

## ON ABELIAN CONFORMAL FIELD THEORY

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Abelian conformal field theory is usually discussed from the view point of the universal Grassmann manifold and Krichever maps ([1]). Here, we consider it from the view point of non-abelian conformal field theory developed in [2]. We take the Heisenberg algebra as a gauge group. In the following we shall show that the main ideas of the paper [2] can be applied to our situation.

We thank A. Tuchiya for pointing out a gap of our original proof of the main theorem and showing us an idea of a proof of Lemma 2.3 below.

## §1. Main Theorem

For a positive even integer  $M$  we let  $H_M$  be a Heisenberg algebra generated by operators  $a(n)$ ,  $n \in \mathbf{Z}$  with commutation relation

$$(1.1) \quad [a(n), a(m)] = Mn\delta_{n+m,0} \cdot id.$$

The Heisenberg algebra is a universal enveloping algebra of an affine Lie algebra  $\{a(n)\}$  associated with a one-dimensional abelian Lie algebra  $\mathbf{C}$  with commutation relation (1.1). For each  $p \in \mathbf{C}$ , by  $\mathcal{F}(p)$  we denote an irreducible highest weight module of  $H_M$  determined by

$$\begin{aligned} a(0)|p\rangle &= p|p\rangle \\ a(n)|p\rangle &= 0, \quad \text{if } n \geq 1, \end{aligned}$$

where  $|p\rangle$  is a highest weight vector. Let  $t_0, t_1, t_2, \dots$  be independent variables. Put

$$\begin{aligned} a(m) &= \frac{\partial}{\partial t_m}, \quad m = 0, 1, 2, \dots \\ a(-n) &= nMt_n, \quad n = 1, 2, 3, \dots \end{aligned}$$

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Then, the Heisenberg algebra  $H_M$  and its irreducible module  $\mathcal{F}(p)$  are realized as

$$H_M = \mathbf{C}[t_1, t_2, \dots, t_n, \dots, \frac{\partial}{\partial t_0}, \frac{\partial}{\partial t_1}, \dots, \frac{\partial}{\partial t_m}, \dots, \frac{\partial}{\partial t_m}, \dots]$$

$$\mathcal{F}(p) = \mathbf{C}[t_1, t_2, \dots, t_n, \dots, e^{pt_0}, e^{-pt_0}],$$

where the highest weight vector  $|p\rangle$  corresponds to  $e^{pt_0}$ . Using there realization, let us introduce an operator  $\hat{q}$  as

$$\hat{q} = Mt_0.$$

Put

$$\phi(z) = \hat{q} + a(0) \log z - \sum_{n \neq 0} \frac{a(n)}{n} z^{-n}$$

$$a(z) = \sum_{n \in \mathbf{Z}} a(n) z^{-n-1}$$

Then we have

$$d\phi(z) = a(z)dz$$

For each integer  $k$ , the Vertex operator  $V_{kM}(z)$  is defined as

$$V_{kM}(z) = \circ \circ e^{k\phi(z)} \circ \circ$$

where  $\circ \circ$  is a normal ordering defined by putting  $a(n)$ ,  $n \geq 0$  the right hand side and  $\hat{q}$ ,  $a(-n)$ ,  $n \geq 1$  the left hand side. Hence, we have

$$V_{kM}(z) = e^{k \sum_{n=1}^{\infty} \frac{a(-n)}{n} z^n} e^{k\hat{q}} e^{ka(0) \log z} e^{-k \sum_{n=1}^{\infty} \frac{a(n)}{n} z^{-n}}$$

The Vertex operator  $V_{kM}(z)$  is an intertwiner between the representations  $\mathcal{F}(p)$  and  $\mathcal{F}(kM + p)$ . Note that in conformal field theory  $a(z)$  behaves as a one-form and  $V_{kM}(z)$  behaves as a  $\frac{k^2}{2}M$ -form. The energy-momentum tensor  $T(z)$  is defined as

$$T(z) = \frac{1}{2M} \circ \circ a(z)a(z) \circ \circ$$

There is a formal expansion

$$T(z) = \sum_{n \in \mathbf{Z}} L_n z^{-n-2},$$

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and  $\{L_n\}$  is a Virasoro algebra. In the following we only consider irreducible highest weight representations of  $H_M$  with highest weight vectors  $|p\rangle$  where  $p$ 's are *integers*.

Let  $\Lambda = \{\bar{0}, \bar{1}, \dots, \overline{M-1}\}$  be representatives of the module  $\mathbf{Z}/M\mathbf{Z}$ . For each  $\bar{p} \in \{\bar{0}, \bar{1}, \dots, \overline{M-1}\}$ , put

$$\mathcal{H}(\bar{p}) := \bigoplus_{p \equiv \bar{p} \pmod{M}} \mathcal{F}(p).$$

Let  $\mathfrak{X} = (C; Q_1, \dots, Q_N; \xi_1, \dots, \xi_N)$  be an  $N$ -pointed stable curve of genus  $g$  with formal neighbourhoods. To each point  $Q_j$  we associate an element  $\bar{p}_j \in \Lambda$  and put

$$\begin{aligned} \vec{p} &= (\bar{p}_1, \bar{p}_2, \dots, \bar{p}_N), \\ \mathcal{H}(\vec{p}) &= \mathcal{H}(\bar{p}_1) \otimes \mathcal{H}(\bar{p}_2) \otimes \dots \otimes \mathcal{H}(\bar{p}_N) \end{aligned}$$

Put also

$$\mathcal{H}^\dagger(\vec{p}) = \text{Hom}_{\mathbf{C}}(\mathcal{H}(\vec{p}), \mathbf{C}).$$

We have a natural pairing

$$\begin{aligned} \mathcal{H}^\dagger(\vec{p}) \times \mathcal{H}(\vec{p}) &\rightarrow \mathbf{C} \\ (\langle \psi |, |\phi \rangle) &\mapsto \langle \psi | \phi \rangle \end{aligned}$$

where  $\langle \psi | \phi \rangle$  means  $\psi(|\phi\rangle)$ .

**Definition 1.1.** The space of vacua  $\mathcal{V}_{\vec{p}}^\dagger(\mathfrak{X})$  attached to the  $N$ -pointed stable curve with formal neighbourhoods  $\mathfrak{X}$  is a subspace of  $\mathcal{H}^\dagger(\vec{p})$  consisting of vectors  $\langle \psi |$  satisfying the following conditions.

(1) For each  $|\phi\rangle \in \mathcal{H}(\vec{p})$ , the data  $\langle \psi | \rho_j(a(\xi_j)) | \phi \rangle d\xi_j$ ,  $j = 1, 2, \dots, N$  are the Laurent expansions of an element  $\omega \in H^0(C, \omega_C(*\sum Q_j))$  at  $Q_j$ 's with respect to the formal coordinates  $\xi_j$ 's,

(2) For each  $|\phi\rangle \in \mathcal{H}(\vec{p})$ , the data  $\langle \psi | \rho_j(V_{\pm M}(\xi_j) | \phi \rangle (d\xi_j)^{\frac{M}{2}}$ ,  $j = 1, 2, \dots, N$ , are the Laurent expansions of an element  $\tau \in H^0(C, \omega_C^{\otimes \frac{M}{2}}(*\sum Q_j))$  at  $Q_j$ 's with respect to the formal coordinates  $\xi_j$ .

**Main Theorem.** We have

$$\dim_{\mathbf{C}} \mathcal{V}_{\vec{p}}^\dagger(\mathfrak{X}) = \begin{cases} M^g, & \text{if } \bar{p}_1 + \dots + \bar{p}_N = \bar{0} \\ 0, & \text{otherwise} \end{cases}$$

where  $g$  is the genus of the stable curve  $C$ .

## §2. Outline of a proof of Main Theorem.

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First we shall rewrite the conditions (1), (2) in Definition 1.1. Note that the condition (1) is equivalent to the condition

$$(1^*) \quad \sum_{j=1}^n \operatorname{Res}_{\xi_j=0} (\langle \psi | \rho_j(a(\xi_j)) | \phi \rangle g(\xi_j) d\xi_j) = 0$$

for every  $g \in H^0(C, \mathcal{O}_C(*\sum Q_j))$ , where  $g(\xi_j)$  is the Laurent expansion of  $g$  at  $Q_j$ . The condition (2) is equivalent to the condition

$$(2^*) \quad \sum_{j=1}^N \operatorname{Res}_{\xi_j=0} (\langle \psi | \rho_j(V_{\pm M}(\xi_j)) | \phi \rangle h(\xi_j) d\xi_j) = 0$$

for every  $h \in H^0(C, \omega_C^{\otimes(1-\frac{M}{2})}(*\sum Q_j))$ , where  $h(\xi_j)(d\xi_j)^{\frac{M}{2}}$  is the Laurent expansion of  $h$  at  $Q_j$ . In the following we choose integers  $p_j$  such that  $p_j \equiv \bar{p}_j \pmod{M}$ . Put

$$|p_1, p_2, \dots, p_N\rangle = |p_1\rangle \otimes |p_2\rangle \otimes \cdots \otimes |p_N\rangle.$$

Apply the condition (1\*) to an element  $\langle \psi | \in \mathcal{V}_{\bar{p}}^+(\mathfrak{X})$  and  $1 \in H^0(C, \mathcal{O}_C(*\sum Q_j))$ . Since we have

$$\operatorname{Res}_{\xi_j=0} \{(a(\xi_j) | p_j \rangle d\xi_j\} = a(0) | p_j \rangle = p_j | p_j \rangle,$$

the condition (1\*) implies that

$$\left( \sum_{j=1}^N p_j \right) \langle \psi | p_1, p_2, \dots, p_N \rangle = 0.$$

Hence, if  $\langle \psi | p_1, p_2, \dots, p_N \rangle \neq 0$ , then  $\sum_{j=1}^N p_j = 0$ .

First let us consider an  $N$ -pointed projective line  $(\mathbf{P}^1(\mathbf{C}); a_1, a_2, \dots, a_N)$  with  $a_1 = 0, a_2 = 1, a_N = \infty$ . Let  $z$  (resp.  $w$ ) be a coordinate of an affine line in  $\mathbf{P}^1(\mathbf{C})$  containing  $0$  (resp.  $\infty$ ) with  $z \cdot w = 1$ . Put

$$(2.1) \quad \xi_j = \begin{cases} z - a_j, & j = 1, 2, \dots, N-1 \\ w, & j = N, \end{cases}$$

and

$$\mathfrak{X} = (\mathbf{P}^1(\mathbf{C}); a_1, a_2, \dots, a_N; \xi_1, \xi_2, \dots, \xi_N).$$

First we shall prove the following proposition.

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**proposition 2.1.**

$$\dim_{\mathbf{C}} \mathcal{V}_{\vec{p}}^{\dagger}(\mathfrak{X}) = \begin{cases} 1, & \text{if } \bar{p}_1 + \bar{p}_2 + \cdots + \bar{p}_N = 0 \\ 0, & \text{otherwise} \end{cases}$$

Let  $F_0\mathcal{H}(\bar{p}_j)$  be a subspace of  $\mathcal{H}(\vec{p})$  spanned by the highest weight vectors  $|lM + p_j\rangle$ ,  $l \in \mathbf{Z}$  over  $\mathbf{C}$ . Put

$$F_0\mathcal{H}(\vec{p}) = F_0\mathcal{H}(\bar{p}_1) \otimes F_0\mathcal{H}(\bar{p}_2) \otimes \cdots \otimes F_0\mathcal{H}(\bar{p}_N).$$

To prove the above proposition we need the following lemma.

**Lemma 2.2.** *Under a natural mapping*

$$j : \text{Hom}_{\mathbf{C}}(\mathcal{H}(\vec{p}), \mathbf{C}) \longrightarrow \text{Hom}_{\mathbf{C}}(F_0\mathcal{H}(\vec{p}), \mathbf{C}),$$

the space of vacua  $\mathcal{V}_{\vec{p}}^{(\dagger)}(\mathfrak{X})$  of the  $N$ -pointed projective line with coordinates (2.1) is mapped injectively.

The lemma and the above consideration imply

$$\mathcal{V}_{\vec{p}}^{\dagger}(\mathfrak{X}) = 0$$

if  $\bar{p}_1 + \bar{p}_2 + \cdots + \bar{p}_N \neq 0$ . Therefore, assume  $\bar{p}_1 + \bar{p}_2 + \cdots + \bar{p}_N = 0$ . Choose  $p_j$ 's in such a way that

$$p_1 + p_2 + \cdots + p_N = 0,$$

and fix them in the following. For an element  $\langle \psi | \in \mathcal{V}_{\vec{p}}^{\dagger}(\mathfrak{X})$ , put

$$\psi_{l_1, l_2, \dots, l_N} = \langle \psi | (|l_1M + p_1\rangle \otimes |l_2M + p_2\rangle \otimes \cdots \otimes |l_NM + p_N\rangle).$$

If  $\psi_{l_1, l_2, \dots, l_N} \neq 0$ , then  $l_1 + l_2 + \cdots + l_N = 0$ . The condition (1\*) implies that  $\psi_{l_1, l_2, \dots, l_N}$  determines uniquely the values

$$\begin{aligned} \langle \psi | (a(-n_1^{(1)}) \cdots a(-n_{k_1}^{(1)}) |l_1M + p_1\rangle \otimes a(-n_1^{(2)}) \cdots a(-n_{k_2}^{(2)}) |l_2M + p_2\rangle \otimes \\ \cdots \otimes a(-n_1^{(N)}) \cdots a(-n_{k_N}^{(N)}) |l_NM + p_N\rangle), \end{aligned}$$

for any positive integers  $n_j^{(i)}$ . Also, the condition (2\*) implies that  $\psi_{l_1, l_2, \dots, l_N}$  can be uniquely determined by the value  $\psi_{0,0, \dots, 0}$ . Thus, we conclude that

$$\dim_{\mathbf{C}} \mathcal{V}_{\vec{p}}^{\dagger}(\mathfrak{X}) = 1$$

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This proves Proposition 2.1.

Let us consider a bigger subspace  $\mathcal{V}_{\vec{p}}^{\dagger}(n)$  of  $\mathcal{H}^{\dagger}(\vec{p})$ . An element  $\langle \psi |$  is in  $\mathcal{V}_{\vec{p}}^{\dagger}(n)$ , if  $\langle \psi |$  satisfies the following two conditions (1<sub>n</sub><sup>\*\*</sup>) and (2<sub>n</sub><sup>\*\*</sup>).

$$(1_n^{**}) \quad \sum_{j=1}^n \text{Res}_{\xi_j=0} (\langle \psi | a(\xi_j) | \phi \rangle g_j(\xi_j) d\xi_j) = 0$$

for all

$$(g_j(\xi_j) d\xi_j) \in \mathbf{C} \bigoplus_{j=1}^N (\mathbf{C}[\xi_j] \xi_j^{-n}), \quad \text{and} \quad |\phi\rangle \in \mathcal{H}(\vec{p}),$$

where an element  $c \in \mathbf{C}$  in the right hand side can be considered as  $(c, c, \dots, c)$ .

$$(2_n^{**}) \quad \sum_{j=1}^n \text{Res}_{\xi_j=0} (\langle \psi | V_{\pm M}(\xi_j) | \phi \rangle g_j(\xi_j) d\xi_j) = 0$$

for all

$$(h_j(\xi_j) (d\xi_j)^{\frac{M}{2}}) \in \bigoplus_{j=1}^N (\mathbf{C}[\xi_j] \xi_j^{-n}) (d\xi_j)^{(1-\frac{M}{2})}, \quad \text{and} \quad |\phi\rangle \in \mathcal{H}(\vec{p}).$$

**Key Lemma.** *Under the above notation we have*

$$\dim \mathcal{V}_{\vec{p}}^{\dagger}(n) < \infty.$$

To prove the Key Lemma we need the following Lemma due to Tuchiya.

**Lemma 2.3.** *Let  $\mathfrak{X}$  be an  $N$ -pointed smooth curve of genus  $g$  with formal neighbourhoods. Then we have*

$$\dim \mathcal{V}_{\vec{p}}^{\dagger}(\mathfrak{X}) \leq n^g.$$

The idea of the proof is as follows. For each non-zero element  $\langle \psi | \in \mathcal{V}_{\vec{p}}^{\dagger}(\mathfrak{X})$  and any element  $|v\rangle \in F_0 \mathcal{H}(\vec{p})$  we can define a meromorphic form

$$\langle \psi | V_{\pm M}(z_1) V_{\pm M}(z_2) \cdots V_{\pm M}(z_m) | v \rangle (dz_1)^{\frac{M}{2}} (dz_2)^{\frac{M}{2}} \cdots (dz_m)^{\frac{M}{2}}$$

on  $\underbrace{C \times C \times \cdots \times C}_m$ . By the operator product expansion of the energy momentum tensor  $T(z)$  we know singularities of this form and we can express the form by means of prime forms. This shows Lemma 2.3. To prove Key Lemma we need also the following lemma.

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**Lemma 2.4.** *For positive integers  $n$  and  $N$  there exist a smooth curve  $D$  of genus  $g$  and points  $Q_1, \dots, Q_N$  on  $D$  with local coordinates  $\xi_1, \xi_2, \dots, \xi_N$  such that*

$$\begin{aligned} Gr_{\bullet}^F H^0(D, \mathcal{O}_D(* \sum Q_j)) &\subset \mathbb{C} \bigoplus \bigoplus_{j=1}^N \mathbb{C}[\xi_j^{-1}] \xi_j^{-n} \\ Gr_{\bullet}^F H^0(D, \omega_D^{\otimes(1-\frac{M}{2})}(* \sum Q_j)) &\subset \mathbb{C} \bigoplus \bigoplus_{j=1}^N \mathbb{C}[\xi_j^{-1}] \xi_j^{-n} (d\xi_j)^{1-\frac{M}{2}} \end{aligned}$$

where the filtration  $F$  can be defined by the order of poles at  $Q_j$ .

The first inclusion can be proved, if the divisor  $n(Q_1 + Q_2 + \dots + Q_N)$  is not special on a curve  $D$ . The second inclusion is trivially true, if we have  $(2g-2)(1-\frac{M}{2}) > nN$ .

Now introducing the filtration on  $\mathcal{H}(\bar{p})$  and  $\mathcal{H}^\dagger(\bar{p})$  compatible with the filtration in Lemma 2.4, we can show *finite dimensionality* of  $\mathcal{V}_{\bar{p}}(\mathfrak{X})$  for all  $N$ -pointed stable curve with formal neighbourhoods.

Now let us consider a semi-stable curve  $C$ . For a double point  $P \in C$  we let  $\pi: \tilde{C} \rightarrow C$  be the normalization at the point  $P$ . Then, the inverse image  $\pi^{-1}(P)$  of the point  $P$  consists of two points  $P_+, P_-$ . Let  $\eta_+, \eta_-$  be formal coordinates of  $P_+$  and  $P_-$  respectively such that  $C$  is defined formally in a neighbourhood of the origin of  $\mathbb{C}^2$  by an equation  $\eta_+ \cdot \eta_- = 0$ . Let  $\mathfrak{X} = (C; Q_1, \dots, Q_N; \xi_1, \dots, \xi_N)$  be an  $N$ -pointed stable curve with formal neighbourhoods whose underlying curve is the semi-stable curve  $C$ . Put

$$\tilde{\mathfrak{X}} = (\tilde{C}; Q_1, \dots, Q_N, P_+, P_-; \xi_1, \dots, \xi_N, \eta_+, \eta_-).$$

Then, we have the following theorem.

**Theorem 2.5.** *Under the above notation and assumptions, we have a canonical isomorphism.*

$$\bigoplus_{\bar{q} \in \mathbb{Z}/M\mathbb{Z}} \mathcal{V}_{\bar{q}, -\bar{q}, \bar{p}}^\dagger(\tilde{\mathfrak{X}}) \simeq \mathcal{V}_{\bar{p}}^+(\mathfrak{X}).$$

From this theorem and Proposition 2.1 we infer the following lemma.

**Lemma 2.6.** *Let  $\mathfrak{X} = (C; Q_1, \dots, Q_N; \xi_1, \dots, \xi_N)$  be an  $N$ -pointed stable curve with formal neighbourhoods. Assume that all the irreducible component of the semi-stable curve  $C$  are  $\mathbb{P}^1(\mathbb{C})$  and the genus of  $C$  is  $g$ . Then, we have*

$$\dim_{\mathbb{C}} \mathcal{V}_{\bar{p}}^\dagger(\mathfrak{X}) = \begin{cases} M^g, & \text{if } \bar{p}_1 + \dots + \bar{p}_N = 0 \\ 0, & \text{otherwise.} \end{cases}$$

Now we need to show that  $\dim_{\mathbb{C}} \mathcal{V}_{\bar{p}}^+(\mathfrak{X})$  depends only on the genus of the underlying curve  $C$ . For that purpose we need to consider the family  $\mathcal{V}_{\bar{p}, N} = \bigcup_{\mathfrak{X}} \mathcal{V}_{\bar{p}}^+(\mathfrak{X})$

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over the moduli space  $\overline{\mathcal{M}}_{g,N}^{(\infty)}$  of  $N$ -pointed curves of genus  $g$  with formal neighbourhoods. By a similar method as the one in [2], we can show that  $\overline{\mathcal{V}}_{\vec{p},N}$  comes from a sheaf  $\mathcal{V}_{\vec{p},N}^{(1)}$  on  $\mathcal{M}_{g,N}^{(1)}$ , the moduli space of  $N$ -pointed curves of genus  $g$  with first order neighbourhoods. Then, by Key Lemma we can show that  $\mathcal{V}_{\vec{p},N}^{(1)}$  is a coherent  $\mathcal{O}_{\overline{\mathcal{M}}_{g,N}^{(1)}}$ -module and it carries a logarithmic projectively flat connection. From these fact we infer that  $\mathcal{V}_{\vec{p},N}^{(1)}$  is locally free on the open part of  $\overline{\mathcal{M}}_{g,N}^{(1)}$  corresponding to non-singular curves.

Again, using a similar arguments as in [2] we can show that  $\mathcal{V}_{\vec{p},N}^{(1)}$  is locally free. By Lemma 2.6 this implies our main theorem.

## REFERENCES

- [1] Kawamoto, N., Y. Namikawa, A. Tuchiya & Y. Yamada, *Geometric realization of conformal field theory on Riemann surfaces*, Commun. Math. Phys. **116** (1988), 247 – 308.
- [2] A.Tsuchiya, K. Ueno & Y. Yamada, *Conformal field theory on universal family of stable curves with gauge symmetries*, Adv. Stud. in Pure Math. **19** (1989), 459–566.