Single crystal growth of the metallic triangular-lattice antiferromagnet PdCrO2

Author(s)
Takatsu, Hiroshi; Maeno, Yoshiteru

Citation
Journal of Crystal Growth (2010), 312(23): 3461-3465

Issue Date
2010-11-15

URL
http://hdl.handle.net/2433/131816

Type
Journal Article

Kyoto University
Single crystal growth of the metallic triangular-lattice antiferromagnet PdCrO$_2$

Hiroshi Takatsu, Yoshiteru Maeno

Department of Physics, Graduate School of Science, Kyoto University, Kyoto 606-8502, Japan

Abstract

We report details of single crystal growth of the metallic triangular-lattice antiferromagnet PdCrO$_2$ consisting of layers of Pd triangles and Cr triangles stacking along the c axis. We used the NaCl flux method and obtained the crystals with the size as large as 2$^3$:5$^3$:3 mm$^3$. We confirmed that single crystals have the delafossite structure with the $R\overline{3}m$ symmetry. The electrical resistivity along the c axis and that in the ab plane exhibit metallic temperature dependence with the anisotropic ratio $\rho_c/\rho_{ab}$ of over 300 at low temperatures. The residual resistivity of as small as $\rho_{0,ab}$ = 45 n$\Omega$ cm and the residual resistivity ratio of over 100 indicate high quality of the crystals investigated. Nevertheless, there is no sign of superconductivity down to 0.3 K. These crystals are useful for the investigation of anisotropic magnetic and transport properties including the unconventional anomalous Hall effect (UAHE).


PACS: 61.66.Fn, 72.15.Eb, 72.80.Ga, 75.30.Gw,

1. Introduction

Delafossites ABO$_2$ (A: noble-metal elements; B: transition-metal and rare-earth elements) crystallize in a layered structure consisting of alternating stacks of triangular lattices (TL) of monovalent $A^{1+}$ ions and trivalent $B^{3+}$ ions along the c axis (Fig. 1). Because of the simple TL structure, members with the magnetic ions at the B site are useful for the clarification of the unresolved issues of ideal TL magnets [1, 2, 3, 4]. In addition, by the choice of A site ions, the metallicity can be controlled. Insulating delafossites are actively investigated as the typical multiferroic materials. For examples, CuFeO$_2$ [5] and CuCrO$_2$ [6, 7] are known to exhibit strong coupling between the spiral magnetic ordering with the proper-screw type spin structure and the ferroelectric polarization. Metallic delafossites with frustrated spin structure are promising for the investigation of the unconventional anomalous Hall effect (UAHE), which emerges not from the conventional order parameter magnetization $\mathbf{M}$ but possibly from the multi-spin quantity “scalar spin chirality” [8].

Although most delafossites are insulators or semiconductors, there are several members of metallic delafossites, such as AgNiO$_2$ [9, 10], PdCoO$_2$ [11, 12], PtCoO$_2$ [13, 14], and PdCrO$_2$ [15]. Among them, only AgNiO$_2$ and PdCrO$_2$ are magnetic metals. In particular, since metallic and magnetic origins are well separated in PdCrO$_2$, this compound should provide unique opportunities to study the interplay between the localized frustrated spins and conduction electrons including UAHE [16].

In order to clarify the mechanism of such unconventional transport properties, determinations of the Fermi surfaces and the detailed magnetic structure are neces-


![Figure 1: (a) Unit-cell crystal structure of delafossite ABO$_2$. Both A and B ions form triangular lattices and stack along the c-axis in the sequence B-O-A-O-B. (b) Side view of the crystal structure. The oxygen ions above and below the A-site ions are not shown in this figure in order to focus on the layered structure of A and B atoms. The black lines indicate the unit cell. The symbols X, Y, and Z represent the stacking patterns of the layers. (c) Schematic drawing of the top view of the crystal structure. The definition of the symbols X, Y, and Z are the same as that in (b). These figures are generated by VESTA [17].](image)
sary. For these measurements, availability of the high-quality single crystals is crucial. Here, we present the details of the growth of high-quality PdCrO$_2$ single crystals. To the best of our knowledge, this is the first report of the details of the single crystal growth of PdCrO$_2$.

Previous studies on powder samples of PdCrO$_2$ have revealed that it exhibits an antiferromagnetic transition at $T_N = 37.5$ K forming a 120$^\circ$ spin structure [18, 15]. This Neel temperature is an order of magnitude lower than the Weiss temperature $\theta_W \approx -500$ K; the frustration parameter $\mathcal{f} = |\theta_W|/T_N$ is about 13, indicating a strong frustration among spins. Single crystals obtained in this study revealed isotropic magnetization above $T_N$, consistent with the Heisenberg spins, while they clarified the emergence of anisotropy below $T_N$, indicating that the 120$^\circ$ spin structure is easy-axis type.

2. Crystal growth

Single crystals of PdCrO$_2$ were grown by a NaCl flux method. Polycrystalline PdCrO$_2$ used for the single crystal growth were prepared in two steps. In the first step, LiCrO$_2$ was prepared as a precursor by the solid-state reaction of the stoichiometric mixture of Li$_2$CO$_3$ (99.99%, Aldrich Chemical Co.) and Cr$_2$O$_3$ (99.99%, Rare Metallic Co. Ltd.) at 850$^\circ$C in air in an alumina crucible for 24 hours. In the second step, PdCrO$_2$ powder was synthesized by the following metathetical reaction: $\text{Pd} + \text{PdCl}_2 + 2\text{LiCrO}_2 \rightarrow 2\text{PdCrO}_2 + 2\text{LiCl}$ [13, 15, 19]. For this step, Pd powder (99.99%, Furuchichi Chem. Co.) and PdCl$_2$ powder (99.999%, Aldrich Chemical Co.) were used. The mixture was ground in a mortar for 30–60 minutes, sealed in an evacuated quartz tube (100 mm × 10 mm), heated to 790$^\circ$C and kept at this temperature for 96 hours. The obtained product was washed with aqua regia and distilled water to remove LiCl, unreacted Pd, and other by-products.

Single crystals were grown from the mixture of polycrystals of PdCrO$_2$ and NaCl (99.99%, Rare Metallic Co. Ltd.) with a mass ratio of typically 1 : 10. We note here that the minimum mass ratio in which single crystals were formed was 1 : 4. The mixture was sealed in an evacuated quartz tube (130 mm × φ20 mm), heated to 880$^\circ$C in 3.5 hours, and kept at this temperature for 24 hours. It was then cooled down to 800$^\circ$C at the cooling rate of 0.25–0.5$^\circ$C/h and to 700$^\circ$C at 1$^\circ$C/h. We should note that the melting point of NaCl is 801$^\circ$C at ambient pressure. After the slow cooling processes, the quartz tube was furnace-cooled down to room temperature. Single crystals were obtained with by-products such as Cr$_2$O$_3$. In order to grow larger crystals, single crystals were added in the mixture as seeds. The largest crystal size obtained is about 2.0 × 3.5 × 0.3 mm$^3$ (see the microscope pictures in Fig. 2). We note here that single crystals of PdCrO$_2$ were not obtained either with PdCl$_2$ flux, Bi$_2$O$_3$ flux, or Bi$_2$O$_3$–B$_2$O$_3$ mixed flux. The starting PdCrO$_2$ polycrystals remained unreacted in PdCl$_2$ flux. They decomposed in Bi$_2$O$_3$ and Bi$_2$O$_3$–B$_2$O$_3$ fluxes but did not produce any crystals. We also note that attempts to grow the single crystals in air using Pt crucibles were not successful.

3. Characterization

Figure 3 represents the powder X-ray diffraction (XRD) spectra of polycrystalline and single-crystalline samples of PdCrO$_2$, with the CuK$_{\alpha1}$ radiation. These
Figure 3: Powder XRD spectra of (a) polycrystalline samples and (b) crushed single-crystalline samples of PdCrO$_2$ at room temperature. The peak pattern can be labeled with the delafossite structure with $R\bar{3}m$ symmetry. The additional peak labeled with * for the single crystal spectrum at about $2\theta = 34^\circ$ originates from an impurity phase of Cr$_2$O$_3$. It is well separated from the PdCrO$_2$ single crystals, but a small amount of powder Cr$_2$O$_3$ was inevitably picked up in the preparation for measurements.

The X-ray Laue patterns for the (001) plane and (100) plane are presented in Figs. 4(a) and (b). Calculated patterns for these conditions shown in Figs. 4(c) and (d) are consistent with the measured patterns. The clear Laue spots guarantee high crystallinity of the crystals. Scanning electron microscope (SEM) images with the energy dispersive X-ray (EDX) analysis mapping are consistent with the measured patterns. The clear patterns for these conditions shown in Figs. 4(c) and (d) are presented in Figs. 4(a) and (b). Calculated patterns for both single crystals and polycrystals.

Figure 4: X-ray Laue patterns of (a) the (001) plane and (b) the (100) plane in back-scattered conditions. The photos exhibit clear spots that agree with the $R\bar{3}m$ symmetry. (c) and (d): the corresponding calculated patterns.

Two spectra exhibit the same patterns, yielding the lattice parameters of $a = b = 2.923$ Å, $c = 18.086$ Å for both single crystals and polycrystals. The X-ray Laue patterns for the (001) plane and (100) plane are presented in Figs. 4(a) and (b). Calculated patterns for these conditions shown in Figs. 4(c) and (d) are consistent with the measured patterns. The clear Laue spots guarantee high crystallinity of the crystals.

The dc magnetization $M$ was measured with a SQUID magnetometer (Quantum Design, MPMS) from 300 K to 1.8 K in magnetic fields $\mu_0H$ between 0.01 and 7 T in both field-cooled (FC) and zero field cooled (ZFC) conditions. Figure 6(b) shows the temperature dependence of the magnetic susceptibilities ($\chi = M/\mu_0H$) measured in the field $\mu_0H = 1$ T along the $c$ axis, and in the $ab$ plane. Above $T_N$, $\chi$ is isotropic, exhibiting the Curie–Weiss behavior above 200 K with the localized moment $\mu_{\text{eff}} = 4.0 \pm 0.2$ $\mu_B$ [15]. These results suggest that PdCrO$_2$ constitutes a Heisenberg spin system, consistent with the expectation of the localized spins of the Cr$^{3+}$ ion (3d$^3$, $S = 3/2$). In contrast, below $T_N$, at which the spins form a 120$^\circ$ structure [15, 18], $\chi$ becomes anisotropic with a sudden decrease at $T_N$; the value of $\chi_{ab}$ becomes larger than that of $\chi_c$. Such anisotropy is similar to the susceptibility of another delafossite CuCrO$_2$ [7], but is different from a similar layered TL magnet LiCrO$_2$ with $\chi_c$ larger than $\chi_{ab}$ [22]. Within the mean field approximation [23], $\chi$ should be isotropic even below $T_N$ if the inter-plane spin configuration is ferromagnetic. The observed anisotropy $\chi_{ab} > \chi_c$ is expected if the inter-plane configuration is...
antiferromagnetic and the spins lie in a plane containing the $c$ axis (the so-called easy-axis type $120^\circ$ structure). We note that the anisotropy $\chi_{ab} < \chi_c$ for LiCrO$_2$ is attributed to a $120^\circ$ structure where the spins lie in the $ab$ plane (the easy-plane type structure) with an antiferromagnetic inter-plane configuration.

4. Conclusion

We have succeeded in growing single crystals of PdCrO$_2$ by a NaCl flux method. It was confirmed that the obtained crystals are in the delafossite structure with the $R3m$ symmetry. The large values of the residual resistivity ratio over 100 and the mean free path in the $ab$ plane as large as 30 $\mu$m guarantee the high quality of the crystals.

The metallic resistivity with a large anisotropy ($\rho_c/\rho_{ab} > 190$) supports the expectation that the conductivity in PdCrO$_2$ mainly originates from the Pd layers. The magnetic susceptibility above $T_N$ confirms that PdCrO$_2$ constitutes a Heisenberg spin system. Moreover, the emergence of the anisotropy at $T_N$ indicates an antiferromagnetic coupling between the layers with the easy-axis type $120^\circ$ spin structure.

These crystals make possible further investigations of the details of both the actual spin configuration and the Fermi surface topology, as well as searches of other fascinating phenomena related to the frustrated magnetism.

Figure 6: (a) Temperature dependence of the electrical resistivity along the $c$ axis ($\rho_c$) and in the $ab$ plane ($\rho_{ab}$). Both $\rho_c$ and $\rho_{ab}$ exhibit clear anomalies at $T_N$, associated with the antiferromagnetic transition. The inset shows the low temperature behavior of $\rho_c$ and $\rho_{ab}$. (b) Temperature dependence of the magnetic susceptibilities in the field at 1 T along the $c$ axis ($\chi_c$) and in the $ab$ plane ($\chi_{ab}$). Above $T_N$, $\chi_c$ and $\chi_{ab}$ are nearly equal. In contrast, below $T_N$, $\chi_c$ becomes smaller than $\chi_{ab}$. The inset focuses on the behavior near and below $T_N$. A slight splitting of the ZFC and FC curves appears at temperatures below about 10 K.

Acknowledgements

We would like to thank J. Ishikawa for his assistance of the single crystal growth. We also thank K. Ishida, S. Yonezawa, M. Kriener, D. Peets, Y. Nakai and S. Kitaoka for useful discussions and for their experimental assistance. This work was supported by the Grant-in-Aid for the Global COE Program “The Next Generation of Physics, Spun from Universality and Emergence” from the Ministry of Education, Culture, Sports, Science and Technology (MEXT) of Japan and Grants-in-Aids for Scientific Research from the Japan Society for the Promotion of Science (JSPS). H.T. is financially supported as a JSPS research fellow.
References

[20] T. Shishidou and T. Oguchi, (Personal communication). Here, the value of \( k_F \) is estimated from the band structure calculation.
[21] The carrier density \( n \) is estimated from the ordinary Hall coefficient \( R_0 \) with the single band approximation, \( R_0 = 1/ne \). It is however difficult to evaluate \( R_0 \) of PdCrO\(_2\) because the Hall resistivity in PdCrO\(_2\) contains the anomalous Hall resistivity in addition to the ordinary Hall resistivity. Thus, \( R_0 \) of PdCrO\(_2\) is approximated by \( R_0 \) of the isostructural nonmagnetic compound PdCoO\(_2\) [16]. This approximation is based on the expectation that one Pd 4d electron per formula unit dominantly contributes to the electric conductivity for both PdCrO\(_2\) and PdCoO\(_2\). Indeed, the electronic specific heat coefficient, related to the density of states of conduction carriers, is nearly the same between these two oxides.