All deformation potentials in GaN determined by reflectance spectroscopy under uniaxial stress: Definite breakdown of the quasicubic approximation

Ryota Ishii,* Akio Kaneta, Mitsuru Funato, and Yoichi Kawakami†
Department of Electronic Science and Engineering, Kyoto University, Kyoto 615-8510, Japan

Atsushi A. Yamaguchi
Research Laboratory for Integrated Technological Systems, Kanazawa Institute of Technology, Tokyo 105-0002, Japan

(Received 7 December 2009; revised manuscript received 1 March 2010; published 8 April 2010)

Reflectance spectroscopy was performed for nonpolar and semipolar bulk GaN substrates under uniaxial stress. The exciton-transition energies and oscillator strengths clearly depended on uniaxial stress. The results were interpreted by group theory on excitons and analyzed in terms of the effective Hamiltonian proposed by Bir and Pikus. In the Hamiltonian, the short-range Coulomb interaction was taken into consideration. This approach allows all exciton deformation potentials to be determined without the quasicubic approximation and we found that the exciton deformation potentials in GaN considerably deviate from the quasicubic approximation.

DOI: 10.1103/PhysRevB.81.155202 PACS number(s): 71.70.Fk, 71.55.Eq, 71.35.—y

I. INTRODUCTION

GaN-related materials have received much attention for optoelectronic applications because they interact with photons in the ultraviolet to near-infrared region. Highly efficient light-emitting diodes and high-power laser diodes (LDs) have been realized in the blue and blue-green regions using InGaN quantum wells (QWs).

However, the external quantum efficiency of c-plane (polar plane) InGaN QWs decreases rapidly as the In composition increases due to the quantum-confined Stark effect caused by the piezoelectric polarization and degradation of the crystal quality. Under these circumstances, a reduction in the piezoelectric polarization in nonpolar and semipolar planes has been theoretically predicted, and is consistent with the experimental observation of fast radiative recombination lifetimes compared to those for polar plane. Recently, many groups including ours have reported crystal growth and device fabrications onto these nonpolar and semipolar QWs. Moreover, green LDs operating at 520 nm at room temperature in the continuous-wave mode have been realized on the (201) plane. To fully understand the device performances and to create novel functionalities, it is indispensable to understand the electronic band structures of nonpolar and semipolar QWs.

It has been theoretically predicted that anisotropic strain will drastically alter the electronic band structures of strained nonpolar and semipolar InGaN or AlGaN QWs. Additionally, a very interesting physical phenomenon, polarization switching was theoretically predicted. Then, our group has experimentally observed polarization switching in (1122) plane InGaN QWs with high In compositions and we noted the importance of strain. However, the analyses in the aforementioned papers relied on the accuracy of GaN physical parameters such as the deformation potentials. Therefore, to fully understand fundamental physics and to optimize device performance, a comprehensive study about the strain-induced effects in GaN-related materials is necessary.

It has been demonstrated that deformation potential theory can well explain the influence of strain on the electronic band structure. In this theory, strain-induced effects are completely determined by the exciton deformation potentials. To date, nearly all papers have experimentally determined exciton deformation potentials in GaN using a combination of x-ray measurements with optical measurements on heteroepitaxial c-plane GaN films. However, errors are unavoidable in these estimated values due to strain inhomogeneity. Furthermore, the quasicubic approximation must be used to extract all the exciton deformation potentials because this approach gives only a part of the exciton deformation potentials or the relations between them. On the other hand, optical measurements under uniaxial stress are very suitable measurement techniques which can create various strained environments. Such techniques have yielded the exciton deformation potentials of II–VI compounds such as ZnO, CdS, and CdSe without the quasicubic approximation. For c-plane GaN thin films grown on sapphire, Yamaguchi et al. have demonstrated reflectance spectroscopy under uniaxial stress. However, the relatively large linewidth due to the heteroepitaxial growth prevented the C-exciton transition (Hopfield notation) from being resolved, and furthermore, all the exciton deformation potentials have not been experimentally obtained because the uniaxial stress was applied only perpendicular to c-axis.

Herein, reflectance spectroscopy is performed for nonpolar and semipolar bulk GaN substrates under uniaxial stress along various directions. The experimental results are interpreted by group theory on excitons and analyzed in terms of the effective Hamiltonian where the short-range Coulomb interaction was taken into consideration. This approach allows us to determine all the exciton deformation potentials without the quasicubic approximation and we found that the exciton deformation potentials in GaN considerably deviate from the quasicubic approximation. Finally, we proposed a set of material parameters such as the crystal-field interaction, the spin-orbit interaction, the spin-exchange interaction, and the exciton deformation potentials in GaN.
TABLE I. Valence-electron and conduction-electron matrix at the Γ point. \([F_z = \Lambda_1 + \Delta_2 + C_1 \epsilon_{xx} + C_2 (\epsilon_{yy} + \epsilon_{yz}) + C_4 (\epsilon_{yy} + \epsilon_{yz}) \pm j; G_z = \Lambda_1 - \Delta_2 + C_1 \epsilon_{xx} + C_2 (\epsilon_{yy} + \epsilon_{yz}) + C_4 (\epsilon_{yy} + \epsilon_{yz}) \pm j/2; j = C_2 (\epsilon_{xx} + \epsilon_{yy} + \epsilon_{yz}) \pm j/2; K = -C_2 (\epsilon_{xx} - \epsilon_{yy} - 2i\epsilon_{yz}); and H = -C_2 (\epsilon_{xx} - i\epsilon_{yz})\].

<table>
<thead>
<tr>
<th>(p_{,\alpha\alpha})</th>
<th>(p_{,\alpha\alpha})</th>
<th>(p_{,\beta\alpha})</th>
<th>(p_{,\alpha\beta})</th>
<th>(p_{,\beta\beta})</th>
<th>(p_{,\alpha\beta})</th>
<th>(p_{,\alpha\beta})</th>
<th>(p_{,\beta\alpha})</th>
<th>(p_{,\beta\alpha})</th>
<th>(p_{,\alpha\alpha})</th>
</tr>
</thead>
<tbody>
<tr>
<td>(F_z)</td>
<td>(K)</td>
<td>0</td>
<td>(-j)</td>
<td>0</td>
<td>0</td>
<td>0</td>
<td>0</td>
<td>0</td>
<td>0</td>
</tr>
<tr>
<td>(K^*)</td>
<td>(G_-)</td>
<td>(\sqrt{2}\Delta_3)</td>
<td>(\lambda_+)</td>
<td>0</td>
<td>0</td>
<td>0</td>
<td>(H^*)</td>
<td>0</td>
<td>0</td>
</tr>
<tr>
<td>(\lambda_+)</td>
<td>0</td>
<td>(G_-)</td>
<td>(K)</td>
<td>(\sqrt{2}\Delta_3)</td>
<td>0</td>
<td>0</td>
<td>0</td>
<td>0</td>
<td>0</td>
</tr>
<tr>
<td>(H^*)</td>
<td>0</td>
<td>0</td>
<td>(K^*)</td>
<td>(F_z)</td>
<td>0</td>
<td>0</td>
<td>0</td>
<td>0</td>
<td>0</td>
</tr>
<tr>
<td>(\lambda_+)</td>
<td>0</td>
<td>0</td>
<td>(H)</td>
<td>(F_z)</td>
<td>0</td>
<td>0</td>
<td>0</td>
<td>0</td>
<td>0</td>
</tr>
<tr>
<td>(H^*)</td>
<td>0</td>
<td>0</td>
<td>(-H^*)</td>
<td>(K^*)</td>
<td>(G_+)</td>
<td>(\sqrt{2}\Delta_3)</td>
<td>0</td>
<td>0</td>
<td>0</td>
</tr>
<tr>
<td>(\lambda_+)</td>
<td>0</td>
<td>0</td>
<td>(H)</td>
<td>(F_z)</td>
<td>0</td>
<td>0</td>
<td>0</td>
<td>0</td>
<td>0</td>
</tr>
<tr>
<td>(\lambda_+)</td>
<td>0</td>
<td>0</td>
<td>(-H^*)</td>
<td>(K^*)</td>
<td>(F_z)</td>
<td>0</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>(H^*)</td>
<td>(-H)</td>
<td>0</td>
<td>0</td>
<td>0</td>
<td>0</td>
<td>0</td>
<td>0</td>
<td>0</td>
<td>0</td>
</tr>
<tr>
<td>(\lambda_+)</td>
<td>0</td>
<td>0</td>
<td>(-j)</td>
<td>(K^*)</td>
<td>(G_+)</td>
<td>(\sqrt{2}\Delta_3)</td>
<td>0</td>
<td>0</td>
<td>0</td>
</tr>
<tr>
<td>(\lambda_+)</td>
<td>0</td>
<td>0</td>
<td>(H)</td>
<td>(F_z)</td>
<td>0</td>
<td>0</td>
<td>0</td>
<td>0</td>
<td>0</td>
</tr>
<tr>
<td>(\lambda_+)</td>
<td>0</td>
<td>0</td>
<td>(-H^*)</td>
<td>(K^*)</td>
<td>(F_z)</td>
<td>0</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>(H^*)</td>
<td>(-H)</td>
<td>0</td>
<td>0</td>
<td>0</td>
<td>0</td>
<td>0</td>
<td>0</td>
<td>0</td>
<td>0</td>
</tr>
</tbody>
</table>

II. THEORY

In this section, the effective exciton Hamiltonian in hexagonal symmetry is described using notations in Ref. 32. We expanded this theory to obtain oscillator strengths for arbitrary directions. This procedure is essential for describing the optical properties of wurtzite materials grown on semipolar planes.

According to the effective Hamiltonian formalism proposed by Bir and Pikus, the exciton Hamiltonian can be written as

\[
H = H_0^b + H_0^c + H^{strain} + H^{exciton},
\]

(1)

where \(H_0^b\) and \(H_0^c\) are the nonstrain conduction and valence-band Hamiltonian, respectively. \(H^{strain}\) is a strain-related term. \(H^{exciton}\) originates from essentially two particle natures and is written as

\[
H^{exciton} = B.E. + \frac{1}{2} j (\sigma_x \cdot \sigma_y),
\]

(2)

where B.E. describes the exciton binding energy and \(j\) is the spin-exchange interaction constant which describes the electron-hole exchange interaction.\(^{33}\)

Considering a hexagonal symmetry, Eq. (1) can be rewritten by the invariant method as

\[
H^{wurtzite} = \Delta_1 J_z \epsilon_{zz} + \Delta_2 J_z \sigma_{zz} + \Delta_3 (\sigma_{xx} + \sigma_{yy} - \sigma_{zz} \epsilon_{zz})
\]

\[
+ (C_1 + C_2 J_z^2) \epsilon_{zz} + (C_2 + C_4 J_z^2) (\epsilon_{xx} + \epsilon_{yy}) - C_5 (J_z^2 \epsilon_{xx} + J_z^2 \epsilon_{yy}) - 2C_6 [(J_z \epsilon_{xx}) \epsilon_{zz} + (J_z \epsilon_{yy}) \epsilon_{zz})
\]

\[
+ B.E. + \frac{1}{2} j (\sigma_x \cdot \sigma_y),
\]

(3)

where

\[
\epsilon_{zz} = \epsilon_{xx} - \epsilon_{yy} \pm 2i \epsilon_{xy}, \quad \epsilon_{zz} = \epsilon_{zz} \pm i \epsilon_{xz}
\]

(4)

and

\[
C_1 = D_1 - a_{cz}, \quad C_2 = D_2 - a_{cz}, \quad C_3 = D_3, \quad C_4 = D_4,
\]

Here \(J\) and \(\sigma\) describe the orbital angular momentum and Pauli-matrix operator, respectively, and their definitions are the same as those in Ref. 28. \(\Delta_1\) describes the crystal-field interaction due to the uniaxial crystal while \(\Delta_2\) and \(\Delta_3\) describe the spin-orbit interaction parallel and perpendicular to the \(c\) axis, respectively. \(C_1\) describes the exciton deformation potentials, which are related to the conduction-band deformation potentials \((a_{cz}, a_{cy})\) and the valence-band deformation potentials \(D_j\) by Eq. (5). \(\epsilon_{ij}\) represents strain tensors. It should be noted that herein a negative sign indicates compressive strain, and we assumed that B.E. and \(\Delta_{1,3}\) are strain independent.\(^{34}\)
ALL DEFORMATION POTENTIALS IN GaN DETERMINED...
To confirm the operation of the uniaxial stress-loading system, we performed polarized reflectance measurements on c-plane GaN films grown on sapphire substrates. Uniaxial stress was applied perpendicular to the a-plane, which is the same condition reported in Ref. 30. It was confirmed that the uniaxial stress-loading system worked properly because our data and that in Ref. 30 quantitatively agree (data not shown).

Before describing the experimental results, the role of the exciton deformation potentials ($C_1 - C_6$) in determining the electronic structures is explained. If the crystal symmetry belongs to the $C_{6v}$ point group, $\epsilon_{xx}$ and $\epsilon_{yy}$ are equivalent and all the nondiagonal terms of the strain tensors are zero. Therefore, the uniaxial stress dependence of the energy difference between exciton transitions (relative exciton-transition energies) are determined by only the exciton deformation potentials of $C_3$ and $C_4$. When the crystal symmetry is reduced from the $C_{6v}$ to the $C_{2v}$ point group, the exciton deformation potential of $C_4$ as well as $C_1$ and $C_4$ are necessary to describe the uniaxial stress dependence of the relative exciton-transition energies because the strain tensors of $\epsilon_{xx}$ and $\epsilon_{yy}$ are no longer equivalent. In addition, applying stress along an off-crystal axis direction further reduces the crystal symmetry, when the exciton deformation potential of $C_6$ is also necessary to describe the uniaxial stress dependence of the relative exciton-transition energies. The exciton deformation potentials of $C_1$ and $C_2$ are only required to explain the uniaxial stress dependence of the absolute exciton-transition energies.

Initially, polarized reflectance measurements were performed on the a-plane bulk GaN substrate under uniaxial stress perpendicular to the c-plane ($P||c$). Figures 4(a) and 5 show the configuration and results, respectively. The uniaxial stress dependence of some reflection anomalies with sharp linewidths were observed for both $E||c$ and $E\perp c$. Using group theory on excitons as discussed above, the resonances observed for $E||c$ were assigned to $\Gamma_{1}^{B}(n=1)$, $\Gamma_{1}^{C}(n=1)$, and $\Gamma_{1}^{C}(n=2)$ on the order of the photon energy. As the uniaxial stress increased, a new anomaly appeared nearby $\Gamma_{1}^{C}(n=1)$, which we assigned to $\Gamma_{1}^{B}(n=2)$ considering the stress dependence and exciton binding energy. Likewise, the resonances observed for $E\perp c$ were assigned to $\Gamma_{1}^{A}(n=1)$, $\Gamma_{2}^{A}(n=1)$, $\Gamma_{3}^{A}(n=2)$, and $\Gamma_{1}^{C}(n=1)$. It should be noted that a weak feature was also observed on the lower energy side of $\Gamma_{3}^{C}(n$
FIG. 6. (Color online) Uniaxial stress dependence of exciton-transition energies for $P\parallel c$. (a) Relative to $\Gamma^2_3(n=1)$ and (b) absolute values. Symbols represent the experimental data whereas the lines are the fitted results.

$\Gamma^2_3(n=1)$ and we assigned it to an impurity-related exciton of $\Gamma^2_3(n=1)$ considering the uniaxial stress dependence. Since we mainly focus on the intrinsic optical properties of GaN in this paper, these features are neglected in the following discussions.

Here, we assumed that the exciton-transition energies correspond to the average energies between those for the reflection maxima and the minima. The exciton-transition energies obtained in this analysis agreed well with those extracted by the fitting based on the exciton-transition model. It should be noted that reflectance spectra should be interpreted by the exciton-polariton model, which includes spatial dispersion. However, it has been reported that the fitting results using the exciton-transition model and the exciton-polariton model are within the experimental error for GaN, which is why we neglected the exciton-polariton effects in this study. Figure 6 shows the uniaxial stress dependence of the relative energies with respect to $\Gamma^2_3(n=1)$ and the absolute exciton-transition energies.

Secondary, polarized reflectance measurements were performed on the $a$-plane bulk GaN substrate under uniaxial stress perpendicular to the $m$-plane ($P \perp c$). Figures 4(b) and 7 show the configuration and results, respectively. Using group theory on excitons, the resonances observed for $E\parallel c$ were assigned to $\Gamma^2_3(n=1)$, $\Gamma^2_3(n=1)$, and $\Gamma^2_3(n=2)$ on the order of photon energy. Under high stress, a new feature, which was assigned to $\Gamma^2_3(n=1)$ by considering group theory, appeared around 3.480 eV. Similarly, the resonances observed for $E \perp c$ were assigned to $\Gamma^2_3(n=1)$, $\Gamma^3_2(n=1)$, $\Gamma^3_2(n=2)$, and $\Gamma^3_2(n=1)$. Figure 8 shows the uniaxial stress dependence of the relative energies with respect to $\Gamma^2_3(n=1)$ and the absolute exciton-transition energies.

Finally, we performed polarized reflectance measurements on the (1122) plane bulk GaN substrate under uniaxial stress $P\parallel[\bar{1}2\bar{1}3]$. Figures 4(c) and 9 show the configuration and

FIG. 7. Reflectance spectra of the $a$-plane GaN bulk substrate under uniaxial stress $P \perp c$. (a) $E\parallel c$ and (b) $E \perp c$.

FIG. 8. (Color online) Uniaxial stress dependence of exciton-transition energies for $P \perp c$. (a) Relative to $\Gamma^2_3(n=1)$ and (b) absolute values. Symbols represent experimental data whereas the lines are the fitted results.

FIG. 9. Reflectance spectra of the (1122) plane GaN bulk substrate under uniaxial stress $P\parallel[\bar{1}2\bar{1}3]$. (a) $E\parallel[\bar{1}2\bar{1}3]$ and (b) $E\parallel[1100]$.
results, respectively. So far, $C_6$ has not been experimentally determined but it has been obtained by the quasicubic approximation. In contrast, the current configuration enables the exciton deformation potential of $C_6$ to be experimentally determined because shear strain is induced. For $E \parallel [\overline{1}1\overline{2}]$, three exciton transitions were resolved as shown in Fig. 9(a) although six exciton transitions are theoretically dipole allowed. This is because three of six exciton transitions with relatively large oscillator strengths were experimentally observed for $E \parallel [\overline{1}1\overline{2}]$. This situation may cause misassignment on excitons. On the other hand, in the case of $E \parallel [\overline{1}0\overline{0}]$, we experimentally observed three exciton transitions as shown in Fig. 9(b) and found that three exciton transitions are theoretically dipole allowed. Therefore, we assigned the exciton-transition energies only for this configuration. Figure 10 shows the uniaxial stress dependence of the absolute exciton-transition energies.

IV. DISCUSSION

Initially, the experimental results for samples under no stress were analyzed in order to obtain the “bulk parameters”: $\Delta_1$, $\Delta_2$, $\Delta_3$, and $j$. In this particular case, the exciton-transition energies were deduced by the exciton-model fitting. The bulk parameters were completely determined from the relative exciton-transition energies. We assumed that our samples were strain free. Table II lists the bulk parameters and the relative exciton-transition energies. $\Delta_2$ and $\Delta_3$ agree well with the reported values but $\Delta_1$ appears somewhat larger. Reference 23 has noted that the residual strain in HVPE samples may cause an overestimation of $\Delta_1$. However, if the short-range Coulomb interaction is not considered, deduction from the relative exciton-transition energies of $\Gamma_{\overline{3}}^A - \Gamma_{\overline{3}}^A$ and $\Gamma_{\overline{3}}^C - \Gamma_{\overline{3}}^C$ indicate that $\Delta_1 = 9.8$ meV, $\Delta_2 = 5.9$ meV, and $\Delta_3 = 6.0$ meV. These values are in excellent agreement with those of a homoepitaxial GaN.37 Therefore, we believe that $\Delta_1$ is larger due to the short-range Coulomb interaction rather than the residual biaxial strain. Hence, the short-range Coulomb interaction is important in determining the bulk parameters. Next, we examined the value of $j$. To date, $j$ has been experimentally extracted by the relative exciton-transition energy between $\Gamma_{\overline{3}}^A$ and $\Gamma_{\overline{3}}^B$. However, as shown in Fig. 11, it should be determined by the relative exciton-transition energy between $\Gamma_{\overline{3}}^C$ and $\Gamma_{\overline{3}}^A$ rather than between $\Gamma_{\overline{3}}^B$ and $\Gamma_{\overline{3}}^C$. Hence, we simultaneously performed the least-squares fitting for all observables and determined the bulk parameters in GaN. Table II shows the results.

Next, the uniaxial stress dependence of the exciton-transition energies was analyzed using analytical expressions of the valence-electron and conduction-electron matrix derived by Gil et al.22,41 due to their usefulness in the least-squares fitting. When a uniaxial stress is applied perpendicular to the c-plane ($P\parallel c$), the relative exciton-transition energy

\begin{table}[h]
\centering
\begin{tabular}{|c|c|c|c|c|c|c|c|c|}
\hline
Experimental & $\Gamma_{\overline{3}}^B - \Gamma_{\overline{3}}^A$ & $\Gamma_{\overline{3}}^C - \Gamma_{\overline{3}}^C$ & $\Gamma_{\overline{3}}^B - \Gamma_{\overline{3}}^A$ & $\Gamma_{\overline{3}}^C - \Gamma_{\overline{3}}^C$ & $\Delta_1$ & $\Delta_2$ & $\Delta_3$ & $j$ \\
\hline
Korona et al. & 5.1 & 23.9 & 5.1 & 23.9 & 9.3 & 6.6 & 6.6 & \\
Yamaguchi et al. & 4.5 & 22.0 & 4.5 & 22.0 & 10.0 & 5.5 & 6.0 & \\
Shikanai et al. & 7.4 & 29.6 & 7.4 & 29.6 & 22.0 & 5.0 & 5.0 & \\
Gil et al. & 4.3 & 21.2 & 3.7 & 23.1 & 10.0 & 5.1 & 6.1 & 2.0 \\
Gil et al. & 4.3 & 21.1 & 4.4 & 21.5 & 8.7 & 5.7 & 5.9 & 0.6 \\
Paskov et al. & 6.1 & 22.0 & 6.2 & 22.5 & 10.0 & 6.2 & 5.5 & 0.6 \\
Paskov et al. & 4.4 & 21.7 & 4.4 & 22.2 & 10.0 & 5.5 & 6.0 & 0.6 \\
Present work & 4.9 & 22.5 & 4.6 & 23.8 & 12.3 & 5.2 & 5.9 & 1.2 \\
\hline
\end{tabular}
\end{table}

$^a$Reference. 37.
$^b$Reference. 23.
$^c$Reference. 21.
$^d$Reference. 41.
$^e$Reference. 42.
$^f$Reference. 43.
All deformation potentials in GaN determined perpendicular to the c-axis and obtained. We performed the least-squares fitting for all the observables as follows:

\[ \alpha'' = C_3 S_{33} + 2 C_4 S_{13}. \]  

(10)

We performed the least-squares fitting for all the observables and obtained \( \alpha'' = 0.0174 \) eV/GPa.

Likewise, when a uniaxial stress is introduced perpendicular to the m-plane \( (P \perp c) \), the relative exciton-transition energies can be determined by the following two parameters:

\[ \beta'' = C_3 S_{13} + C_4 (S_{11} + S_{12}), \]

\[ \beta'' = -C_5 (S_{11} - S_{12}). \]  

(11)

Using the same procedure, we obtained \( \beta'' = -0.0136 \) eV/GPa and \( \beta'' = 0.0116 \) eV/GPa.

Also, for the \( (11 \bar{2} 2) \) plane with \( P \parallel [\bar{T} \overline{2} 3] \), the following three parameters are necessary:

\[ \gamma'' = C_3 (S_{13} \cos^2 \theta + S_{33} \sin^2 \theta) \]

\[ + C_4 [(S_{11} + S_{12}) \cos^2 \theta + 2 S_{13} \sin^2 \theta], \]

\[ \gamma'' = -C_3 (S_{11} - S_{12}) \cos^2 \theta, \]

\[ \gamma'' = \frac{1}{2} C_4 S_{44} \sin \theta \cos \theta. \]  

(12)

The least-squares fitting gave \( \gamma'' = 0.0086 \) eV/GPa, \( \gamma'' = 0.0026 \) eV/GPa, and \( \gamma'' = 0.0081 \) eV/GPa. At this stage, \( \alpha, \beta, \) and \( \gamma \) were determined separately from their respective experiments.

Then, to obtain the exciton deformation potentials from Eqs. (10)–(12), the values of the elastic constants must be determined. Although numerous sets of the elastic constants have been proposed using both first-principles calculations and experiments such as brillouin scattering or resonance ultrasound method,44 the values are largely scattered. Herein, we chose the elastic constants proposed by Yamaguchi et al.45 because this study determined the elastic constants using bulk GaN substrates. It should be noted that the elastic constants in GaN were assumed to be independent of temperature.46

With the selected set of elastic constants, we obtained three equations for the relationship between \( C_3 \) and \( C_4 \) \( (\alpha'', \beta'', \) and \( \gamma'' \) \), two equations for \( C_3 \) \( (\beta'' \) and \( \gamma'' \) \), and one equation for \( C_5 \theta' \). To extract the most plausible four parameters \( (C_3 \) to \( C_6) \) from these 12 equations, \( C_3 \) and \( C_4 \) were initially determined because these were experimental observables in all the experiments. Figure 12 plots the three relations between \( C_1 \) and \( C_4 \) using the elastic constants in Ref. 45. These three relations form a triangle and selecting the position close to the center simultaneously minimizes the errors in all the experiments. Thus, the values for \( C_3 \) and \( C_4 \) were 4.9 and \( -5.0 \) eV, respectively.

Next, \( \beta'' \), \( \gamma'' \), and \( \theta' \) were re-evaluated using the least-squares fitting where \( \beta'' \) and \( \gamma'' \) were constants using the above values of \( C_3 \) and \( C_4 \). Consequently, \( C_5 = -2.8 \) eV and \( C_6 = -3.1 \) eV were obtained. To determine the sign of \( C_5 \) and \( C_6 \), the uniaxial stress dependence of the oscillator strengths as well as the exciton-transition energies were considered.

Finally, \( C_1 \) and \( C_2 \) were determined from the absolute exciton-transition energies. Similar to \( C_3 \), and \( C_4 \), \( C_1 \) and \( C_2 \) were experimental observables in all the experiments. Therefore, the same procedure was used to evaluate \( C_1 \) and \( C_2 \) as 6.5 and 11.2 eV, respectively.

Using the exciton deformation potentials obtained by the current analysis, theoretical calculations were performed and the results are shown in Figs. 6, 8, and 10. The accuracies of the exciton deformation potentials of \( C_{i=1,2,3,4,5} \) and \( C_6 \) were about 10% and 20%, respectively. It must be emphasized that the experimental uniaxial stress dependence of the exciton oscillator strengths can also be qualitatively described by theoretical calculations using our values. Table III lists the exciton deformation potentials. If the elastic constants proposed by Polian et al.56 were used then the exciton-deformation potentials were \( C_1 = 7.6 \) eV, \( C_2 = 11.8 \) eV, \( C_3 \)
It is found that our exciton deformation potentials significantly differ from previous studies, especially the relation between $C_3$ and $C_4$. The well-known approximation in the hexagonal symmetry, that is, the quasicubic approximation, can be written as

$$D_1 - D_2 = -2D_4, \quad D_3 + 4D_5 = \sqrt{2}D_6. \quad (13)$$

Using our notation, Eq. (13) can be written as

$$C_1 - C_2 - (a_{cz} - a_{ct}) = -C_3 = 2C_4,$$

To explain the quasicubic approximation in hexagonal wurtzite structures, the ideal $sp^3$ tetrahedral bond configuration in cubic zincblende structures must be recalled. In “ideal” hexagonal wurtzite structures, the ideal $sp^3$ tetrahedron is preserved, and $c/a$ must equal 1.63299, when the quasicubic approximation is strictly applicable. However, if the $c/a$ ratio deviates from the ideal value, the quasicubic approximation need not to be rigorously satisfied.

To date, the quasicubic approximation has been considered a good approximation in GaN which has a $c/a$ ratio of

$$C_3 + 4C_5 = \sqrt{2}C_6. \quad (14)$$
1.6259. Gil and co-workers\cite{58} have suggested that this approximation is experimentally supported because significant differences between the hydrostatic stress dependence of heavy hole and other bands in the absorption spectra are not present. In fact, $C_1 = -2C_4$ must be satisfied as long as the hydrostatic strain does not vary the energy difference between the heavy-hole band and the other bands. However, the absorption spectra\cite{59,60} that Gil et al. relied on do not exhibit clear exciton-related transitions and the relation between $C_3$ and $C_4$ is generally different between under hydrostatic stress and hydrostatic strain. Thus, their assertion does not seem firm.

Table III also lists the deformation potentials in II-VI compounds. It should be noted that the valence-band ordering of ZnO is still controversial.\cite{61,62} Langer et al.\cite{28} have reported that the irreducible representation of the topmost valence band is $\Gamma^0_\alpha$, while Wrzesinski and Fröhlich\cite{29} have assigned that the valence-band ordering is $\Gamma^7_\alpha-\Gamma^8_\beta-\Gamma^4_\alpha$. The different assignment of the valence bands may cause the differences between the deformation potentials in ZnO in Refs. 28 and 29. Because it cannot be concluded which arrangement is correct at this stage, both the deformation potentials are listed in Table III. Herein, we evaluated the deviation from the quasicubic approximation by comparing $C_3/(-2C_4)$. The deformation potentials in ZnO and CdS do not satisfy the quasicubic approximation most likely because they have “nonideal” wurtzite structures. Figure 13 plots the $c/a$ ratio for select wurtzite semiconductors. It is found that the deviation from the quasicubic approximation is correlated with the $c/a$ ratio. Therefore, we concluded that our exciton deformation potentials in GaN show the breakdown of the quasicubic approximation.

V. CONCLUSION

Reflectance spectroscopy for nonpolar and semipolar bulk GaN substrates under uniaxial stress was performed. Bir-Pikus theory was expanded to describe the optical properties of wurtzite materials grown on semipolar planes, where the short-range Coulomb interaction was considered. Fitting the experimental results with calculations provided all the exciton deformation potentials as $C_1 = 6.5$ eV, $C_2 = 11.2$ eV, $C_3 = 4.9$ eV, $C_4 = -5.0$ eV, $C_5 = -2.8$ eV, and $C_6 = -3.1$ eV as well as bulk parameters. The obtained values clearly indicate a breakdown of the quasicubic approximation in GaN.

ACKNOWLEDGMENTS

This research was partially supported by the Ministry of Education, Culture, Sports, Science, and Technology of Japan, Grant-in-Aid for Scientific Research on Priority Areas, under Grant No. 21016006.

APPENDIX: DERIVATION OF THE UNITARY TRANSFORMATION MATRIX

In this appendix, the unitary transformation matrix, which is used to obtain oscillator strengths for arbitrary directions, is derived. In the $(x,y,z)$ orthonormal coordinates, the Schrödinger equation can be written as

$$H_{\text{Wurtzite}} |\Psi_{\text{exciton}}\rangle = E |\Psi_{\text{exciton}}\rangle.$$  (A1)

Here, we introduce an exciton wave function as $|\Psi_{\text{exciton}}\rangle' = U |\Psi_{\text{exciton}}\rangle$, where $U$ is the unitary transformation matrix. $|\Psi_{\text{exciton}}\rangle'$ is defined as

$$|\Psi_{\text{exciton}}\rangle' = a_{11} |p_+, \alpha\alpha\rangle + a_{22} |p_-, \alpha\alpha\rangle + a_{33} |p_-, \beta\alpha\rangle + a_{44} |p_+, \beta\alpha\rangle + a_{55} |p_-, \beta\beta\rangle + a_{66} |p_+, \alpha\beta\rangle + a_{77} |p_+, \alpha\beta\rangle + a_{88} |p_-, \alpha\beta\rangle + a_{99} |p_-, \beta\alpha\rangle + a_{10} |p_+, \beta\alpha\rangle + a_{11} |p_-, \beta\alpha\rangle + a_{12} |p_-, \alpha\alpha\rangle.$$  (A2)

$$\begin{equation}
|p_+, \alpha\alpha\rangle = \frac{1}{\sqrt{2}} |p_+ + ip_y\rangle, \quad |p_-, \alpha\alpha\rangle = \frac{1}{\sqrt{2}} |p_+ - ip_y\rangle. \quad (A3)
\end{equation}$$

Equation (A1) can be rewritten under the $(x', y', z')$ coordinates as

$$UH_{\text{Wurtzite}} U^{-1} |\Psi_{\text{exciton}}\rangle' = E |\Psi_{\text{exciton}}\rangle'. \quad (A4)$$

The unitary transformation matrix $U$ can be derived considering the coordinate transformation from $(x,y,z)$ to $(x',y',z')$ to give the following matrix:

![Graphical representation of the unitary transformation matrix](image-url)
\[
U = \frac{1}{2} \begin{pmatrix}
\cos \theta & 0 & 0 & 0 & 0 & 0 & 0 & 0 & \frac{\sin \theta}{\sqrt{2}} \\
0 & \frac{\cos \theta}{\sqrt{2}} & 0 & 0 & 0 & 0 & 0 & 0 & -\frac{\sin \theta}{\sqrt{2}} \\
0 & 0 & \frac{\cos \theta}{\sqrt{2}} & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & \frac{\cos \theta}{\sqrt{2}} & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & \frac{\cos \theta}{\sqrt{2}} & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & \frac{\cos \theta}{\sqrt{2}} & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & \frac{\cos \theta}{\sqrt{2}} & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & \frac{\cos \theta}{\sqrt{2}} & 0 \\
\end{pmatrix}
\]

\[\frac{1}{2} \begin{pmatrix}
\cos \theta & 0 & 0 & 0 & 0 & 0 & 0 & 0 & \frac{\sin \theta}{\sqrt{2}} \\
0 & \frac{\cos \theta}{\sqrt{2}} & 0 & 0 & 0 & 0 & 0 & 0 & -\frac{\sin \theta}{\sqrt{2}} \\
0 & 0 & \frac{\cos \theta}{\sqrt{2}} & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & \frac{\cos \theta}{\sqrt{2}} & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & \frac{\cos \theta}{\sqrt{2}} & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & \frac{\cos \theta}{\sqrt{2}} & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & \frac{\cos \theta}{\sqrt{2}} & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & \frac{\cos \theta}{\sqrt{2}} & 0 \\
\end{pmatrix}
\]

---

8yota.ishii@optomater.kuee.kyoto-u.ac.jp
9kawakami@kuee.kyoto-u.ac.jp


