(pss-logo will be inserted here by the publisher)

Bulk crystal growth and characterization of ZnSnP₂ compound semiconductor by flux method

Shigeru Nakatsuka¹, Hiroshi Nakamoto¹, Yoshitaro Nose^{*,1,2}, Tetsuya Uda¹ and Yasuharu Shirai¹

¹ Department of Materials Science and Engineering, Kyoto University, Kyoto 606-8501, Japan
² JST-PRESTO, Japan

Received ZZZ, revised ZZZ, accepted ZZZ Published online ZZZ (Dates will be provided by the publisher.)

Keywords. Chalcopyrite, Crystal growth, Flux method, Order-disorder

ZnSnP₂ is a promising candidate for solar absorber materials from the viewpoint of high absorption and earthabundant constitution elements. In this paper, we fabricated ZnSnP₂ crystals by flux method based on the phase diagram of Sn-ZnP₂ pseudo-binary system and investigated their properties for an application to photovoltaics. The crystal growth experiments with the cooling rate of 0.7 and 12 °C/h were carried out and we successfully obtained ZnSnP₂ crystals with the diameter of 8mm and the thickness of a few mm by a slow cooling rate. The structure of grown crystals studied by X-ray diffraction was indicated to be chalcopyrite-type ZnSnP₂. In addition, the

1 Introduction Recently, photovoltaic devices using various type of materials have been investigated for the increasing energy demand and low-cost. Above all, solar cells using CdTe with II-IV type semiconductor and Cu(In,Ga)Se₂ (CIGS) with chalcopyrite-type structure have high efficiencies, 20.4% and 20.8%, respectively [1]. Although the above solar cells are in practical use in late years, it is not desirable to use toxic elements, Cd and Te, and minor elements, In and Ga. On these backgrounds, solar absorbing materials consisting of earth-abundant and low-toxic elements have been researched and solar cells using Cu₂ZnSnS_{4-x}Se_x achieved a conversion efficiency of 12.6% [2], which is insufficient compared to CIGS.

In this work, we thus concentrated on a ternary compound semiconductor ZnSnP₂ with chalcopyrite structure as a new solar absorber. This compound also consists of earth-abundant and low-toxic elements. It was also reported that ZnSnP₂ thin films fabricated by coevaporation had a p-type conduction, a direct bandgap of 1.62 eV and a high absorption coefficient ($\alpha > 10^5$ cm⁻¹) in visible light range [3]. Therefore, ZnSnP₂ is a promising decrease of the degree of order was observed with the increase of cooling rate. The lattice constants of a and c axes are 5.649 and 11.295 Å, respectively. The composition of grown crystals is a near stoichiometric ratio of ZnSnP₂ by EDX analysis. The bandgaps of ZnSnP₂ crystals obtained by cooling rate of 0.7 and 12 °C/h were estimated to be 1.61 and 1.48 eV, respectively, which is caused by the difference of the degree of order. The hall-resistivity measurement showed that ZnSnP₂ crystals with a slow cooling rate has a p-type conduction. The resistivity, the hole concentration and the mobility are 10~70 Ω cm, 6·10¹⁶~2·10¹⁷ cm⁻³, and 1~3 cm²V⁻¹s⁻¹. The obtained properties are suitable for an absorber of photovoltaics.

Copyright line will be provided by the publisher

material for a solar absorbing layer in photovoltaic devices. For the investigation of ZnSnP2 properties, several studies on crystal growth by flux method using Sn solution were reported [4-10]. In these reports, it was shown that ZnSnP₂ crystals with chalcopyrite or sphalerite structures were obtained depending the cooling rate during crystal growth. The change of the crystal structure is attributed to the orderdisorder transition, where the ordered and the disordered phases are chalcopyrite and sphalerite phases, respectively [11]. In the disordered phase, Zn and Sn atoms ramdomly occupy on their sites. The reported properties of ZnSnP2 are summarized in Table 1. The bandgap of ZnSnP2 with sphalerite structure is in the range of 1.22 to 1.32 eV, while that of chalcopyrite phase is significantly different from literatures: 1.2 to 1.67 eV. In the reported works, ZnSnP₂ with chalcopyrite structure was obtained with the cooling rate of 5 or 10 °C/h, which is relatively rapid in the viewpoint of crystal growth. This rapid cooling rates led to the precipitation of ZnSnP₂ crystals in the form of plates with the small dimensions of about $4 \times 4 \times 0.3$ mm [5]. In addition, it is considered that one of the reasons for wide

^{*} Corresponding author: e-mail nose.yoshitaro.5e@kyoto-u.ac.jp, Phone: +81-75-753-5472, Fax: +81-75-753-3579

range of reported bandgaps is the intermediate structure between chalcopyrite and sphalerite, which is a partly-disordred state.

In this report, we thus aimed to fabricate $ZnSnP_2$ crystals with a larger size to be available for various evaluations and investigate the effect of cooling rate on their properties. Figure 1 shows the phase diagram of the Sn-ZnP₂ pseudobinary system established by our previous work [12]. In the crystal growth of ZnSnP₂, it is difficult to obtain crystals with a single-phase from congruent solidification because the peritectic reaction between ZnP₂ and liquid Sn is observed. Therefore, it is understood that the usage of Sn solution with the composition of more than 90 mol%Sn is necessary for ZnSnP₂ crystal growth.

Table 1 Properties of ZnSnP ₂ crystals grown by flux metr

Pha	se	Coo ra / °C	ling te h ⁻¹	Bandgap / eV	Hole concentration / cm ⁻³	Mobility / cm ² V ⁻¹ sec ⁻¹	Ref.
		4	5	1.2	$7.3 \cdot 10^{16}$	68	[4]
		1	0			30	[5]
					$(5.0 \sim 7.3) \cdot 10^{16}$	48~68	[6]
		4	5	1.46			[7]
Chalcop	e		1.45	$(3 \sim 5) \cdot 10^{16}$	50~60	[8]	
				1.67	$(1.2 \sim 1.6) \cdot 10^{18}$	5~10	[9]
		4	5	1.64			[10]
		0	07	1.61	$6 \cdot 10^{16}$	1~3	*
		0.	. /	1.01	$\sim 2 \cdot 10^{17}$		
		2	5		$6.8 \cdot 10^{16}$	24	[4]
Partly disorder		5	0		$7.2 \cdot 10^{16}$	9.3	[4]
		7	5		$6.6 \cdot 10^{16}$	8	[4]
		12	00		$8.8 \cdot 10^{16}$	4.5	[4]
		1	2	1.48			*
Sphalerite					$7.2 \cdot 10^{16}$	(27)	[7]
					$\sim 1.8 \cdot 10^{17}$	0.3~7.0	[0]
		7	5	1.22			[7]
				1.32	$(1 \sim 2) \cdot 10^{16}$	6~10	[8]
		13		1.25			[10]
* this work							
	1100						
	1000				Liq.		
	ပ	900-				· · · · · ·	
0							
	Temperature	800-		Liq. + ZhP ₂			
		700-	i ⁻		+ ZnSnP₂		
		600-					
		500-		Liq		not investi- gated	
		400-					

Figure 1 Phase diagram of Zn-SnP₂ pseudo binary system [12].

50

Composition of Sn / mol%

ZnŚ

300 | 4 100

Sn

2 Experimental procedure Raw materials, Zn shots (99.9999%, Kojundo Kagaku), Sn shots (99.99%, Kojundo Kagaku) and red phosphorus flacks (99.9999%, Kojundo Kagaku) were sealed in evacuated quartz ampule with the inner diameter of 6 or 10 mm under the pressure below 10⁻² Pa. Zn and Sn were chemical-etched using 0.1M HCl solution before sealed in the ampule. The nominal composition was controlled to be 92 mol%Sn determined by the phase diagram of Sn-ZnP₂ system shown in Fig. 1. Then, the quartz ampule was set in the growth furnaces shown in Fig. 2, which has three heaters to control temperature gradient precisely. The sealed ampule was located at the position around 700 °C for homogenization. After that, the furnaces were lifted up by the rate of 0.2 or 3.5 mm/h and the sample was unidirectionally solidified from the bottom. The temperature of the bottom of the ampule was monitored by Ktype thermocouple during crystal growth. After the growth experiments, bulk crystals obtained by the rate of 0.2 mm/h were cut into several wafers in the direction perpendicular to the growth direction. The surface of each wafer was mechanically polished with a series of emery papers and finally with 1 µm diamond slurry on a buff sheet. Some grown crystals were ground into a fine powder by mortar and pestle. On the other hand, crystals by the rate of 3.5 mm/h were extracted from the sample by dissolving Sn in 0.1 M HCl solution and formed into powder. The product phase was identified by X-ray diffraction (XRD, Panalytical X'Pert Pro) and the composition of the crystals was analysed by SEM (KEYENCE VE-7800) and EDX (EDAX VE-9800). The lattice constant was evaluated from XRD profiles of powder specimens by Rietveld analysis using a software X'Pert HighScore Plus. The bandgap was estimated by diffuse reflectance spectra of powder specimens using Spectrophotometer (HITACHI U-3500). The Hall effect and the resistivity measurements were carried out at room temperature based on van der Pauw method by Hall-Resistivity measurement system (Toyo Technica, ResiTest 8300) and the density and the mobility of carrier were evaluated.



Figure 2 Schematic illustration of furnace for crystal growth.

3 Results and discussion

3-1. Dependence of cooling rate on precipitation behavior

Figure 3 shows the photographs of samples after crystal growth by the furnace lifting with 0.2 and 3.5 mm/h, which correspond to a cooling rate of 0.7 and 12 °C/h, respectively. In the sample obtained by the slow cooling rate, the precipitation of ZnSnP2 crystal was observed at the bottom of the sample as shown in Fig. 3(a). The obtained crystal was a polycrystal with a diameter of 8 mm and a length of 10 mm which was larger compared to reported works. The grain size was about 1mm as shown in Fig. 3(c). On the other hand, the rapid cooling rate led to a dispersed precipitation of small crystals shown in Fig. 3(b). In general, a supersatuation of solute is necessary for a precipitation of crystal from solution. In this study, it is needed that Zn and P diffuse to the interface between the crystal and the solution until enough supersaturation for crystal growth. The diffusion coefficient of Zn in liquid Sn at 700 °C is calculated to be about $5 \times 10^{-5} \text{ cm}^2 \text{sec}^{-1}$ [13]. The mean diffusion length of Zn is evaluated to be 4 mm/h using \sqrt{Dt} . The slow lifting speed of 0.2 mm/h leads to an enough time for diffusion of Zn to the bottom of the sample and the precipitation can occur at the bottom. In contrast, the solidification of the solution occurred before the solute was supersaturated at the bottom in the case of the lifting speed of 3.5 mm/h because the velocity of both solidification and diffusion of Zn in solution are comparable. Consequently, the dispersed precipitations of ZnSnP₂ were observed in the sample by the rapid cooling rate.

3-2. Investigation on some properties

Figure 4 shows the powder XRD profiles of the grown crystals, which indicate that the grown crystals are $ZnSnP_2$ with a single phase. Here, the intensity of superlattice reflection in the crystals decrease with the increase of the cooling rate. The degree of order was evaluated using the ratio of integrated intensities of (101) and (112) peaks, which are superlattice and fundamental reflections, respectively [14]. In this work, the ratio decreases from 0.076



Figure 3 Photographs of samples after crystal growth with the cooling rate of (a) 0.7 and (b) 12 °C/h, and (c) a wafer of grown crystal with the rate of 0.7 °C/h.

to 0.057 with increasing cooling rate, which means the decrease of the degree of order. The reported works also show a similar trend [4, 6-8, 10]. Consequently, the lattice constants of a and c axes were evaluated to be 5.649 and 11.295 Å, respectively, using the XRD profiles of the crystal obtained by the slow cooling rate. The ratio c/a is 2. The composition of the samples was an almost stoichiometric ratio of ZnSnP₂, which was analysed by EDX.

In order to evaluate bandgaps of the crystals, the diffuse reflectance, R, is convert to Kubelka-Munk function, F(R), which represents the light absorption as expressed in equation (1) [15],

$$F(R) = \frac{(1-R)^2}{2R} = \frac{\alpha}{S},$$
 (1)

where α and *S* are the absorption and the scattering coefficients, respectively. In the case of a direct transition, the absorption coefficient is correlated to photon energy by Tauc's equation,

$$\alpha = \frac{A(E - E_g)^{1/2}}{E}, \qquad (2)$$

where E and A are photon energy and a constant which depends on properties of materials. Consequently, the energy bandgap, E_g , is evaluated by Tauc's plot of the Kubelka-Munk function based on equation (3) [16].



Figure 4 Powder XRD profiles of grown crystals.



Figure 5 Tauc's plot of Kubelka-Munk transform using the diffuse reflectance measurements of powder samples.

$$\left[F(R)E\right]^2 = \left(\frac{\alpha E}{S}\right)^2 = \left(\frac{A}{S}\right)^2 (E - E_g).$$
(3)

The bandgaps of ZnSnP₂ crystals prepared by the cooling rate of 0.7 and 12 °C/h were estimated to be about 1.61 and 1.48 eV, respectively, as shown in Fig. 5. The former is close to the value reported by Ryan *et al.* [5]. In addition, they also reported that bandgap of ZnSnP₂ with sphalerite structure is 1.25 eV, while Rubenstein *et al.* reported the bandgap of ZnSnP₂ with chalcopyrite is 1.2 eV. In this case, the structure of ZnSnP₂ might be close to sphalerite. Moreover, Berkovskii *et al.* [7] and Goryunova *et al.* [8] reported the bandgap of chalcopyrite is 1.46 and 1.45 eV, respectively, which are close to the value of the ZnSnP₂ crystal obtained by the cooling rate of 12 °C/h in this work. Therefore, their reported crystals are also considered to be an intermediate state between chalcopyrite and sphalerite structure, which is a partly-disordered state.

The electric properties of $ZnSnP_2$ crystals with the cooling rate of 0.7 °C/h were evaluated using the Hall and resistivity measurements. The crystals obtained by the cooling rate of 12 °C/h was too small to evaluate the electric properties. The Hall effect measurement indicates that $ZnSnP_2$ prepared in this work shows a p-type conduction. The resistivity, the hole concentration and the mobility are 10~70 Ω cm, $6 \cdot 10^{16} \sim 2 \cdot 10^{17}$ cm⁻³, and $1 \sim 3$ cm²V⁻¹s⁻¹, respectively. The hole concentration of the crystal in this work is close to that in reported works, which is suitable for solar cell devices. On the other hand, the mobility is about one order of magnitude smaller than the literature data. It is likely that the difference of the mobility is necessary to clarify.

4 Conclusions In this paper, ZnSnP₂ crystal growth was carried out by flux method with the cooling rate of 0.7 and 12 °C/h. The, polycrystals with the diameter of 8mm and the thickness of a few mm were obtained with the slow cooling rate, while the rapid cooling rate led to the dispersed precipitation of small crystals. The precipitation behaviours can be understood based on the diffusion of Zn in liquid Sn. The precipitated crystals with the cooling rate of 0.7 and 12 °C/h were identified as ZnSnP2 by XRD, however the decrease of the degree of order evaluated from XRD proriles was observed with increasing cooling rate. The bandgaps of ZnSnP2 with the cooling rate of 0.7 and 12 °C/h were calculated to be 1.61 and 1.48 eV using diffuse reflectance spectra, which was due to the difference of the degree of order. The hall-resistivity measurement showed that ZnSnP2 crystals with slow cooling rate have a p-type conduction, and the electric properties suggest that ZnSnP2 with chalcopyrite structure is suitable for a solar absorbing material.

Acknowledgements The authors thanks to Dr. H. Hayashi for his support on hall-resistivity measurements. This work was partly supported by JSPS KAKENHI Grant Number 21760599, and Elements Science and Technology project from MEXT, Japan, and PRESTO from Japan Science and Technology Agency.

References

- M. A. Green, K. Emery, Y. HIshikawa, W. Warta and E. D. Dunlop, Prog. Photovolt: Res. Appl. 22, 1 (2013).
- [2] W. Wang, M. T. Winkler, O. Gunawan, T. Gokmen, T. K. Todorov, Y. Zhu and D. B. Mitzi, Adv. Energy Mater., doi: 10.1002/aenm.201301465 (2013).
- [3] H. Y. Shin and P. K. Ajmera, Mater. Lett. 5, 211 (1987).
- [4] A. A. Vaipolin, N. A. Goryunova, L. I. KleShohinskii, G. V. Loshakova and E. O. Osmanov, Phys. Stat. Sol. 29, 435 (1968).
- [5] M. Rubenstein and R. W. Ure, J. Phys. Chem. Solids, 29, 551 (1968).
- [6] N. A. Goryunova, F. P. Kesamanly and G. V. Loshakova, Sov. Phys. Semicond., 1, 844 (1968).
- [7] F. M. Berkovskii, D. Z. Garbuzov, N. A. Goryunova, G. V. Loshakova, S. M. Ryvkin and G. P. Shpen'kov, Sov. Phys. Semicond., 2, 618 (1968).
- [8] N. A. Goryunova, M. L. Belle, L. B. Zlatkin, G. V. Loshakova, A. S. Poplavnoi and V. A. Chaldyshev, Sov. Phys. Semicond., 2, 1126 (1969).
- [9] A. A. Abdurakhimov, L. V. Kradinova, Z. A. Parimbekov and Yu. V. Rud', Sov. Phys. Semicond. 16, 156 (1982).
- [10] M. A. Ryan, M. W. Peterson, D. L. Williamson, J. S. Frey, G. E. Maciel and B. A. Parkinson, 2, 528 (1987).
- [11] G. A. Seryogin, S. A. Nikishin, H. Temkin, A. M. Mintairov, J. L. Merz and M. Holtz, Appl. Phys. Lett., 74, 2128 (1999).
- [12] Y. Nose and T. Uda, Technical Digest of 21 st. International Photovoltaic Science and Engineering Conference, 2D-3P-35 (2011).
- [13] C. H. Ma and R. A. Swalin, Acta. Met., 8, 388 (1960).
- [14] B. E. Warren, X-ray diffraction (Addison-Wesley, Boston, 1969).
- [15] H. Lin, C. P. Huang, W. Li, C. Ni, S. I. Shah and Y. H. Tseng, Appl. Catal. B: Environ., 68, 1 (2006).
- [16] J. Tauc, R. Grigorovici and A. Vancu, Phys. Satatus Solidi, 15, 627 (1966).