Div - curl estimates with critical power weights

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1 Introduction

This note is based on [11]. But, the proof of the main result in [11] is not correct, though the result is true. Aim of this note is to give a proof of it, along the talk. It applies Bogovskii formula instead of the Green function for the Neumann problem of Poisson equation, which was used in [11]. I learned this formula from Professor Hideo Kozono, and would like to thank him for teaching me it. It is should be emphasize that thanks to the formula, the result on 3-dimension in [11] is generalized to all dimension.

Div - curl estimate is a inequality of the form, which was firstly studied by Coifman-Linons-Meyer-Semmes [2]: for $p, q \in (n/(n+1), \infty)$ and 1/r = 1/p + 1/q < 1 + 1/n,

$$\|(u \cdot \nabla)v\|_{H^r} \lesssim \|u\|_{H^q} \|\nabla v\|_{H^p},$$
 (1)

where u and v are vector valued functions; $u = \{u_j\}_{j=1}^n, v = \{v_j\}_{j=1}^n$, and div $u = \nabla \cdot u = 1$. Here

$$(u \cdot \nabla)v = \left(\sum_{j=1}^n u_j \partial_j v_1, \cdots, \sum_{j=1}^n u_j \partial_j v_n\right).$$

 H^p , $(p \in (0, \infty))$ is the Hardy spaces: for any $\phi \in C_0^{\infty}(\mathbb{R}^n)$ with supp $\phi \subset\subset B(0, 1)$ and $\int \varphi dx = 0$,

$$||f||_{H^p} := ||M_{\phi}[f]||_{L^p}$$
, where $M_{\phi}[f](x) := \sup_{t>0} |f * \phi_t(x)|$,

where $g_t(x) := t^{-n}g(x/t)$. In the case $p = \infty$, define $H^{\infty} := L^{\infty}$.

As we mentioned in [11], we make use of a real interpolation spaces between weighted Hardy spaces.

Definition 1.1. Let $p, q \in (0, \infty]$ and $\alpha \in \mathbb{R}$. Define Hardy spaces associated with Herz spaces $H^{p,q}_{\alpha}(\mathbb{R}^n)$ as

$$H^{p,q}_{\alpha}(\mathbb{R}^n) := \left\{ f \in \mathcal{S}'; \|f\|_{H^{p,q}_{\alpha}} := \|M_{\phi}f\|_{L^{p,q}_{\alpha}} < \infty \right\},$$

where

$$||f||_{L^{p,q}_{\alpha}} := ||\{2^{\alpha k}||f||_{L^{p}(A_{k})}\}_{k \in \mathbb{Z}}||_{l^{q}}.$$

We explain notations. $\mathcal{S}(\mathbb{R}^n)$ and $\mathcal{S}'(\mathbb{R}^n)$ denote the Schwartz spaces of rapidly decreasing smooth functions and tempered distributions on \mathbb{R}^n , respectively. For a measurable subset $E \subset \mathbb{R}^n$, |E| and χ_E are the volume and the characteristic function of E, respectively. For any integers j, A_j denotes a annulus $\{x \in \mathbb{R}^n; 2^{j-1} \leq |x| < 2^j\}$, and χ_j is the characteristic function of A_j . B(x,r) is a ball in \mathbb{R}^n , centered at x of radius r. $\langle g \rangle_B := |B|^{-1} \int_B f dy$. Also, $A \lesssim B$ means $A \leq cB$ with positive constant c, and $A \approx B$ means $A \lesssim B$ and $B \lesssim A$.

To give $(u \cdot \nabla)v$ a definition as a tempered distribution, we define Y by a space of all locally integrable functions f satisfying that there exist $c_f > 0$ and a seminorm $|\cdot|_{\mathcal{S}}$ of \mathcal{S} so that $\int |f(x)\varphi(x)|dx \leq c_f|\varphi|_{\mathcal{S}}$, for all $\varphi \in \mathcal{S}$.

The main result reads as follows.

Theorem 1.1. For n/(n+1) , it holds

$$\|(u\cdot\nabla)v\|_{H^{p,\infty}_{\alpha(p)}} \le c\|u\|_{L^{\infty}}\|\nabla v\|_{H^{p,n/(n+1)}_{\alpha(p)}},$$

for $u \in L^{\infty}(\mathbb{R}^n)^n$ with div u = 0 and $v \in (Y \cap W_{loc}^{1,r}(\mathbb{R}^n))^n$ for some $r \in (1, \infty)$, where $\alpha(p) := n(1 - 1/p) + 1$.

Remark 1.1. The same argument as the proof of Theorem 1.1, we can also show a weak type estimate:

$$\|(u \cdot \nabla)v\|_{H^{(n/(n+1),\infty)}} \lesssim \|u\|_{L^{\infty}} \|\nabla v\|_{H^{n/(n+1)}},$$
 (2)

because $\alpha(n/(n+1)) = 0$. Here, $f \in H^{(p,q)} \iff M_{\phi}[f] \in L^{(p,q)}$, where $L^{(p,q)}$ is the Lorentz spaces. Similar estimates were established by Miyakawa [8]. This can be regarded as an endpoint case with p = n/(n+1) and $q = \infty$ of [2]. The ingredient of the proof of (2) is the pointwise estimate:

$$||N[v]||_{L^{(n/(n+1),\infty)}} \lesssim ||\nabla v||_{H^{n/(n+1)}}$$

instead of (4). This is achieved from the pointwise estimate (7) and a Fefferman-Stein's vector valued inequality (2) of Theorem 1 in [3].

Motivation of this research comes from the optimal L^2 -energy decay for the incompressible Navier-Stokes equations. Wiegner [12] constructed global weak solutions u having

$$||u(t)||_{L^2} \lesssim t^{-(n+2)/4},$$

assuming that initial data $a \in L^2$ satisfying $||e^{t\Delta}a||_{L^2} \lesssim t^{-(n+2)/4}$. By Miyakawa-Schonbek [9], it is well-known that the decay order (n+2)/4 is optimal. $H^p_{\alpha(p)}$ is relevant to this order (n+2)/4, because one has that for $p \in (0,2]$,

$$||e^{t\Delta}a||_{L^2} \lesssim t^{-(n+2)/4} ||a||_{H^p_{\alpha(p)}},$$

see [10] for the proof. The present author [10] investigated the L^2 decay of mild solutions by Kato [5] and constructed solutions whose decay order of L^2 energy is $\gamma < (n+2)/4$. One of reasons why the order γ in [10] did not reach to the optimal order (n+2)/4 is that div-curl estimate in [10] cannot allow us to deal with the critical exponent $\alpha = \alpha(p)$. As mentioned in Remark 7.3 in [7], the bilinear term $(u \cdot \nabla)v$ does not belong to $H^p_{\alpha(p)}$. This observation tells us that if we try to establish div-curl estimate with $\alpha = \alpha(p)$, we has to replace $H^p_{\alpha(p)}$ in the left hand side by some larger spaces. For the purpose, we use Hardy spaces associated to Herz spaces, as in [6] and [7]. Although, a critical div - curl lemma is proved in this article, the author does not know whether or not it is possible to construct global solutions having optimal L^2 decay from the similar argument as the previous paper [10].

2 Proof

Before we start the proof of Theorem 1.1, we point out the mistake in [?]. In [?], the Green function for the Neumann problem of Poisson equation:

$$-\Delta h = g$$
 in B , $g = 0$ on ∂B

for $g \in C_0^{\infty}(B)$ with $\int g dx = 0$, is applied. I treated the solution h as a $C_0^2(B)$ function in [11]. Although $h \in C^2(\mathbb{R}^n)$, this is not true in general. To overcome this difficulty, as we mentioned above, we apply Bogovskii formula. This is a representation, with a kernel function, of solutions to the divergence equation.

2.1 Bogovskii formula

Let B be a ball in \mathbb{R}^n and $g \in C_0^{\infty}(B)$ with $\int g dx = 0$. We refer Lemma III.3.1 in [4] for next lemma.

Lemma 2.1. There exists a vector function $\mathbf{K} = \{K_j\}_{j=1}^n$ on $\mathbb{R}^n \times \mathbb{R}^n \setminus \{(x,y) : x = y\}$ so that $\mathbf{G}_B(x) := \int \mathbf{K}(x,y)g(y)dy \in C_0^{\infty}(B)^n$ is a solution to the divergence equation $\nabla \cdot \mathbf{G}_B = g$ on B satisfying that for $q \in (1,\infty)$

$$\|\mathbf{G}_B\|_{L^q} \lesssim |B|^{1/n} \|g\|_{L^q}$$
 and $\|\nabla \cdot \mathbf{G}_B\|_{BMO} \lesssim \|g\|_{L^{\infty}}$.

Remark 2.1. The $L^{\infty}-BMO$ estimate above is deduced from the fact that the operator $g \mapsto \mathbf{G}_B$ is a Calderón-Zygmund operator.

2.2 Vector valued restricted weak type inequality

Another ingredient for the proof of Theorem 1.1 is a vector-valued "restricted weak" type inequality for Hardy-Littlewood maximal operator; for $r \in (0, \infty)$

$$M_r f(x) := \sup_{B \ni x} \langle |f|^r \rangle_B^{1/r}, ,$$

where the supremum is taken over all ball B containing x. Define $Mf(x) := M_1f(x)$. The following is a generalization of the result of Fefferman-Stein [3], and the proof is found in [11].

Proposition 2.1. For $1 < r, p < \infty$ and $\alpha = n(1 - 1/p)$,

$$\left\| \left(\sum_{l=1}^{\infty} \left(M f_l \right)^r \right)^{1/r} \right\|_{L^{p,\infty}_{\alpha}} \lesssim \left\| \left(\sum_{l=1}^{\infty} |f_l|^r \right)^{1/r} \right\|_{L^{p,1}_{\alpha}}.$$

This can be rewritten as the following form.

Corollary 2.1. For $0 < r < 1 < p < \infty \ and \ \alpha = n(1/r - 1/p)$,

$$\left\| \sum_{l=1}^{\infty} M_r f_l \right\|_{L^{p,\infty}_{\alpha}} \lesssim \left\| \sum_{l=1}^{\infty} |f_l| \right\|_{L^{p,r}_{\alpha}}.$$

2.3 Complete of the proof of Theorem 1.1

The proof is almost same as that in [11] except for applying Lemma 2.1 instead of the Green function for the Neumann problem of Poisson equations.

Because

$$\|(u \cdot \nabla)v\|_{H^{p,\infty}_{\alpha(p)}} = \sum_{k=1}^n \left\| \sum_{j=1}^n u_j \partial_j v_k \right\|_{H^{p,\infty}_{\alpha(p)}} = \sum_{k=1}^n \left\| M_\phi \left(\sum_{j=1}^n u_j \partial_j v_k \right) \right\|_{L^{p,\infty}_{\alpha(p)}},$$

it is enough to show the inequality

$$\left\| M_{\phi} \left(\sum_{j=1}^{n} u_{j} \partial_{j} v \right) \right\|_{L_{\alpha(p)}^{p,\infty}} \lesssim \|u\|_{L^{\infty}} \|\nabla v\|_{H_{\alpha(p)}^{p,n/(n+1)}},$$

for all divergence free vector fields u and functions $v \in Y \cap W_{loc}^{1,r}$. Firstly, we give a definition of $\sum_{j=1}^{n} u_j \partial_j v$ as a tempered distribution as follows; for $\varphi \in \mathcal{S}$

$$\left\langle \sum_{j=1}^n u_j \partial_j v, \varphi \right\rangle := -\sum_{j=1}^n \int u_j(y) v(y) \partial_j \varphi(y) dy.$$

Our assumption ensures that the integral in the right hand side absolutely converges. Then, it follows

$$\sum_{j=1}^{n} u_j \partial_j v * \phi_t(x) = -C_\phi \|u\|_{L^\infty} \int v(y) \left[\sum_{j=1}^{n} \tilde{u}_j(y) \partial_{y_j} \phi_t(x-y) \right] dy,$$

where C_{ϕ} is a constant depending on ϕ , and $\tilde{u}_{j}(y) = \frac{u_{j}(y)}{C_{\phi}||u||_{L^{\infty}}}$. Owing to the divergence free condition on u, we see that for every $x \in \mathbb{R}^{n}$

$$\sum_{j=1}^{n} \tilde{u}_{j}(y) \partial_{y_{j}} \phi_{t}(x-y) = \sum_{j=1}^{n} \partial_{y_{j}} \left(\tilde{u}_{j}(y) \phi_{t}(x-y) \right) \quad \text{in} \quad \mathcal{S}'(\mathbb{R}_{y}^{n}). \tag{3}$$

Hence, we obtain the pointwise estimate

$$M_{\phi}\left(\sum_{j=1}^{n} u_{j} \partial_{j} v\right)(x) \leq C_{\phi} \|u\|_{L^{\infty}} N[v](x),$$

where

$$N[v](x) := \sup_{t>0} \left| \int v(y)g(y)dy \right| \text{ and } g(y) = g(y;x,t) := \sum_{j=1}^{n} \tilde{u}_{j}(y)\partial_{y_{j}}\phi_{t}(x-y).$$

It is enough to prove that

$$||N[v]||_{L^{p,\infty}_{\alpha(p)}} \lesssim ||\nabla v||_{H^{p,n/(n+1)}_{\alpha(p)}}.$$
(4)

To show this from a pointwise estimate, we make use of Bogovskii formula Lemma 2.1 and the atomic decomposition in $H_{\alpha(p)}^{p,n/(n+1)}$ due to Miyachi [7].

Fix $x \in \mathbb{R}^n$ and $t \in (0, \infty)$. There exists $\varepsilon_0 > 0$ so that for all $\varepsilon \in (0, \varepsilon_0)$, supp $g_*\phi_\varepsilon \subset\subset B(x,t)$. Take $\eta_0 \in C_0^\infty(\mathbb{R}^n)$ such that $\eta(y) = \eta(y;x,t) := \eta_0\left(\frac{y-x}{t}\right)$ satisfies that $0 \le \eta \le 1$ and for $\varepsilon \in (0, \varepsilon_0)$

supp
$$g * \phi_{\varepsilon} \subset\subset \text{supp } \eta \subset\subset B(x,t)$$
, and $\eta \equiv 1$ on supp $g * \phi_{\varepsilon}$

Remark that $\|\eta\|_{L^p} = ct^{n/p}$ for all $p \in [1, \infty]$ with c independent of x, t and ε . Next we see that

 $\sigma_{\varepsilon} := \|\eta\|_{L^1}^{-1} \int g * \phi_{\varepsilon} dy \to 0 \text{ as } \varepsilon \searrow 0.$

In fact, for a test function $\rho \in C_0^{\infty}(B(x,2t))$ with $\rho \equiv 1$ on B(x,t), we have from (3)

$$\|\eta\|_{L^1}\sigma_{\varepsilon} = \langle g, \rho * \phi_{\varepsilon} \rangle \to \langle g, \rho \rangle = 0.$$

For $\varepsilon \in (0, \varepsilon_0)$, letting $g^{\varepsilon} := g * \phi_{\varepsilon} - \sigma_{\varepsilon} \eta \in C_0^{\infty}(B(x, t))$, we obtain $\int g^{\varepsilon} dy = 0$. From Lemma 2.1,

$$\mathbf{G}^{\varepsilon}(y) := \int \mathbf{K}(y, z) g^{\varepsilon}(z) dz \in C_0^{\infty}(B(x, t))$$

solves the Drichlet problem of Poisson equation:

$$\nabla \cdot \mathbf{G}^{\varepsilon} = g^{\varepsilon} \text{ in } B(x,t), \quad \mathbf{G}^{\varepsilon} = 0 \text{ on } \partial B(x,t).$$

Further, \mathbf{G}^{ε} fulfills the following estimates: for all $q \in (1, \infty)$,

$$\|\mathbf{G}^{\varepsilon}\|_{L^{q}} \lesssim t^{-n+n/q} \quad \text{and} \quad \|\partial_{j}\mathbf{G}^{\varepsilon}\|_{BMO} \lesssim t^{-(n+1)}.$$
 (5)

Indeed, these follows that $\|g^{\varepsilon}\|_{L^{q}} \lesssim t^{-1-n+n/q}$ and $\|g^{\varepsilon}\|_{L^{\infty}} \lesssim t^{-(n+1)}$, respectively. Integration by parts yields

$$\int vgdy = -\lim_{\varepsilon \to 0} \int \nabla v \cdot \mathbf{G}^{\varepsilon} dy.$$

Therefore, one obtains

$$\left| \int vgdy \right| \le \sum_{k=1}^n \limsup_{0 < \varepsilon < \varepsilon_0} \left| \int \partial_k v \mathbf{G}_k^{\varepsilon} dy \right|.$$

Since $\partial_k v \in H^{p,n/(n+1)}_{\alpha(p)}$, following Miyachi [7], it can be decomposed as

$$\partial_k v = \sum_{j=1}^{\infty} a_j^{(k)}$$

where supp $a_j^{(k)} \subset B_j = B(x_j, r_j), \ a_j^{(k)} \in L^{\infty} \text{ and } \int x^{\beta} a_j^{(k)}(x) dx = 0 \text{ for } |\beta| \leq 1, \text{ also }$

$$\left\| \sum_{j=1}^{\infty} \|a_j^{(k)}\|_{L^{\infty}} \chi_{B_j} \right\|_{L^{p,n/(n+1)}_{\alpha(p)}} \lesssim \|\partial_k v\|_{H^{p,n/(n+1)}_{\alpha(p)}}.$$

and one obtains

$$\left| \int vgdy \right| \le \sum_{k=1}^{n} \sum_{j=1}^{\infty} \limsup_{0 < \varepsilon < \varepsilon_0} \left| \int a_j^{(k)} \mathbf{G}_k^{\varepsilon} dy \right|.$$

From (5), we immediately see that

$$\left| \int a_j^{(k)} \mathbf{G}_k^{\varepsilon} dy \right| \le \|a_j^{(k)}\|_{L^{\infty}} |B(x,t)|^{1-1/q} \|\mathbf{G}_k^{\varepsilon}\|_{L^q} \lesssim \|a_j^{(k)}\|_{L^{\infty}}.$$

When $x \notin 4B_j$, if $Ct < |x - x_j|$ with C > 8/3, then it holds $B_j \cap B(x,t) = \emptyset$ and $\int a_j^{(k)} \mathbf{G}_k^{\varepsilon} dy = 0$. On the other hand, if $Ct \ge |x - x_j|$, then we can derive the decay estimate

$$\lim_{0 < \varepsilon < \varepsilon_0} \left| \int a_j^{(k)} \mathbf{G}_k^{\varepsilon} dy \right| \lesssim \|a_j^{(k)}\|_{L^{\infty}} \left(\frac{r_j}{|x - x_j|} \right)^{n+1}. \tag{6}$$

We may assume $x \neq x_j$. Using the moment condition on $a_i^{(k)}$ twice, one has

$$\int a_j^{(k)}(y)\mathbf{G}_k^{\varepsilon}(y)dy = \int a_j^{(k)}(y)\left(\mathbf{G}_k^{\varepsilon}(y) - \mathbf{G}_k^{\varepsilon}(x_j)\right)dy$$

$$= \sum_{s=1}^n \int_0^1 \int a_j^{(k)}(y)(y - x_j)_s(\partial_s \mathbf{G}_k^{\varepsilon})(\theta y + (1 - \theta)x_j)dyd\theta$$

$$= \sum_{s=1}^n \int_0^1 \int a_j^{(k)}(y)(y - x_j)_s \left[(\partial_s \mathbf{G}_k^{\varepsilon})(\theta y + (1 - \theta)x_j) - \langle \partial_s \mathbf{G}_k^{\varepsilon} \rangle_{B(x_j, \theta r_j)}\right]dyd\theta.$$

From this representation, the decay estimate (6) is derived as follows from (5);

$$\left| \int a_j^{(k)}(y) \mathbf{G}_k^{\varepsilon}(y) dy \right| \lesssim r_j \|a_j^{(k)}\|_{L^{\infty}} \sum_{s=1}^n \int_0^1 \theta^{-n} \int_{B(x_j, \theta r_j)} \left| \partial_s \mathbf{G}_k^{\varepsilon}(y) - \langle \partial_s \mathbf{G}_k^{\varepsilon} \rangle_{B(x_j, \theta r_j)} \right| dy d\theta$$

$$\lesssim r_j^{n+1} \|a_j^{(k)}\|_{L^{\infty}} \sum_{s=1}^n \|\partial_s \mathbf{G}_k^{\varepsilon}\|_{BMO}$$

$$\lesssim \left(\frac{r_j}{t}\right)^{n+1} \|a_j^{(k)}\|_{L^{\infty}}$$

$$\lesssim \left(\frac{r_j}{|x - x_j|}\right)^{n+1} \|a_j^{(k)}\|_{L^{\infty}}.$$

As mentioned in [7], because $\left(\frac{1}{1+|x-x_j|/r_j}\right)^{n+1} \approx M_{n/(n+1)}(\chi_{B_j})(x)$, as a consequence it follows that for all $x \in \mathbb{R}^n$,

$$N[v](x) = \sup_{t>0} \left| \int v(y)g(y;x,t)dy \right| \lesssim \sum_{k=1}^{n} \sum_{j=1}^{\infty} \|a_j^{(k)}\|_{L^{\infty}} M_{n/(n+1)}(\chi_{B_j})(x).$$
 (7)

Now, we apply Corollary 2.1 with r = n/(n+1) and obtain

$$||N[v]||_{L^{p,\infty}_{\alpha(p)}} \lesssim \sum_{k=1}^{n} \left\| \sum_{j=1}^{\infty} ||a_{j}^{(k)}||_{L^{\infty}\chi_{B_{j}}} \right\|_{L^{p,n/(n+1)}_{\alpha(p)}} \lesssim \sum_{k=1}^{n} ||\partial_{k}v||_{H^{p,n/(n+1)}_{\alpha(p)}} = ||\nabla v||_{H^{p,n/(n+1)}_{\alpha(p)}}.$$

Here we have used $n(1-1/p) + n((n+1)/n - 1) = n(1-1/p) + 1 = \alpha(p)$. The proof is completed.

Remark 2.2. In [1], the pointwise estimate (6) with $n+1-\varepsilon$ replaced by n+1 was proved.

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