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Atomic level observation of octahedral distortions at the perovskite oxide heterointerface

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Many functional properties observed in perovskite oxides, ABO₃, exhibit close couplings to slight structural distortions in the perovskite lattice, that consists of a three-dimensional network of corner-sharing BO₆ octahedra. The distortions can typically be categorized into three types: (i) cation displacements within octahedra, (ii) deformations (changes in size), and (iii) cooperative tilting (rotations) of the octahedra. Identification of such distortions and elucidation of their influence on properties have been important topics of investigation in fundamental materials science research as well as for applications of oxide materials in electronic devices.

It has been demonstrated that a heterostructure consisting of different ABO₃ can be used as a platform to control and tailor functionalities, as it can modify the degree of the coupling between lattices, electrons, and spins in the oxides. Heteroepitaxial thin films are formed by coherently growing oxides in such a way that the in-plane lattice parameters of the film are identical to those of the substrate. This introduces additional structural distortions as a result of the accommodation of elastic strain energy due to structural mismatch at the heterointerface. It has thus been widely believed that mismatch-induced strain is a primary cause for new properties emerging in the heterostructures. However, an atomic-scale understanding of how structural distortions are introduced to accommodate strain at the heterointerface is still missing and remains experimentally unaddressed because this requires precise measurements of the positions of both cations and oxygen.

The recent development of annular bright-field (ABF) imaging in aberration-corrected scanning transmission electron microscopy (STEM) allows simultaneous imaging of both light and heavy elements. ABF imaging combined with high angle annular dark-field (HAADF) imaging in which the contrast strongly depends on atomic number (Z)^2 would be appropriate to visualize both cation and oxygen atomic positions in the oxide heterostructures. The images acquired with STEM, however, were difficult to use for precise structural analysis because of image distortions due to drifts of both the specimen and incident probe. In order to overcome this problem and obtain structural images with minimized image distortion, we used the fast multiple-image acquisition and drift correction techniques using cross correlation of the image. In addition, we were able to precisely determine the atomic positions using Bragg filtering and cubic interpolation techniques.

Here, we focus on a SrRuO₃ (SRO) epitaxial thin film grown on a GdScO₃ (GSO) substrate. As shown in Fig. 1, bulk SRO and GSO with a Ptmn orthorhombic structure show in-phase octahedral rotation around the [001] axis and out-of-phase rotation around the [110] axis, which are described as a ∝ b ∝ c in the Glazer notation. This indicates that the images projected along the [001] direction are suitable for...
across the heterointerface. dated by displacing only the oxygen atoms within only four unit cells revealed that the mismatch of the octahedral tilt angles is accommodated at the interface despite the film being under tensile strain. We also see that except in the vicinity of the interface (yellow region), each confirmed, by X-ray reciprocal space mapping, that the fabricated (110) SRO film has a distorted orthorhombic structure with no crystallographic twins. Figure 2a shows the HAADF image of the heterostructure of the SRO film and GSO substrate taken along the [001]_ortho direction. The HAADF image which produces contrast depending on the atomic number (Z) indicates the constituent atomic positions as bright spots. The insets are simulated HAADF images for the bulk SRO and GSO, which are in good accorded with the contrast of the observed image. In the GSO substrate area, the image contrast of the Gd atomic columns is slightly distorted, reflecting the projection of displaced Gd atoms (see Fig. 1). This indicates that a slight atomic shift can be detected precisely from the HAADF image. The image also shows that there are no misfit dislocations in the heterostructure, consistent with our previous X-ray structural characterizations. Figure 2b shows intensity profiles of the HAADF image across the heterointerface. The intensity profile of A-site cations (along the red line in Fig. 2a) gives atomic positions of Sr (Z = 38) in the film (green in Fig. 2b) and Gd (Z = 64) in the substrate (purple in Fig. 2b). On the other hand, the intensity profile of B-site cations (along the blue line in Fig. 2a) gives atomic positions of Ru (Z = 44) in the film (gray in Fig. 2b) and Sc (Z = 21) in the substrate (pink in Fig. 2b). From the point where the HAADF intensities change significantly, we can identify the interface of the heterostructure. Termination of the GSO substrate is identified to be a ScO$_2$ layer and the SRO epitaxial thin film begins from the SrO layer. Figure 2c shows the ABF image taken from the same region as the HAADF image (Fig. 2a), where the atomic positions are visualized as dark contrast. We can clearly see oxygen columns in the image, which was verified from the inserted simulation images. This image allows determination of the shapes of corner-shared oxygen octahedra projected along the [001]_ortho direction as denoted by the red solid open squares. We are thus successful in investigating the connectivity of the tilted octahedra across the interface between the SRO thin film and GSO substrate.

To quantitatively analyze lattice distortions induced by the structural mismatch, we measured the positions of A-site cations from the HAADF image at sub-pixel resolution using the "Find Peaks" option incorporated in the Peak Pairs Analysis software. Figure 3a shows the variation of the cation lattice spacing extracted from the HAADF image (Fig. 2a) of which a portion is shown on the left. The interatomic distances along the out-of-plane (red solid squares) and in-plane (blue solid circles) directions were determined by averaging over 18 unit cells of the pseudocubic perovskite lattice along the [100]_ortho direction (the in-plane direction). The measured in-plane lattice parameters in the SRO film (3.957 ± 0.050 Å) are nearly equal to that of the GSO substrate (3.96 Å, purple dotted line), which is consistent with the fact that the film was coherently grown on the substrate with 1.0% tensile strain. On the other hand, the measured out-of-plane lattice parameters in the SRO film (3.915 ± 0.039 Å) have almost the same value as that of the bulk SRO (3.92 Å, green dotted line). These results indicate that the pseudocubic unit cell of the SRO film is distorted by the epitaxial strain and elongates only along the in-plane direction. Interestingly, the out-of-plane lattice parameters drastically change across the interface. We note that the ScO$_2$ layer at the interface (orange dotted line) is heavily distorted with the out-of-plane lattice parameter identical to that of the film region. This suggests that the distorted ScO$_2$ layer plays an important role in the mismatch accommodation at the interface.

Figure 3b shows the variation of the oxygen octahedral tilt angles $\theta$, which were extracted from the oxygen atomic positions in the ABF image (Fig. 2c). The projected ScO$_6$ (purple boxes) and RuO$_6$ (green boxes) octahedra are also drawn in Fig. 3b. The tilt angles were determined by averaging alternately over 18 unit cells of pseudocubic perovskite lattice along the [100]_ortho direction (the in-plane direction) because of the characteristic orthorhombic lattice. We see that except in the vicinity of the interface (yellow region), each

**Results**

A 15 nm-thick SRO thin film was grown on a (110) GSO substrate by pulsed laser deposition, as described in a previous paper. We
octahedral tilt angle in the SRO film and GSO substrate is nearly constant in the value of each bulk counterpart. Thus, the mismatch in the oxygen octahedral tilt angle is adjusted only near the SRO/GSO heterointerface, indicating that the rotation pattern of the substrate penetrates poorly into the epitaxial thin film.

The above results lead us to further investigations of the octahedral connectivity at the interface in the perovskite oxide heterostructure. The perovskite oxide epitaxial thin films often accommodate the additional octahedral rotations or tilts due to the octahedral connection between the film and substrate1–5,30,34–44. Hence, the oxygen octahedra across the heterointerface are expected to be deformed drastically. Figure 4a shows the variation of the oxygen octahedral tilt angles $\theta$ across the SRO/GSO interface. Within only four unit cells around the interface, the tilt angle $\theta$ changes from 156° for the GSO substrate to 168° for the SRO film, and the tilt angles for the film in the vicinity of the heterointerface (yellow region) are slightly smaller than those of the SRO film away from the interface. Given that the oxygen octahedra in the perovskite structure are connected by sharing the oxygen atom in its corner, the drastic changes in the octahedral tilt angle observed at the interface can be ascribed to modifications of the oxygen atom positions. We plotted the displacements of apical oxygen atoms of the octahedra along the in-plane direction, $\Delta x$, across the heterointerface in Fig. 4c. The variation of $\Delta x$ shows a similar trend to that of the tilt angles (Fig. 4a). The large displacement of 43 pm of the apical oxygen atoms in the GSO substrate is adjusted to the displacements ($\sim 21$ pm) for the SRO thin film within only four unit cells around the heterointerface. Note also that the bottom apical oxygen atoms in the topmost ScO$_6$ octahedra at the interface are also displaced about 5 pm. The octahedral tilt angle between the topmost ScO$_6$ octahedral layers in the GSO substrate (indicated by the red dashed circle in Fig. 4a) thus slightly increases from that of the bulk GSO. These results indicate the significant distortion of the topmost ScO$_6$ octahedra of the substrate, which correlates with the lattice distortion of the ScO$_2$ layer at the heterointerface as mentioned above (see Fig. 3a).

**Discussion**

Based on our observations of the cation lattice and the oxygen octahedra, we constructed an atomic structure model of the heterointerface between the SRO film and GSO substrate (Fig. 4b). Accommodation of the octahedral tilt angle mismatch between the SRO and GSO results in significant distortion of the topmost ScO$_6$ octahedral layer of the substrate as well as the in-plane displacement of the apical oxygen atoms near the heterointerface. It is worth pointing out that the in-plane cation lattice mismatch in the SRO thin film is still accommodated in the entire film region. This suggests that the mechanisms for accommodation of the cation lattice mismatch and the octahedral tilt angle mismatch are different.
lattice mismatch can be accommodated by a slight shift of the cations, and as a result, the thin films maintain the framework of the perovskite structure under the lattice mismatch-induced strain. On the other hand, the octahedral tilt is predominantly caused by the size mismatch of the constituent A- and B-site cations according to the Goldschmidt tolerance factor. Thus, the oxygen octahedral tilt angles are constant in the film and substrate, and their difference should be accommodated only near the heterointerface at which the size balance of the A- and B-site cations changes. Such mismatch accommodation in the oxygen tilt angle is achieved mainly by in-plane displacements of the apical oxygen atoms of the octahedra. This explains our observations and demonstrates that the oxygen atomic positions in the perovskite framework play a dominant role in the accommodation of the octahedral tilt angle mismatch.

In summary, from complementary HAADF- and ABF-STEM images with minimized image distortions, we were able to reveal octahedral distortions resulting from accommodation of structural mismatch at the interface between the SRO film and GSO substrate. The octahedral tilt angle in the strained SRO thin film is comparable to that in the bulk SRO, and mismatch of the octahedral tilt angle between the film and substrate is accommodated within only four unit cells around the heterointerface. This accommodation results in significant distortions of the topmost ScO$_6$ octahedral layer of the substrate as well as in-plane displacement of the apical oxygen atoms near the interface. These results highlight the fact that the oxygen arrangement around the heterointerface plays a critical role in the epitaxial strain accommodation in the perovskite oxide heterostructures.

**Methods**

**Sample preparation.** A 15 nm thick SrRuO$_3$ (SRO) epitaxial thin film was grown on a (110)$_{GSO}$ GdScO$_3$ (GSO) substrate by pulsed laser deposition (the subscript ortho denotes the orthorhombic structure). The lattice constants of the GSO substrate are $a_{\text{ortho}} = 5.45\ \text{Å}$, $b_{\text{ortho}} = 5.75\ \text{Å}$, and $c_{\text{ortho}} = 7.93\ \text{Å}$. Thus the averaged lattice mismatch between SRO ($a_{\text{ortho}} = 5.57\ \text{Å}$, $b_{\text{ortho}} = 5.53\ \text{Å}$, and $c_{\text{ortho}} = 7.85\ \text{Å}$ in bulk) and the substrate, $(a_{\text{sub}} - a_{\text{SRO}})/a_{\text{SRO}}$, where $a_{\text{sub}}$ and $a_{\text{SRO}}$ are the pseudocubic lattice parameters of the GSO substrate and the SRO film, is $+1.0\%$ (tensile strain). Details of the thin film fabrication process are given in our previous report. Briefly, the substrate temperature and partial oxygen pressure during the deposition of the SRO layer were kept at 700°C and 100 mTorr, respectively. The film thickness was determined from the period of the Laue oscillation observed in the X-ray 20-θ diffraction pattern. We confirmed that the fabricated SRO thin film has a slightly distorted orthorhombic structure and that there are no crystallographic twins.

**Microscopy.** For cross-sectional TEM observations, the thin film specimen was thinned down to electron transparency by mechanical polishing and Ar-ion milling. STEM images were acquired at room temperature in a spherical aberration corrected STEM (JEM-9900T/KP1; accelerating voltage = 200 kV, Cs = 0.025 mm, C$_\text{ab}$ = 15 mm) equipped with a cold field emission gun. The annular detection angle for HAADF was 50–133 mrad and that for ABF was 11–23 mrad because the convergent semi-angle of the incident probe was 23 mrad. 50 HAADF and 50 ABF images were acquired from the same area with a short dwell time (ca. 4.2 μs/pixel). Then, the multiple images were superimposed after correcting for the relative drifts. This procedure provides high resolution STEM images with improved signal-to-noise (SN) ratio and with minimized image distortion. Measurements of the atomic positions were carried out in the STEM images at sub-pixel resolution using Bragg filtering and cubic interpolation techniques in the “Find Peaks” option (Peak Pairs Analysis software package by HREM Research). The simulated images were obtained using an aberration-corrected probe with C$_\text{ab}$ = −0.025 mm, C$_\text{ab}$ = 15 mm and convergent.

**Simulations.** For reliable interpretation of the HAADF and ABF image contrasts, STEM image simulation was performed using multislice simulation software (WinHREM by HREM Research). The simulated images were obtained using an aberration-corrected probe with C$_\text{ab}$ = −0.025 mm, C$_\text{ab}$ = 15 mm and convergent.
The annular detection angle for HAADF in the simulations was 50–133 mrad, and that for ABF was 11–23 mrad.


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Author contributions
R.A. and D.K. conceived the idea and initiated the project. D.K. and Y.S. fabricated the samples. R.A. and H.K. collected and analyzed the STEM data. Y.S. and H.K. supervised the project. All authors discussed the experimental data and co-wrote the manuscript.

Additional information
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