact guinea pig ventricular cells and reconstructed mathematical models. It was suggested that the rise of [Mg\(^{2+}\)]\(_i\) remarkably modulates conductance of gap junction.

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FIGURES’ LEGEND

Fig.1. Electrophysiological recording of gap junction channel conductance
A. Diagram of the experiment setup. B. Pulse protocol. The current response to ± 5 mV from a holding potential of 0 mV was recorded every 500 ms. C. Equivalent circuits before and after laser perforation. ELv: electrode for voltage recording, ELi: electrode for current feeding, Gj: gap junction channels conductance, Gm,1: non-junctional membrane conductance of cell 1, Gm,2: non-junctional membrane conductance of cell 2, Gs: conductance through the holes made in cell 2.

Fig.2. Time course of Gj change
A. Representative current records. B. Time course of Gj change. a: before perforation, b: immediately after perforation, c: steady state at 10 mM d: Mg2+, steady state at 0.1 mM Mg2+. Mg2+ concentrations are denoted above the graph.
Fig. 3. Comparison of time course of $G_j$ decay at three different distances between the junction and holes. The myocytes were superfused with a solution containing 10 mM Mg$^{2+}$. $G_j$ was normalized to the one immediately after the laser perforation.

Fig. 4. Kinetics of Mg$^{2+}$ gate of GJC.
Single exponential fit of $G_j$ modulation Mg$^{2+}$. Data (black circles, A: 10 mM Mg$^{2+}$, B: 0.6 mM Mg$^{2+}$) are fitted with single exponential function (gray circles). Relationships between $[\text{Mg}^{2+}]_i$ and $\tau^{-1}$ (D) and between $[\text{Mg}^{2+}]_i$ and $G_j$ (C). In C, the data were normalized to the value at 0.1 mM $[\text{Mg}^{2+}]_i$. The data numbers were three at 3 and 10 mM $[\text{Mg}^{2+}]_i$ in (C), and two at 3 mM $[\text{Mg}^{2+}]_i$ and four at 10 mM $[\text{Mg}^{2+}]_i$, respectively. Others are one. Gray lines are least square fit. In B, the equation is $y = y_0 + a\cdot c^b/(c^b+[\text{Mg}^{2+}]_b)$, with $a=0.957$, $b=3.4$, $c=0.6$, $y_0=0.0217$. $G_j$ was normalized to the one at 0.1 mM Mg$^{2+}$.

Fig. 5. Schemes of Mg gate model. A. Two-state model. B. Four-state model.

Fig. 6. Simulation results.
A. A relationship between $[\text{Mg}^{2+}]_i$ and $G_j$. B. A relationship between $[\text{Mg}^{2+}]_i$ and $\tau^{-1}$. Black circles are experimental data (same as in fig. 4C) and gray circles are model simulation. C. and D. Time course. Experimental data (black circles: same as fig. 4D) were Gj normalized to the one at 0.1mM Mg$^{2+}$ and open probability was plotted for model simulation (gray circles).

Fig. 7. Action potential propagation in the 1-D cable model at various $[\text{Mg}^{2+}]_i$.
A. Diagram of 1-D cable model. B. Relationship between $[\text{Mg}^{2+}]_i$ and conduction velocity (CV) in the 1-D cable model. CV was calculated at cell 15. C. Action potentials of 30 cells at three different $[\text{Mg}^{2+}]_i$: 0.1 (left), 0.7 (middle), 2.0 (right) mM.

Fig. 8. Time course of gap junction modulation by Mg$^{2+}$. A. Open probability. $[\text{Mg}^{2+}]_i$ of cell 11-30 was increased to 2.6 mM at time of 4.0 sec. B. Action potentials of the 30 cells 76 sec after the $[\text{Mg}^{2+}]_i$ increase.
Fig. 1

A

Paired cells

cell 1  cell 2

chamber

Laser irradiation (10 Hz)

B

0 mV
+5 mV
-5 mV
0 mV

100 ms
50 ms
150 ms

Before the perforation  After the perforation

C

ELv  ELl

Gj

Gm,1  Gm,2

Before the perforation  After the perforation

Gm,1  Gs
Fig. 2

A

B

10 mM Mg^{2+}  0.1 mM Mg^{2+}

Laser irradiation

G_i (nS)

Time (s)

0  50  100  150  200  250  300

0  500  1000  1500  2000  2500  3000

1 nA

10 nA

100 ms
Fig. 3

Paired cells

L (μm)

Laser irradiation

Normalized \( G_j \)

Time (s)

\( L = \{ \quad \}

- ○ 48
- ● 18
- □ 11
Fig. 4

A. 10 mM Mg\(^{2+}\)

B. 0.6 mM Mg\(^{2+}\)

C. Normalized Gj vs. [Mg\(^{2+}\)]i (mM)

D. 1/\(\tau\) vs. [Mg\(^{2+}\)]i (mM)
Fig. 5

A: simple two state model

\[ 3\text{Mg}^{2+} \]

\[ \xrightarrow[\alpha]{\text{O}} \text{C} \xleftarrow[\beta]{\text{C}} \]

B-1: RLS in 1\textsuperscript{st} step

\[ \text{Mg}^{2+} \]

\[ \xrightarrow[\alpha_1]{\text{O}} \text{C} \xleftarrow[\beta_1]{\text{C}} \text{C} \xrightarrow[\alpha_2]{\text{C}} \text{C} \xleftarrow[\beta_2]{\text{C}} \]

B-2: RLS in 2\textsuperscript{nd} step

\[ \text{Mg}^{2+} \]

\[ \xrightarrow[\alpha_1]{\text{O}} \text{O} \xleftarrow[\beta_1]{\text{O}} \text{C} \xrightarrow[\alpha_2]{\text{C}} \text{C} \xleftarrow[\beta_2]{\text{C}} \text{C} \xrightarrow[\alpha_3]{\text{C}} \]

B-3: RLS in 3\textsuperscript{rd} step

\[ \text{Mg}^{2+} \]

\[ \xrightarrow[\alpha_1]{\text{O}} \text{O} \xleftarrow[\beta_1]{\text{O}} \text{O} \xrightarrow[\alpha_2]{\text{O}} \text{O} \xleftarrow[\beta_2]{\text{O}} \text{O} \xrightarrow[\alpha_3]{\beta_3} \text{C} \]
Fig. 6

A) Open probability as a function of $[\text{Mg}^{2+}]_i$ (mM) plotted against time (s).

B) $\tau$ inverse as a function of $[\text{Mg}^{2+}]_i$ (mM).

C) Normalized $G_i$ as a function of time (s).

D) Normalized $G_i$ as a function of time (s).