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On local newforms for unramified $U(2, 1)$

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1 Introduction

Local newforms play an important role in the theory of automorphic representations. Roughly speaking, a newform for an irreducible generic representation $\pi$ of a $p$-adic group is a vector which attains the $L$-factor of $\pi$ via Rankin-Selberg integral. The existence of newforms was known only for $GL(n)$ until the work of Roberts and Schmidt [11] for GSp(4). In this note, we study newform theory for unramified $U(2, 1)$.

This note is a survey of the author’s work [6], [7], [9], [8] on newforms for unramified $U(2, 1)$. Let $G$ denote the unramified unitary group in three variables defined over a $p$-adic field of odd residual characteristic. Newforms for an irreducible generic representation $(\pi, V)$ of $G$ is defined by using a family of open compact subgroups $\{K_n\}_{n \geq 0}$ of $G$, which is an analog of paramodular subgroups of GSp(4). For each non-negative integer $n$, we denote by $V(n)$ the space of $K_n$-fixed vectors. The smallest integer such that $V(n)$ is not trivial is called the conductor of $\pi$. We write $N_\pi$ for the conductor of $\pi$, and call $V(N_\pi)$ the space of newforms for $\pi$. An algebraic structure of $V(n)$ was studied in [6] and [9], for example, the multiplicity one theorem for newforms and the dimension formula for $V(n)$, $n \geq N_\pi$.

Our main concern is the relation of newforms and Rankin-Selberg factors. Gelbart and Piatetski-Shapiro in [4] attached a family of Rankin-Selberg integrals to an irreducible generic representation $\pi$ of $G$, and defined $L$ and $\varepsilon$-factors for $\pi$. In loc. cit. they showed that the spherical vector attains the $L$-factor of $\pi$ when $\pi$ is an unramified principal series representation. But there were no results for ramified representations. In this note, we establish a theory of newforms for Gelbart and Piatetski-Shapiro’s integral. We see that (i) the newform for an irreducible generic representation $\pi$ of $G$ attains the $L$-factor of $\pi$ (Theorem 4.1) (ii) the conductor of $\pi$ coincides with the exponent of $q^{-\alpha}$ of the $\varepsilon$-factor, where $q$ denotes the cardinality of the residue field (Theorem 4.3).

We summarize the contents of this paper. In section 2, we recall from [1] the theory of Rankin-Selberg integral introduced by Gelbart, Piatetski-Shapiro and Baruch. In section 3, we define newforms for $G$ and recall their basic properties. In section 4, we show Theorems 4.1 and 4.3 assuming Lemma 4.2, which is proved in section 5.
2 Rankin-Selberg integral

In this section, we recall from [1] the theory of Rankin-Selberg integral for U(2, 1) introduced by Gelbart, Piatetski-Shapiro and Baruch.

2.1 Notation

We use the following notation. Let \( F \) be a non-archimedean local field of characteristic zero, \( \mathfrak{o}_F \) its ring of integers, and \( \mathfrak{p}_F \) the maximal ideal in \( \mathfrak{o}_F \). We fix a uniformizer \( \varpi_F \) in \( F \), and denote by \( | \cdot |_F \) the absolute value of \( F \) normalized so that \( |\varpi_F| = q_F^{-1} \), where \( q_F \) is the cardinality of the residue field \( \mathfrak{o}_F/\mathfrak{p}_F \). Throughout this paper, we assume that the characteristic of \( \mathfrak{o}_F/\mathfrak{p}_F \) is different from two.

Let \( E = F[\sqrt{\epsilon}] \) be the quadratic unramified extension over \( F \), where \( \sqrt{\epsilon} \) is a non-square unit in \( \mathfrak{o}_F \). We denote by \( \mathfrak{o}_E, \mathfrak{p}_E \) the analogous objects for \( E \). Then \( \varpi_F \) is a uniformizer of \( E \), and the cardinality of \( \mathfrak{o}_E/\mathfrak{p}_E \) is equal to \( q_F^2 \). So we abbreviate \( \varpi = \varpi_F \) and \( q = q_F \).

We realize (the group of \( F \)-points of) the unramified unitary group in three variables defined over \( F \) as

\[
G = U(2, 1) = \{ g \in \text{GL}_3(E) \mid ^t\bar{g}Jg = J \}.
\]

Here we denotes by \( - \) the non-trivial element in \( \text{Gal}(E/F) \) and

\[
J = \begin{pmatrix}
0 & 0 & 1 \\
0 & 1 & 0 \\
1 & 0 & 0
\end{pmatrix}.
\]

Let \( B \) be the upper triangular Borel subgroup of \( G \), \( U \) its unipotent radical, and \( T \) the group of the diagonal elements in \( G \). For a non-trivial additive character \( \psi_E \) of \( E \), we also denote by \( \psi_E \) the character of \( U \) defined by \( \psi_E(u) = \psi_E(u_{12}) \), for \( u = (u_{ij}) \in U \). For an irreducible generic representation \((\pi, V)\) of \( G \), we write \( \mathcal{W}(\pi, \psi_E) \) for the Whittaker model of \( \pi \) associated to \((U, \psi_E)\).

2.2 Zeta integrals

Let \( \mathcal{C}_c^\infty(F^2) \) be the space of locally constant, compactly supported functions on \( F^2 \). For an irreducible generic representation \((\pi, V)\) of \( G \), Gelbart and Piatetski-Shapiro introduced a family of zeta integrals which has the form \( Z(s, W, \Phi) (W \in \mathcal{W}(\pi, \psi_E), \Phi \in \mathcal{C}_c^\infty(F^2)) \) as follows:

We identify the subgroup

\[
H = \left\{ \begin{pmatrix}
a & 0 & b \\
0 & 1 & 0 \\
c & 0 & d
\end{pmatrix} \in G \right\}
\]
of $G$ with $U(1,1)$. Since $SU(1,1)$ is isomorphic to $SL_2(F)$, we can write any element $h$ in $H$ as

$$h = \begin{pmatrix} b & 0 \\ 0 & b^{-1} \end{pmatrix} \begin{pmatrix} \sqrt{\epsilon} & 0 \\ 0 & 1 \end{pmatrix} h_1 \begin{pmatrix} \sqrt{\epsilon^{-1}} & 0 \\ 0 & 1 \end{pmatrix},$$

where $b \in E^\times$ and $h_1 \in SL_2(F)$. For $\Phi \in C^\infty_c(F^2)$ and $h \in H$, we define a function $f(s, h, \Phi)$ on $\mathbb{C}$ by

$$f(s, h, \Phi) = |b|^s_E \int_{F^\times} \Phi((0, r) h_1) |r|^s_E d^\times r$$

by using the decomposition of $h$ in (2.1). We note that the definition of $f(s, h, \Phi)$ is independent of the choices of $b \in E^\times$ and $h_1 \in SL_2(F)$.

Set $B_H = B \cap H$ and $U_H = U \cap H$. Then $U_H$ is the unipotent radical of the Borel subgroup $B_H$ of $H = U(1,1)$. For $W \in \mathcal{W}(\pi, \psi_E)$ and $\Phi \in C^\infty_c(F^2)$, we define zeta integral

$$Z(s, W, \Phi) = \int_{U_H \backslash H} W(h) f(s, h, \Phi) dh.$$

Then $Z(s, W, \Phi)$ absolutely converges to a function in $\mathbb{C}(q^{-2s})$ if $\text{Re}(s)$ is sufficiently large.

### 2.3 $L$ and $\varepsilon$-factors

We recall the definition of $L$ and $\varepsilon$-factors attached to an irreducible generic representation $(\pi, V)$ of $G$. Set

$$I_\pi = \langle Z(s, W, \Phi) | W \in \mathcal{W}(\pi, \psi_E), \Phi \in C^\infty_c(F^2), \psi_E : \text{non-trivial} \rangle.$$

Then $I_\pi$ is a fractional ideal of $\mathbb{C}[q^{-2s}, q^{2s}]$ which contains 1. Thus, there exists a polynomial $P(X)$ in $\mathbb{C}[X]$ such that $P(0) = 1$ and $I_\pi = (1/P(q^{-2s}))$. We define the $L$-factor $L(s, \pi)$ of $\pi$ by

$$L(s, \pi) = \frac{1}{P(q^{-2s})}.$$
for all $W \in \mathcal{W}(\pi, \psi_E)$ and $\Phi \in C_c^\infty(F^2)$.

By using the above functional equation, we define the $\epsilon$-factor $\varepsilon(s, \pi, \psi_F, \psi_E)$ of $\pi$ by

$$\varepsilon(s, \pi, \psi_F, \psi_E) = \gamma(s, \pi, \psi_F, \psi_E) \frac{L(s, \pi)}{L(1-s, \tilde{\pi})},$$

where $\tilde{\pi}$ is the contragradient representation of $\pi$. By [7], we obtain $L(s, \tilde{\pi}) = L(s, \pi)$, and hence

$$\varepsilon(s, \pi, \psi_F, \psi_E) = \gamma(s, \pi, \psi_F, \psi_E) \frac{L(s, \pi)}{L(1-s, \pi)}.$$  

Thus, we can show the following proposition by the standard argument:

**Proposition 2.3.** The $\epsilon$-factor $\varepsilon(s, \pi, \psi_F, \psi_E)$ is a monomial in $q^{-2s}$ of the form

$$\varepsilon(s, \pi, \psi_F, \psi_E) = \pm q^{-2n(s-1/2)},$$

with some $n \in \mathbb{Z}$.

# 3 Newforms

In this section, we introduce a family of open compact subgroups of $G$, and define the notion of newforms for irreducible generic representations of $G$. We summarize the basic properties of newforms for $G$, which are an analog of those for $GL(n)$ and $GSp(4)$.

## 3.1 Newforms

Newforms for $G$ are defined by the following open compact subgroups $\{K_n\}_{n \geq 0}$ of $G$. For each non-negative integer $n$, we define an open compact subgroup $K_n$ of $G$ by

$$K_n = \left( \begin{array}{ccc} o_E & o_E & p_E^{-n} \\ p_E^n & 1 + p_E^2 & o_E \\ p_E^n & p_E^n & o_E \end{array} \right) \cap G.$$ 

**Remark 3.1.** The definition of $K_n$ is inspired by the paramodular subgroups of $GSp(4)$, which is used in [11]. We also note that the group $\left( \begin{array}{ccc} o_E & o_E & p_E^{-n} \\ p_E^n & 1 + p_E^2 & o_E \\ p_E^n & p_E^n & o_E \end{array} \right)^\times$ is a conjugate of the subgroup of $GL_3(E)$ which is used to define newforms for $GL_3(E)$ in [5].
For an irreducible generic representation $(\pi, V)$ of $G$, we set

$$V(n) = \{ v \in V \mid \pi(k)v = v, \ k \in K_n \}, \ n \geq 0.$$  

Then it follows from [9] that there exists a non-negative integer $n$ such that $V(n)$ is not zero.

**Definition 3.2.** We define *the conductor of $\pi$* by

$$N_\pi = \min\{ n \geq 0 \mid V(n) \neq \{0\} \}.$$  

We call $V(N_\pi)$ the space of *newforms* for $\pi$ and $V(n)$ that of *oldforms*, for $n > N_\pi$.

### 3.2 Basic properties of newforms

We recall some basic properties of newforms from [6] and [9]. Firstly, the growth of dimensions of oldforms for generic representations $\pi$ is independent of $\pi$, as in the cases of $GL(n)$ and $GSp(4)$ (see [2], [10], [11]). The following dimension formula for oldforms holds:

**Proposition 3.3** ([6], [9]). *Let $(\pi, V)$ be an irreducible generic representation of $G$. For $n \geq N_\pi$, we have*

$$\dim V(n) = \left\lfloor \frac{n - N_\pi}{2} \right\rfloor + 1.$$  

*In particular, $V(N_\pi)$ and $V(N_\pi + 1)$ are one-dimensional.*

Secondly, newforms for $G$ are test vectors for the Whittaker functional. We say that a function $W$ in $W(\pi, \psi_E)$ is a *newform* if $W$ is fixed by $K_{N_\pi}$. The following proposition is important to the application to the theory of zeta integral:

**Proposition 3.4** ([6]). *Suppose that the conductor of $\psi_E$ is $\mathfrak{o}_E$. Then for all nonzero newforms $W$ in $W(\pi, \psi_E)$, we have*

$$W(1) \neq 0.$$  

### 3.3 Zeta integral of newforms

We apply newforms for $G$ to the theory of zeta integral. We suppose that the conductor of $\psi_E$ is $\mathfrak{o}_E$. One of the nice properties of the subgroups $\{K_n\}_{n \geq 0}$ is that $K_{n,H} = K_n \cap H$ is
a maximal compact subgroup of $H$ for all $n$. Set $T_H = T \cap H$. Then we have an Iwasawa decomposition $H = U_H T_H K_{n,H}$, for any $n$. There exists an isomorphism
\[ t : E^\times \simeq T_H; a \mapsto \begin{pmatrix} a \\ 1 \\ a^{-1} \end{pmatrix} \]
For $W \in \mathcal{W}(\pi, \psi_E)$ and $\Phi \in C_c^\infty(F^2)$, we obtain
\[ Z(s, W, \Phi) = \int_{E^\times} \int_{K_{n,H}} W(t(a)k) f(s, k, \Phi)|a|_E^{-1}dkd^\times a. \]
For $n \geq 0$, we denote by $\Phi_n$ the characteristic function of $p_F^n \oplus \mathfrak{o}_F$. If $W$ a newform in $\mathcal{W}(\pi, \psi_E)$, then we have
\begin{equation} \label{eq:3.5} Z(s, W, \Phi_{N}) = \text{vol}(K_{n,H}) Z(s, W)L_E(s, 1). \end{equation}
Here $L_E(s, 1) = 1/(1 - q^{-2s})$ is the $L$-factor of the trivial representation of $E^\times$ and
\begin{equation} \label{eq:3.6} Z(s, W) = \int_{E^\times} W(t(a))|a|_E^{-1}d^\times a. \end{equation}
We note that Proposition 3.4 implies that the integral $Z(s, W)$ does not vanish for any non-zero newforms in $\mathcal{W}(\pi, \psi_E)$.
If $\psi_F$ has conductor $\mathfrak{o}_F$, then we have $\hat{\Phi}_n = q^{-N_a}ch_{\mathfrak{p}_F \mathfrak{o}_F^{-N_a}}$, and hence
\begin{equation} \label{eq:3.7} Z(1 - s, W, \hat{\Phi}_n) = q^{-2N_a(s-1/2)}Z(1 - s, W, \Phi_n) \end{equation}
by (3.5).

4 Main results

In this section, we show our two main theorems, which describe $L$ and $\epsilon$-factors of irreducible generic representations of $G$ in terms of newforms and conductors.

4.1 $L$-factors and newforms

We show that zeta integrals of newforms attain $L$-factors. We normalize Haar measures on $E^\times$ and $K_{n,H}$ so that the volumes of $\mathfrak{o}_E^\times$ and of $K_{n,H}$ are one respectively. Then the following holds:
Theorem 4.1 ([8]). Suppose that $\psi_E$ has conductor $\mathfrak{o}_E$. Let $\pi$ be an irreducible generic representation of $G$ and $W$ the newform in $\mathcal{W}(\pi; \psi_E)$ such that $W(1) = 1$. Then we have

$Z(s,W,\Phi_{N_\pi}) = L(s,\pi)$.

Theorem 4.1 is reduced to the following lemma:

Lemma 4.2. With the notation as above, we have

$Z(s,W,\Phi_{N_\pi})/L(s,\pi) = 1/L_E(s,1)$ or $1/L_E(s,1)$.

We postpone the proof of Lemma 4.2 to the next section.

Proof of Theorem 4.1. We further assume that $\psi_F$ has conductor $\mathfrak{o}_F$. Suppose that $Z(s,W,\Phi_{N_\pi})/L(s,\pi) = 1/L_E(s,1)$. Then by (2.2), we obtain

$\varepsilon(s,\pi,\psi_F,\psi_E) = \gamma(s,\pi,\psi_F,\psi_E)L(s,\pi)_{L(1-s,\pi)}$

$= \frac{Z(1-s,W,\Phi_{N_\pi})}{Z(s,W,\Phi_{N_\pi})} L(s,\pi)_{L(1-s,\pi)}$

$= q^{-2N_\pi(s-1/2)} \frac{L_E(s,1)}{L_E(1-s,1)}$.

The last equality follows from (3.7). This contradicts Proposition 2.3 which implies that $\varepsilon(s,\pi,\psi_F,\psi_E)$ is monomial. Thus we get $Z(s,W,\Phi_{N_\pi}) = L(s,\pi)$, as required. $\square$

4.2 $\varepsilon$-factors and conductors

We show that the exponent of $q^{-2s}$ of the $\varepsilon$-factor of an irreducible generic representation $\pi$ of $G$ coincides with the conductor of $\pi$. Applying the argument in the proof of Theorem 4.1, we obtain the following:

Theorem 4.3 ([8]). Suppose that $\psi_E$ and $\psi_F$ have conductors $\mathfrak{o}_E$ and $\mathfrak{o}_F$ respectively. For any irreducible generic representation $\pi$ of $G$, we have

$\varepsilon(s,\pi,\psi_F,\psi_E) = q^{-2N_\pi(s-1/2)}$.

5 Proof of Lemma 4.2

In this section, we explain how to prove Lemma 4.2.
5.1 Evaluation of $L$-factors

We shall evaluate $L(s, \pi)$, for each irreducible generic representation $(\pi, V)$ of $G$. The $L$-factor $L(s, \pi)$ is defined as the greatest common divisor of the zeta integrals $Z(s, W, \Phi)$. For $W \in \mathcal{W}(\pi, \psi_{E})$ and $\Phi \in C_{c}^{\infty}(F^{2})$, there exist $W_{i} \in \mathcal{W}(\pi, \psi_{E})$ and $\Phi_{i} \in C_{c}^{\infty}(F^{2})$ $(1 \leq i \leq m)$ such that

$$Z(s, W, \Phi) = \sum_{i=1}^{m} Z(s, W_{i}) f(s, 1, \Phi_{i}).$$

By the theory of zeta integral for $GL(1)$, we have

$$f(s, 1, \Phi_{i}) \in L_{E}(s, 1) C[q^{-2s}, q^{2s}].$$

Recall that we defined

$$Z(s, W) = \int_{E^{x}} W(t(a)) |a|_{E}^{s-1} d^{	imes}a,$$

for $W \in \mathcal{W}(\pi, \psi_{E})$. To estimate $Z(s, W)$, we can apply the theory of Kirillov model for $GL(2)$.

An irreducible generic representation of $G$ is supercuspidal, or else a subrepresentation of a parabolically induced representation from $B$. The Levi component $T$ of $B$ is isomorphic to $E^{x} \times U(1)$. For a quasi-character $\mu_{1}$ of $E^{x}$ and a character $\mu_{2}$ of $U(1)$, we denote by $\text{Ind}_{B}^{G} \mu_{1} \otimes \mu_{2}$ the corresponding parabolically induced representation. According to the classification of representations of $G$, we have the following evaluation of the shape of $L$-factors:

**Proposition 5.1.** Let $\pi$ be an irreducible generic representation of $G$.

(i) If $\pi$ is supercuspidal, then $L(s, \pi)$ divides $L_{E}(s, 1)$.

(ii) If $\pi$ is a proper submodule of $\text{Ind}_{B}^{G} \mu_{1} \otimes \mu_{2}$, then $L(s, \pi)$ divides $L_{E}(s, \mu_{1}) L_{E}(s, 1)$. 

(iii) If $\pi = \text{Ind}_{B}^{G} \mu_{1} \otimes \mu_{2}$, then $L(s, \pi)$ divides $L_{E}(s, \mu_{1}) L_{E}(s, \overline{\mu_{1}}^{-1}) L_{E}(s, 1)$.

5.2 Calculation of zeta integral of newforms

Let $W$ be the newform in $\mathcal{W}(\pi, \psi_{E})$ such that $W(1) = 1$. We shall compute $Z(s, W, \Phi_{N_{\pi}})$. Suppose that $\pi$ has conductor zero. Then $\pi = \text{Ind}_{B}^{G} (\mu_{1} \otimes 1)$, for some unramified quasi-character $\mu_{1}$ of $E^{x}$. In this case, newforms in $\mathcal{W}(\pi, \psi_{E})$ are just spherical Whittaker functions. In [4], Gelbart and Piatetski-Shapiro showed that

$$Z(s, W, \Phi_{0}) = L_{E}(s, \mu_{1}) L_{E}(s, \overline{\mu_{1}}^{-1}) L_{E}(s, 1)$$
by using Casselman-Shalika’s formula for spherical Whittaker functions in [3]. We therefore obtain $Z(s, W, \Phi_0) = L(s, \pi)$ because of Proposition 5.1.

From now on, we assume that $N_\pi$ is positive. By (3.5), we have

$$Z(s, W, \Phi_{N_\pi}) = Z(s, W)L_E(s, 1),$$

and hence it is enough to compute $Z(s, W)$. One can easily observe that

$$Z(s, W) = \int_{E^\times} W(t(a)|a|_{E}^{s-1}d^\times a = \sum_{i=0}^{\infty} W(t(\varpi^i))q^{2i(1-s)}.$$  \hspace{1cm} (5.2)

So we shall give a recursion formula for $W(t(\varpi^i)), i \geq 0$, in terms of two “Hecke eigenvalues” $\lambda$ and $\nu$.

We abbreviate $N = N_\pi$. Let us define the eigenvalue $\lambda$. We define a level raising operator $\theta' : V(N) \rightarrow V(N+1)$ by

$$\theta'v = \int_{K_{N+1}} \pi(k)vdk, v \in V(N),$$

and a level lowering operator $\delta : V(N+1) \rightarrow V(N)$ by

$$\delta w = \int_{K_N} \pi(k)wdk, w \in V(N+1).$$

Since dim $V(N) = 1$, there exists $\lambda \in \mathbb{C}$ such that

$$\lambda v = \delta \theta' v,$$

for all $v \in V(N)$.

Next, we define the eigenvalue $\nu$. Put

$$\zeta = \begin{pmatrix} \varpi & 1 \\ \varpi^{-1} & 1 \end{pmatrix} \in G.$$  \hspace{1cm} (6.1)

We define the Hecke operator $T$ on $V(N+1)$ by

$$Tv = \int_{K_{N+1}\zeta K_{N+1}} \pi(k)vdk, v \in V(N+1).$$

Because dim $V(N+1) = 1$, there exists $\nu$ in $\mathbb{C}$ such that

$$Tv = \nu v,$$

for all $v \in V(N+1)$.

With the notation as above, we obtain the following recursion formula for $W(t(\varpi^i)), i \geq 0$. 
Proposition 5.3. Let $(\pi, V)$ be an irreducible generic representation of $G$ whose conductor $N_\pi$ is positive. For any newform $W$ in $\mathcal{W}(\pi, \psi_E)$, we have

$$(\nu + q^2 - \lambda)c_i + q(\nu + q^2 - q^3)c_{i+1} = q^5c_{i+2}, \quad i \geq 0,$$

$$(\nu - q^3)c_0 = q^4c_1,$$

where $c_i = W(t(w^i)), \ i \geq 0$.

By (5.2) and Proposition 5.3, we can describe the zeta integral of newforms in terms of $\lambda$ and $\nu$:

Proposition 5.4. Let $(\pi, V)$ be an irreducible generic representation of $G$ whose conductor $N_\pi$ is positive and $W$ its newform in $\mathcal{W}(\pi, \psi_E)$ such that $W(1) = 1$. Then we have

$$Z(s, W) = \frac{1 - q^{-2s}}{1 - \frac{\nu + q^2 - q^3}{q^2}q^{-2s} - \frac{\nu + q^2 - \lambda}{q}q^{-4s}}.$$ 

In particular,

$$Z(s, W, \Phi_{N_\pi}) = \frac{1}{1 - \frac{\nu + q^2 - q^3}{q^2}q^{-2s} - \frac{\nu + q^2 - \lambda}{q}q^{-4s}}.$$ 

5.3 Proof of Lemma 4.2

We have seen that Lemma 4.2 holds for the unramified principal series representations.

Let $\pi$ be an irreducible generic representation of $G$. We assume that $N_\pi$ is positive. Proof of Lemma 4.2 is done by comparing Propositions 5.1 and 5.4. Suppose that $\pi$ is supercuspidal or a subrepresentation of $\text{Ind}_{B}^{G}\mu_1 \otimes \mu_2$, for some ramified quasi-character $\mu_1$ of $E^\times$. Then it follows from Proposition 5.1 that $L(s, \pi) = 1$ or $L_E(s, 1)$. By definition, we have $Z(s, W, \Phi_{N_\pi})/L(s, \pi) \in \mathbb{C}[q^{-2s}, q^{2s}]$. So we get

$$Z(s, W, \Phi_{N_\pi})/L(s, \pi) = 1 \text{ or } 1/L_E(s, 1),$$

by Proposition 5.4.

Suppose that $\pi$ is a subrepresentation of $\text{Ind}_{B}^{G}\mu_1 \otimes \mu_2$, for unramified $\mu_1$. Then we can regard newforms for $\pi$ as functions in $\text{Ind}_{B}^{G}\mu_1 \otimes \mu_2$. Due to [6], non-zero newforms $f$ in $\pi$ satisfy $f(1) \neq 0$. By using this property of newforms, we can compute the eigenvalues $\nu$ and $\lambda$ explicitly, and Lemma 4.2 follows.
References


