# Stochastic differential equations related to random matrix theory

By

Hirofumi Osada\* and Hideki Tanemura\*\*

#### Abstract

In this note we review recent results on existence and uniqueness of solutions of infinite-dimensional stochastic differential equations describing interacting Brownian motions on  $\mathbb{R}^d$ .

# § 1. Introduction

Let  $\mathbf{X}^{N}(t) = (X_{j}^{N}(t))_{j=1}^{N}$  be a solution of the stochastic differential equation (SDE)

(1.1) 
$$dX_j^N(t) = dB_j(t) + \frac{\beta}{2} \sum_{k=1, k \neq j}^N \frac{dt}{X_j^N(t) - X_k^N(t)}$$

or the SDE with Ornstein-Uhlenbeck's type drifts

(1.2) 
$$dX_j^N(t) = dB_j(t) - \frac{\beta}{4N} X_j^N(t) dt + \frac{\beta}{2} \sum_{k=1, k \neq j}^N \frac{dt}{X_j^N(t) - X_k^N(t)},$$

where  $B_j(t), j = 1, 2, ..., N$  are independent one-dimensional Brownian motions. These are called Dyson's Brownian motion models with parameters  $\beta > 0$  [4]. They were introduced to understand the statistics of eigenvalues of random matrix ensembles as

Received January 31, 2016. Revised May 14, 2016.

 $<sup>2010\ \</sup>mathrm{Mathematics}\ \mathrm{Subject}\ \mathrm{Classification}(s);\ 15B52,\ 30C15,\ 47D07,\ 60G55,\ 82C22$ 

Key Words: Random matrices, Interacting particle systems, Stochastic differential equations.

Supported by a Grant-in-Aid for Scientific Research (KIBAN-A, No.24244010) and a Grant-in-Aid for Scientific Research (KIBAN-C, No.15K04916)

<sup>\*</sup>Faculty of Mathematics, Kyushu University, Fukuoka 819-0395 Japan. e-mail: osada@math.kyushu.ac.jp

<sup>\*\*</sup>Department of Mathematics and Informatics, Faculty of Science, Chiba University, 1-33 Yayoi-cho, Inage-ku, Chiba 263-8522, Japan.

e-mail: tanemura@math.s.chiba-u.ac.jp

<sup>© 2016</sup> Research Institute for Mathematical Sciences, Kyoto University. All rights reserved.

distributions of particle positions in one-dimensional Coulomb gas systems with logpotential.

The solution of (1.2) is a natural reversible stochastic dynamics with respect to  $\check{\mu}^N_{\mathsf{bulk},\beta}$ :

(1.3) 
$$\check{\mu}_{\mathsf{bulk},\beta}^{N}(d\mathbf{x}_{N}) = \frac{1}{Z} h_{N}(\mathbf{x}_{N})^{\beta} e^{-\frac{\beta}{4N}|\mathbf{x}_{N}|^{2}} d\mathbf{x}_{N},$$

where  $d\mathbf{x}_N = dx_1 dx_2 \cdots dx_N$ ,  $\mathbf{x}_N = (x_i) \in \mathbb{R}^N$ , and

$$h_N(\mathbf{x}_N) = \prod_{i < j}^N |x_i - x_j|.$$

Throughout, Z denotes a normalizing constant. Gaussian ensembles are called Gaussian orthogonal/unitary/symplectic ensembles (GOE/GUE/GSE) according to their invariance under conjugation by orthogonal/unitary/symplectic groups, which correspond to the inverse temperatures  $\beta = 1, 2$  and 4, respectively [9, 2]. It is natural to believe that the N-limit of the process  $\mathbf{X}^N(t)$  solves the infinite-dimensional stochastic differential equation (ISDE)

(1.4) 
$$dX_j(t) = dB_j(t) + \frac{\beta}{2} \lim_{\substack{r \to \infty \\ |X_k(t)| < r}} \sum_{\substack{k=1, k \neq j \\ |X_k(t)| < r}}^{\infty} \frac{dt}{X_j(t) - X_k(t)}.$$

The result was not proved rigorously until a few years ago when it was shown for  $\beta = 2$  in [17], for  $\beta = 1, 2, 4$  in [8], and for  $\beta \ge 1$  in [23].

Set  $Y_j^N(t) = N^{1/6}(X_j^N(t) - 2\sqrt{N}), j = 1, 2, ..., N$  for the solution  $\mathbf{X}^N$  of (1.2). It has also been shown that the N-limit of the process  $\mathbf{Y}^N(t)$  solves the ISDE

(1.5) 
$$dY_j(t) = dB_j(t) + \frac{\beta}{2} \lim_{r \to \infty} \left\{ \sum_{\substack{k=1, k \neq j \\ |Y_k(t)| < r}}^{\infty} \frac{1}{Y_j(t) - Y_k(t)} - \int_{-r}^{r} \frac{\widehat{\rho}(x)dx}{-x} \right\} dt,$$

with  $\widehat{\rho}(x) = \pi^{-1} \sqrt{-x} \mathbf{1}(x < 0)$ , for  $\beta = 2$  [17] and for  $\beta = 1, 2, 4$  [8].

One of the key parts of proving the above results is the existence and uniqueness of solutions of an ISDE of the form

(1.6) 
$$dX_j(t) = dB_j(t) - \frac{1}{2} \nabla \Phi(X_j(t)) dt - \frac{1}{2} \sum_{k=1, k \neq j}^{\infty} \nabla \Psi(X_j(t), X_k(t)) dt$$

with free potential  $\Phi$  and interaction (pair) potential  $\Psi$ . In ISDEs (1.4) and (1.5),  $\Psi$  is given by the log pair potential  $-\beta \log |x-y|$ . The present note is a short summary of results on existence and uniqueness of solutions for ISDE (1.6).

### § 2. Quasi-Gibbs measure

Let S be a closed set in  $\mathbb{R}^d$  such that  $0 \in S$  and  $\overline{S^{\text{int}}} = S$ , where  $S^{\text{int}}$  denotes the interior of S. The configuration space  $\mathfrak{M}$  of unlabelled particles is given by

(2.1) 
$$\mathfrak{M} = \left\{ \xi : \xi \text{ is a nonnegative integer valued Radon measure in } S \right\}$$
$$= \left\{ \xi(\cdot) = \sum_{j \in \mathbb{I}} \delta_{x_j}(\cdot) : \sharp \{ j \in \mathbb{I} : x_j \in K \} < \infty, \text{ for any } K \text{ compact} \right\},$$

where  $\mathbb{I}$  is a countable set and  $\delta_a$  is the Dirac measure at  $a \in S$ . Thus  $\mathfrak{M}$  is a Polish space with the vague topology. We also introduce a subset  $\mathfrak{M}_{s,i}$  of  $\mathfrak{M}$ :

$$\mathfrak{M}_{s,i} = \{ \xi \in \mathfrak{M} : \xi(\{x\}) \le 1 \text{ for all } x \in S, \, \xi(S) = \infty \},$$

that is, the set of configurations of an infinite number of particles without collisions. For Borel measurable functions  $\Phi: S \to \mathbb{R} \cup \{\infty\}$  and  $\Psi: S \times S \to \mathbb{R} \cup \{\infty\}$  and a given increasing sequence  $\{b_r\}$  of  $\mathbb{N}$ , we introduce the Hamiltonian

(2.3) 
$$H_r(\xi) = H_r^{\Phi,\Psi}(\xi) = \sum_{x_j \in S_r} \Phi(x_j) + \sum_{x_j, x_k \in S_r, j < k} \Psi(x_j, x_k), \quad \xi = \sum_{j \in \mathbb{I}} \delta_{x_j},$$

where  $S_r = \{x \in S : |x| < b_r\}$ . We call  $\Phi$  a free potential, and call  $\Psi$  an interaction potential. Let  $\Lambda_r^m$  be the restriction of a Poisson random measure with intensity measure dx on  $\mathfrak{M}_r^m = \{\xi \in \mathfrak{M} : \xi(S_r) = m\}$ . We define maps  $\pi_r, \pi_r^c : \mathfrak{M} \to \mathfrak{M}$  such that  $\pi_r(\xi) = \xi(\cdot \cap S_r)$  and  $\pi_r^c(\xi) = \xi(\cdot \cap S_r^c)$ . For two measures  $\nu_1, \nu_2$  on a measurable space  $(\Omega, \mathcal{F})$  we write  $\nu_1 \leq \nu_2$  if  $\nu_1(A) \leq \nu_2(A)$  for any  $A \in \mathcal{F}$ . We can now state the definition of a quasi-Gibbs measure [13, 14].

**Definition 2.1.** A probability measure  $\mu$  on  $\mathfrak{M}$  is said to be a  $(\Phi, \Psi)$ -quasi Gibbs measure if its regular conditional probabilities

$$\mu_{r,\xi}^m(d\zeta) = \mu(d\zeta|\pi_r^c(\zeta) = \pi_r^c(\xi), \zeta(S_r) = m), \quad r, m \in \mathbb{N},$$

satisfy that, for  $\mu$ -a.s.  $\xi$ ,

$$c^{-1}e^{-H_r(\eta)}\Lambda_r^m(\pi_{S_r} \in d\eta) \le \mu_{r,\xi}^m(\pi_{S_r} \in d\eta) \le ce^{-H_r(\eta)}\Lambda_r^m(\pi_{S_r} \in d\eta).$$

Here,  $c = c(r, m, \xi)$  is a positive constant depending on r, m, and  $\xi$ .

It is readily seen that the quasi-Gibbs property is a generalized notion of the canonical Gibbs property. If  $\mu$  is a  $(\Phi, \Psi)$ -quasi Gibbs measure, then  $\mu$  is also a  $(\Phi + \Phi_{\text{loc.bdd}}, \Psi)$ -quasi Gibbs measure for any locally bounded measurable function  $\Phi_{\text{loc.bdd}}$ . In this sense, the notion of "quasi-Gibbs" seems to be robust. Information about the free potential of  $\mu$  is determined from its logarithmic derivative [12].

A function f on  $\mathfrak{M}$  is called a polynomial function if

$$(2.4) f(\xi) = Q(\langle \phi_1, \xi \rangle, \langle \phi_2, \xi \rangle, \dots, \langle \phi_{\ell}, \xi \rangle)$$

with  $\phi_k \in C_c^{\infty}(S)$  and a polynomial function Q on  $\mathbb{R}^{\ell}$ , where  $\langle \phi, \xi \rangle = \int_S \phi(x) \xi(dx)$  and  $C_c^{\infty}(S)$  is the set of smooth functions with compact support. We denote by  $\mathcal{P}$  the set of all polynomial functions on  $\mathfrak{M}$ .

**Definition 2.2.** We call  $\mathbf{d}^{\mu} \in L^1_{loc}(S \times \mathfrak{M}, \mu^{[1]})$  the logarithmic derivative of  $\mu$  if

$$\int_{S\times\mathfrak{M}} \mathbf{d}^{\mu}(x,\eta) f(x,\eta) d\mu^{[1]}(x,\eta) = -\int_{S\times\mathfrak{M}} \nabla_x f(x,\eta) d\mu^{[1]}(x,\eta)$$

is satisfied for  $f \in C_c^{\infty}(S) \otimes \mathcal{P}$ . Here  $\mu^{[k]}$  is the Campbell measure of  $\mu$ 

$$\mu^{[k]}(A \times B) = \int_A \mu_{\mathbf{x}}(B) \rho^k(\mathbf{x}) d\mathbf{x}, \quad A \in \mathcal{B}(S^k), B \in \mathcal{B}(\mathfrak{M}),$$

 $\mu_{\mathbf{x}}$  is the reduced Palm measure conditioned at  $\mathbf{x} \in S^k$ 

(2.5) 
$$\mu_{\mathbf{x}} = \mu \left( \cdot - \sum_{j=1}^{k} \delta_{x_j} \middle| \xi(x_j) \ge 1 \text{ for } j = 1, 2, \dots, k \right),$$

and  $\rho^k$  is the k-correlation function for  $k \in \mathbb{N}$ .

Quasi-Gibbs measures inherit the following property from canonical Gibbs measures [19, Lemma 11.2]. Let  $\mathcal{T}(\mathfrak{M})$  be the tail  $\sigma$ -field

$$\mathcal{T}(\mathfrak{M}) = \bigcap_{r=1}^{\infty} \sigma(\pi_r^c)$$

and let  $\mu_{\text{Tail}}^{\xi}$  be the regular conditional probability defined as

(2.6) 
$$\mu_{\text{Tail}}^{\xi} = \mu(\cdot | \mathcal{T}(\mathfrak{M}))(\xi).$$

Then the following decomposition holds:

(2.7) 
$$\mu(\cdot) = \int_{\mathfrak{M}} \mu_{\text{Tail}}^{\xi}(\cdot) \mu(d\xi).$$

Furthermore, there exists a subset  $\mathfrak{M}_0$  of  $\mathfrak{M}$  satisfying  $\mu(\mathfrak{M}_0) = 1$  and, for all  $\xi, \eta \in \mathfrak{M}_0$ :

(2.8) 
$$\mu_{\text{Tail}}^{\xi}(A) \in \{0, 1\} \text{ for all } A \in \mathcal{T}(\mathfrak{M}),$$

(2.9) 
$$\mu_{\text{Tail}}^{\xi}(\{\zeta \in \mathfrak{M} : \mu_{\text{Tail}}^{\xi} = \mu_{\text{Tail}}^{\zeta}\}) = 1,$$

(2.10) 
$$\mu_{\text{Tail}}^{\xi}$$
 and  $\mu_{\text{Tail}}^{\eta}$  are mutually singular on  $\mathcal{T}(\mathfrak{M})$  if  $\mu_{\text{Tail}}^{\xi} \neq \mu_{\text{Tail}}^{\eta}$ .

# § 3. General theory of solutions of ISDEs

A polynomial function f on  $\mathfrak{M}$  is a local function, that is, a function satisfying  $f(\xi) = f(\pi_r(\xi))$  for some  $r \in \mathbb{N}$ . When  $\xi \in \mathfrak{M}_r^m$ ,  $m \in \mathbb{N} \cup \{0\}$  and  $\pi_r(\xi)$  is represented by  $\sum_{j=1}^m \delta_{x_j}$ , we can regard  $f(\xi) = f(\sum_{j=1}^m \delta_{x_j})$  as a permutation invariant smooth function on  $S_r^m$ . For  $f, g \in \mathcal{P}$ , define

$$\mathbb{D}(f,g)(\xi) = \frac{1}{2} \sum_{j=1}^{\infty} \nabla_{x_j} f(\xi) \cdot \nabla_{x_j} g(\xi).$$

For a probability  $\mu$  on  $\mathfrak{M}$ , we denote by  $L^2(\mathfrak{M}, \mu)$  the space of square integrable functions on  $\mathfrak{M}$  with the inner product  $\langle \cdot, \cdot \rangle_{\mu}$  and the norm  $\| \cdot \|_{L^2(\mathfrak{M}, \mu)}$ . We consider the bilinear form  $(\mathcal{E}^{\mu}, \mathcal{P}^{\mu})$  on  $L^2(\mathfrak{M}, \mu)$  defined by

(3.1) 
$$\mathcal{E}^{\mu}(f,g) = \int_{\mathfrak{M}} \mathbb{D}(f,g) d\mu, \quad \mathcal{P}^{\mu} = \{ f \in \mathcal{P} : ||f||_{1}^{2} < \infty \},$$

where  $||f||_1^2 \equiv \mathcal{E}^{\mu}(f, f) + ||f||_{L^2(\mathfrak{M}, \mu)}^2$ .

We make the following assumptions

- (A.0)  $\mu$  has a locally bounded n-correlation function  $\rho^n$  for each  $n \in \mathbb{N}$ .
- (A.1) There exist upper semi-continuous functions  $\Phi_0: S \to \mathbb{R} \cup \{\infty\}$  and  $\Psi_0: S \times S \to \mathbb{R} \cup \{\infty\}$  that are locally bounded from below, and c > 0 such that

$$c^{-1}\Phi_0(x) \le \Phi(x) \le c\Phi_0(x), \quad c^{-1}\Psi_0(x,y) \le \Psi(x,y) \le c\Psi_0(x,y).$$

(A.2) There exists a T > 0 such that for each R > 0

$$\liminf_{r \to \infty} \operatorname{Erf}\left(\frac{r}{(r+R)T}\right) \int_{|x| \le r+R} \rho^{1}(x) dx = 0,$$

where  $\text{Erf}(t) = (2\pi)^{-1/2} \int_t^{\infty} e^{-x^2/2} dx$ .

Note that  $\mathcal{P}^{\mu} = \mathcal{P}$  and  $(\mathcal{E}^{\mu}, \mathcal{P}^{\mu}) = (\mathcal{E}, \mathcal{P})$  under condition (A.0).

**Theorem 3.1** ([12, 13, 14, 11, 16]). Suppose that  $\mu$  is a  $(\Phi, \Psi)$ -quasi Gibbs measure satisfying (A.0) and (A.1). Then

- (i)  $(\mathcal{E}, \mathcal{P})$  is closable and its closure  $(\mathcal{E}^{\mu}, \mathcal{D}^{\mu})$  is a quasi regular Dirichlet form and there exists the diffusion process  $(\Xi(t), P^{\xi}_{\mu})$  associated with  $(\mathcal{E}^{\mu}, \mathcal{D}^{\mu})$ .
- (ii) Furthermore, assume conditions (A.2) and (A.3):
- (A.3)  $\operatorname{Cap}^{\mu}((\mathfrak{M}_{s.i})^c) = 0$  and  $\operatorname{Cap}^{\mu}(\xi(\partial S) \ge 1) = 0$ ,

where  $\operatorname{Cap}^{\mu}$  is the capacity of the Dirichlet form. If there exists a logarithmic derivative  $\mathbf{d}^{\mu}$ , then there exists  $\tilde{\mathfrak{M}} \subset \mathfrak{M}$  such that  $\mu(\tilde{\mathfrak{M}}) = 1$ , and for any  $\xi = \sum_{j \in \mathbb{N}} \delta_{x_j} \in \tilde{\mathfrak{M}}$ , there exists an  $S^{\mathbb{N}}$ -valued continuous process  $\mathbf{X}(t) = (X_j(t))_{j=1}^{\infty}$  satisfying  $\mathbf{X}(0) = \mathbf{x} = (x_j)_{j=1}^{\infty}$  and

$$dX_j(t) = dB_j(t) + \frac{1}{2} \mathbf{d}^{\mu} \left( X_j(t), \sum_{k: k \neq j} \delta_{X_k(t)} \right) dt, \quad j \in \mathbb{N}.$$

Let  $\mathfrak{l}$  be a label map from  $\mathfrak{M}_{s.i.}$  to  $S^{\mathbb{N}}$ , that is, for each  $\xi \in \mathfrak{M}_{s.i.}$ ,  $\mathfrak{l}(\xi) = (\mathfrak{l}(\xi)_j)_{j=1}^{\infty} \in S^{\mathbb{N}}$  satisfies  $\xi = \sum_{j=1}^{\infty} \delta_{\mathfrak{l}(\xi)_j}$ . The map  $\mathfrak{l}$  can be lifted to the map from  $C([0,\infty),\mathfrak{M}_{s.i.})$  to  $C([0,\infty),S^{\mathbb{N}})$ . For  $\Xi \in C([0,\infty),\mathfrak{M}_{s.i.})$  we put

$$\Xi^{\diamond m}(t) = \sum_{j=m+1}^{\infty} \delta_{X_j(t)}$$

for each  $m \in \mathbb{N}$ , where  $(X_j)_{j=1}^{\infty} = \mathfrak{l}(\Xi) \in C([0,\infty), S^{\mathbb{N}})$ . We make the following assumption.

(A4) There exists a subset  $\mathfrak{M}_{SDE}$  of  $\mathfrak{M}_{s,i}$  such that

$$P_{\mu}^{\xi}(\Xi(t) \in \mathfrak{M}_{SDE}) = 1$$
 for any  $\xi \in \mathfrak{M}_{SDE}$ ,

and for each  $\Xi \in C([0,\infty),\mathfrak{M}_{\mathrm{SDE}})$  and each  $m \in \mathbb{N}$ ,

$$(3.2) \quad dY_{j}^{(m)}(t) = dB_{j}(t) - \frac{1}{2}\nabla\Phi(Y_{j}^{(m)}(t))dt - \frac{1}{2}\sum_{k=1, k\neq j}^{m}\nabla\Psi(Y_{j}^{(m)}(t), Y_{k}^{(m)}(t))dt - \frac{1}{2}\int_{\mathfrak{M}}\nabla\Psi(Y_{j}^{(m)}(t), X(t))\Xi^{\diamond m}(dX)dt, \quad 1 \leq j \leq m,$$

$$(3.3) Y_j^{(m)}(0) = \mathfrak{l}(\Xi(0))_j, 1 \le j \le m,$$

has a unique strong solution  $\mathbf{Y}^{(m)} = (Y_1^{(m)}, Y_2^{(m)}, \dots, Y_m^{(m)}).$ 

We also make the following assumptions about the probability measure  $\mu$  (A5) For each  $r, T \in \mathbb{N}$ , there exists a positive constant c such that

$$\int_{S} \operatorname{Erf}\left(\frac{|x|-r}{\sqrt{cT}}\right) \rho^{1}(x) dx < \infty.$$

(A6) The tail  $\sigma$ -field  $\mathcal{T}(\mathfrak{M})$  is  $\mu$ -trivial, that is,  $\mu(A) \in \{0,1\}$  for  $A \in \mathcal{T}(\mathfrak{M})$ .

**Definition 3.2.** Let  $\mu$  be a probability measure on  $\mathfrak{M}$  and let  $\Xi(t)$  be an  $\mathfrak{M}$ -valued process. We say that  $\Xi(t)$  satisfies the  $\mu$ -absolute continuity condition if  $\mu \circ \Xi(t)^{-1}$  is absolutely continuous with respect to  $\mu$  for  $\forall t > 0$ . We say that an  $S^{\mathbb{N}}$ -valued process  $\mathbf{X}(t)$  satisfies the  $\mu$ -absolute continuity condition if  $\mathfrak{u}(\mathbf{X}(t))$  satisfies the  $\mu$ -absolute continuity condition, where  $\mathfrak{u}$  is the map from  $S^{\mathbb{N}}$  to  $\mathfrak{M}$  defined by  $\mathfrak{u}((x_j)_{j=1}^{\infty}) = \sum_{j=1}^{\infty} \delta_{x_j}$ .

Then we have the following theorem.

**Theorem 3.3** ([19]). Suppose that the assumptions in Theorem 3.1 are satisfied. Furthermore assume (A4)–(A6). Then, for  $\mu$ -a.s.  $\xi$ , ISDE (1.6) with  $\mathbf{X}(0) = \mathfrak{l}(\xi)$  has a strong solution satisfying the  $\mu$ -absolute continuity condition, and that pathwise uniqueness holds for ISDE (1.6) with the  $\mu$ -absolute continuity condition.

#### § 4. Applications

Theorems 3.1 and 3.3 can be applied to quite general class of ISDEs. In this section we give some important examples.

**Example 4.1** (Canonical Gibbs measures). Let  $S = \mathbb{R}^d$ ,  $d \in \mathbb{N}$ . Assume that  $\Phi = 0$  and that  $\Psi_0$  is a super stable and regular in the sense of Ruelle [22], and is smooth outside the origin. Let  $\mu$  be a canonical Gibbs measure with the interaction  $\Psi_0$ . Then its logarithmic derivative is

(4.1) 
$$\mathbf{d}^{\mu}\left(x, \sum_{k: k \neq j} \delta_{y_k}\right) = -\sum_{k=1, k \neq j}^{\infty} \nabla \Psi_0(x - y_k).$$

Assume that (A.2) is satisfied. In the case  $d \geq 2$ , there exists a diffusion process associated with  $\mu$  and the labeled process solves

(4.2) 
$$dX_j(t) = dB_j(t) - \frac{1}{2} \sum_{k=1, k \neq j}^{\infty} \nabla \Psi_0(X_j(t) - X_k(t)) dt.$$

In the case  $d=1, \Psi_0$  needs to be sufficient repulsive at the origin to satisfy (A.3).

Assume that (A.5) is satisfied and that, for each  $n \in \mathbb{N}$ , there exist positive constants c, c' satisfying

(4.3) 
$$\sum_{1}^{\infty} \frac{\int_{|x|>r} \rho^{1}(x) dx}{r^{c}} < \infty,$$

(4.4) 
$$\sum_{i,j=1}^{d} \left| \frac{\partial^2}{\partial x_i \partial x_j} \Psi_0(x) \right| \le \frac{c'}{(1+|x|)^{c'+1}},$$

for all  $|x| \ge 1/n$ . In [19, Theorem 3.3] it was proved that, for  $\mu$ -a.s.  $\xi$ , ISDE (4.2) with  $\mathbf{X}(0) = \mathfrak{l}(\xi)$  has a strong solution satisfying the  $\mu_{\mathrm{Tail}}^{\xi}$ -absolute continuity condition, and that pathwise uniqueness holds for ISDE (1.6) with the  $\mu_{\mathrm{Tail}}^{\xi}$ -absolute continuity condition.

**Example 4.2** (Sine random point fields). Let  $\check{\mu}_{\mathsf{bulk},\beta}^N$  be the probability measure defined in (1.3). We denote by  $\mu_{\mathsf{bulk},\beta}^N$  the distribution of  $\sum_{j=1}^N \delta_{x_j}$  under  $\check{\mu}_{\mathsf{bulk},\beta}^N$ . For  $\beta > 0$  the existence of the limit of  $\mu_{\mathsf{bulk},\beta}^N$  as  $N \to \infty$  was shown in Valko-Virág[24]. We denote the limit by  $\mu_{\mathsf{bulk},\beta}$ . In particular, when  $\beta = 2$ ,  $\mu_{\mathsf{bulk},2}$  is the determinantal point process (DPP) with the sine kernel

(4.5) 
$$K_{\sin,2}(x,y) = \frac{\sin(x-y)}{\pi(x-y)},$$

and when  $\beta = 1, 4$ , it is a quaternion determinantal point process [2]. It was shown that  $\mu_{\mathsf{bulk},\beta}$  for  $\beta = 1, 2, 4$  is a quasi-Gibbs measure in [13], and that its logarithmic derivative is

(4.6) 
$$\mathbf{d}^{\mu}\left(x, \sum_{k:k\neq j} \delta_{y_k}\right) = \beta \lim_{\substack{k:k\neq j \\ |y_k| < r}} \frac{1}{x - y_k}$$

in [12]. In [19, Theorem 3.1] it was shown that for  $\mu_{\mathsf{bulk},\beta}$ -a.s.  $\xi$ , ISDE (1.4) with  $\mathbf{X}(0) = \mathfrak{l}(\xi)$  has a strong solution satisfying the  $\mu_{\mathsf{bulk},\beta,\mathsf{Tail}}^{\xi}$ -absolute continuity condition, and that pathwise uniqueness holds for ISDE (1.4) with the  $\mu_{\mathsf{bulk},\beta,\mathsf{Tail}}^{\xi}$ -absolute continuity condition. In the case  $\beta = 2$ , the facts that  $\mathcal{T}(\mathfrak{M})$  is  $\mu_{\mathsf{bulk},2}$ -trivial and  $\mu_{\mathsf{bulk},2,\mathsf{Tail}}^{\xi} = \mu_{\mathsf{bulk},2}$  were shown in [15].

Tsai [23] proved the existence and uniqueness of solutions of ISDE (1.4) for  $\beta \geq 1$  by a different method. Thus it is conjectured that  $\mu_{\mathsf{bulk},\beta}$  is a quasi-Gibbs measure and has a logarithmic derivative of the form (4.6) for  $\beta \geq 1$ .

**Example 4.3** (Airy random point fields). We denote by  $\mu_{\mathsf{soft},\beta}^N$  the distribution of  $\sum_{j=1}^N \delta_{N^{1/6}(x_j-2\sqrt{N})}$  under  $\check{\mu}_{\mathsf{bulk},\beta}^N$ . For  $\beta>0$ , the existence of the limit of  $\mu_{\mathsf{soft},\beta}^N$  as  $N\to\infty$  was shown in Ramírez-Rider-Virág [21]. We denote the limit by  $\mu_{\mathsf{soft},\beta}$ . In particular, when  $\beta=2$ ,  $\mu_{\mathsf{soft},2}$  is the DPP with the Airy kernel

(4.7) 
$$K_{\operatorname{Ai},2}(x,y) = \frac{\operatorname{Ai}(x)\operatorname{Ai}'(y) - \operatorname{Ai}'(x)\operatorname{Ai}(y)}{x - y},$$

where Ai denotes the Airy function and Ai' its derivative [9]. When  $\beta = 1, 4$ , it is a quaternion determinantal point process [2]. In the cases  $\beta = 1, 2, 4$ , it has been proved that the random point field is quasi-Gibbsian [14], and that its logarithmic derivative is

(4.8) 
$$\mathbf{d}^{\mu}\left(x, \sum_{k:k\neq j} \delta_{y_k}\right) = \beta \lim_{r \to \infty} \left\{ \sum_{\substack{k:k\neq j \\ |y_k| < r}} \frac{1}{x - y_k} - \int_{-r}^{r} \frac{\widehat{\rho}(x)dx}{-x} \right\},$$

and for  $\mu_{\mathsf{soft},\beta}$ -a.s.  $\xi$ , ISDE (1.5) with  $\mathbf{X}(0) = \mathfrak{l}(\xi)$  has a strong solution satisfying the  $\mu_{\mathsf{soft},\beta,\mathrm{Tail}}^{\xi}$ -absolute continuity condition, and pathwise uniqueness holds for ISDE (1.5) with the  $\mu_{\mathsf{soft},\beta,\mathrm{Tail}}^{\xi}$ -absolute continuity condition [18, Theorem 2.3]. In the case  $\beta = 2$  the facts that  $\mathcal{T}(\mathfrak{M})$  is  $\mu_{\mathsf{soft},2}$ -trivial and that  $\mu_{\mathsf{soft},2,\mathrm{Tail}}^{\xi} = \mu_{\mathsf{soft},2}$  were shown in [15].

Determining whether  $\mu_{\mathsf{soft},\beta}$  has the quasi-Gibbs property for general  $\beta$  and finding its logarithmic derivative is (4.8) are interesting and important problems.

**Example 4.4** (Bessel random point field). Let  $S = [0, \infty)$  and  $1 \le \alpha < \infty$ . Let  $\mu_{\mathsf{hard},2}$  be the determinantal point process with Bessel kernel

(4.9) 
$$K_{J_{\alpha}}(x,y) = \frac{J_{\alpha}(\sqrt{x})\sqrt{y}J_{\alpha}'(\sqrt{y}) - \sqrt{x}J_{\alpha}'(\sqrt{x})J_{\alpha}(\sqrt{y})}{2(x-y)}.$$

In [6] it was shown that  $\mu_{\mathsf{hard},2}$  is a quasi-Gibbs measure and that the related process is the unique strong solution of the ISDE

$$dX_j(t) = dB_j(t) + \left\{ \frac{\alpha}{2X_j(t)} + \sum_{k=1, k \neq j}^{\infty} \frac{1}{X_j(t) - X_k(t)} \right\} dt$$

with the  $\mu_{\mathsf{hard},2}$ -absolute continuity condition.

**Example 4.5** (Ginibre random point field). Let  $S = \mathbb{R}^2$  be identified as  $\mathbb{C}$ . Let  $\mu_{\mathsf{Gin}}$  be the DPP with the kernel  $K_{\mathsf{Gin}} : \mathbb{C} \times \mathbb{C} \to \mathbb{C}$  defined by

(4.10) 
$$K_{\mathsf{Gin}}(x,y) = \frac{1}{\pi} e^{-|x|^2/2 - |y|^2/2} e^{x\overline{y}}.$$

In [13] it was shown that  $\mu_{Gin}$  is a quasi-Gibbs measure, and in [12] that the related process is a solution of the ISDE

(4.11) 
$$dX_j(t) = dB_j(t) - X_j(t)dt + \lim_{r \to \infty} \sum_{\substack{k: k \neq j \\ |X_k(t)| < r}} \frac{X_j(t) - X_k(t)}{|X_j(t) - X_k(t)|^2} dt.$$

The pathwise uniqueness of solutions of (4.11) with the  $\mu_{Gin}$ -absolute continuity condition was shown in [19].

#### § 5. Remarks

In the previous section we gave some examples of DPPs that are not canonical Gibbs measures but quasi-Gibbs measures. It is expected that quite general DPPs have the quasi-Gibbs property. We thus present examples of DPPs related to random matrix theory or non-colliding Brownian motions, whose quasi-Gibbs property have not been shown.

**Example 5.1** (Pearcey process). Consider 2N noncolliding Brownian motions, in which all particles start from the origin and N particles end at  $\sqrt{N}$  at time t = 1, and the other N particles end at  $-\sqrt{N}$  at t = 1. We denote the system by  $(X_1^N(t), \ldots, X_{2N}^N(t))$ ,  $0 \le t \le 1$ . When N is very large, there is a cusp at  $x_0^N = 0$  when  $t_0 = \frac{1}{2}$ , that is, before time  $t_0$  particles are in one interval with high probability, while after time  $t_0$  they are separated into two intervals by the origin. We denote the distribution

$$\sum_{j=1}^{2N} \delta_{2^{3/2}(2N)^{1/4}X_{j}^{N}(\frac{1}{2})}$$

on  $\mathfrak{M}$  by  $\mu_{\mathsf{pearcey}}^N$ . It was proved in Adler-Orantin-von Moerbeke [1] that

$$\mu_{\text{pearcey}}^N \to \mu_{\text{pearcey}}, \quad \text{weakly as } N \to \infty$$

and that  $\mu_{pearcey}$  is the DPP  $K_{pearcey}(x, y)$  given by

$$K_{\mathsf{pearcey}}(x,y) = \frac{P(x)Q''(y) - P'(x)Q'(y) + P''(x)Q(y)}{x - y}, \quad x, y \in \mathbb{R},$$

with

$$Q(y) = \frac{i}{2\pi} \int_{-i\infty}^{i\infty} e^{-u^4/4 - uy} du$$
 and  $P(x) = \frac{1}{2\pi i} \int_C e^{v^4/4 + vx} dv$ ,

where the contour C is given by the ingoing rays from  $\pm \infty e^{i\pi/4}$  to 0 and the outgoing rays from 0 to  $\pm \infty e^{-i\pi/4}$ . These integrals are known as Pearcey's integrals [20].

**Example 5.2** (Tacnode process). Consider two groups of non-colliding pinned Brownian motions  $(X_1^N(t), \ldots, X_{2N}^N(t))$  in the time interval  $0 \le t \le 1$ , where one group of N particles starts and ends at  $\sqrt{N}$  and the other group of N particles starts and ends at  $-\sqrt{N}$ . The distribution  $(N^{1/6}X_1^N(\frac{1}{2}), N^{1/6}X_2^N(\frac{1}{2}), \ldots, N^{1/6}X_{2N}^N(\frac{1}{2}))$  on the Weyl chamber of type  $A_{2N-1}$ 

$$\mathbb{W}_{2N} = \left\{ \mathbf{x} = (x_1, x_2, \cdots, x_{2N}) : x_1 < x_2 < \cdots < x_{2N} \right\},\,$$

is given by

$$m_{\text{tac}}^{2N}(d\mathbf{x}_{2N}) = \frac{1}{Z} \left[ \det_{1 \le i,j \le 2N} \left( e^{-2|x_i - a_j|^2} \right) \right]^2,$$

where  $a_j = -\sqrt{N}$  for  $1 \le j \le N$  and  $a_j = \sqrt{N}$  for  $N+1 \le j \le 2N$ . We denote the distribution of  $\sum_{j=1}^{2N} \delta_{N^{1/6}x_j}$  under  $m_{\sf tac}^{2N}$  by  $\mu_{\sf tac}^N$ . It was proved in Delvaux-Kuijlaars-Zhang [3] and Johansson [7] that

$$\mu_{\mathsf{tac}}^N \to \mu_{\mathsf{tac}}, \quad \text{weakly as } N \to \infty$$

and that  $\mu_{tac}$  is the DPP with the correlation kernel

$$K_{\mathsf{tac}}(x,y) \equiv L_{\mathsf{tac}}(x,y) + L_{\mathsf{tac}}(-x,-y), \quad x,y \in \mathbb{R},$$

where

$$L_{\mathsf{tac}}(x,y) = K_{\mathsf{Ai},2}(x,y)$$

$$+2^{1/3} \int_{(0,\infty)^2} du dv \ \mathsf{Ai}(y+2^{1/3}u) R(u,v) \mathsf{Ai}(x+2^{1/3}v)$$

$$-2^{1/3} \int_{(0,\infty)^2} du dv \ \mathsf{Ai}(-y+2^{1/3}u) \mathsf{Ai}(u+v) \mathsf{Ai}(x+2^{1/3}v)$$

$$-2^{1/3} \int_{(0,\infty)^3} du dv dw \ \mathsf{Ai}(-y+2^{1/3}u) R(u,v) \mathsf{Ai}(v+w) \mathsf{Ai}(x+2^{1/3}w).$$

Here, R(x,y) is the resolvent operator for the restriction of the Airy kernel to  $[0,\infty)$ , that is, the kernel of the operator

(5.1) 
$$R = (I - K_{Ai})^{-1} K_{Ai}$$

on  $L^2[0,\infty)$ .

In [3, 7] it was also shown that

$$\Xi^{N}(t) \equiv \sum_{j=1}^{2N} \delta_{N^{1/6}X_{j}(\frac{1}{2}+N^{-1/3}t)} \to \Xi(t), \quad \text{as } N \to \infty,$$

in the sense of finite-dimensional distributions, where  $\Xi(t)$  is a reversible process with reversible measure  $\mu_{\text{tac}}$ . We expect that  $\Xi(t)$  is the diffusion process associated with the Dirichlet form  $(\mathcal{E}^{\mu_{\text{tac}}}, \mathcal{D}^{\mu_{\text{tac}}})$ .

#### References

- [1] Adler, M., Orantin, N. and von Moerbeke, P., Universality for the Pearcey process, *Physica D* **239** (2010), 924–941.
- [2] Anderson, G. W., Guionnet, A. and Zeitouni, O., An Introduction to Random Matrices, Cambridge university press, 2010.
- [3] Delvaux, S., Kuijlaars, B.J. and Zhang, L., Critical Behavior of Nonintersecting Brownian motions at a Tacnode, *Comm. Pure Appl. Math.* **64** (2011), 1305–1383.
- [4] Dyson, F. J., A Brownian-motion model for the eigenvalues of a random matrix, *J. Math. Phys.* **3** (1962), 1191–1198.
- [5] Fukushima, M., Oshima, Y. and Takeda, M., Dirichlet forms and symmetric Markov processes, 2nd ed., Walter de Gruyter, 2011.
- [6] Honda, R. and Osada, H., Infinite-dimensional stochastic differential equations related to the Bessel random point fields, *Stochastic Processes and their Applications* **125** (2015), 3801–3822.
- [7] Johansson, K., Non-colliding Brownian motions and the extended Tacnode process, *Comm. Math. Phys.*, **269** (2012), 571–609.
- [8] Kawamoto, Y. and Osada, H., Finite particle approximations of interacting Brownian motions in infinite dimensions and SDE gaps, (in preparation).
- [9] Mehta, M. L., Random Matrices. 3rd edition, Amsterdam: Elsevier, 2004
- [10] Osada, H., Dirichlet form approach to infinite-dimensional Wiener processes with singular interactions, *Commun. Math. Phys.*, **176** (1996), 117–131.
- [11] Osada, H., Tagged particle processes and their non-explosion criteria, J. Math. Soc. Japan, **62** (2010), 867–894.
- [12] Osada, H., Infinite-dimensional stochastic differential equations related to random matrices, *Probability Theory and Related Fields*, **153** (2012), 471–509.
- [13] Osada, H., Interacting Brownian motions in infinite dimensions with logarithmic interaction potentials, Ann. of Probab.. 41 (2013), 1–49.

- [14] Osada, H., Interacting Brownian motions in infinite dimensions with logarithmic interaction potentials II: Airy random point field, *Stochastic Processes and their Applications*, **123** (2013), 813–838.
- [15] Osada, H. and Osada, S., Discrete approximations of determinantal point processes on continuous space: tree representations and tail triviality, (preprint) arXiv:1517677 [math.PR].
- [16] Osada, H. and Tanemura, H., Cores of Dirichlet forms related to Random Matrix Theory, *Proc. Jpn. Acad.*, Ser. A, **90** (2014), 145–150.
- [17] Osada, H. and Tanemura, H., Strong Markov property of determinantal processes with extended kernels, *Stochastic Processes and their Applications*, **126** (2016), 186–208.
- [18] Osada, H. and Tanemura, H., Infinite-dimensional stochastic differential equations arising from Airy random point fields, (preprint) arXiv:1408.0632 [math.PR].
- [19] Osada, H. and Tanemura, H., Infinite dimensional stochastic differential equations and tail  $\sigma$ -fields, (preprint) arXiv:1412.8674 [math.PR].
- [20] Pearcey, T., The structure of an electomagnetic field in the neighbourhood of a cusp of a caustic, *Phil. Mag.* **37** (1946), 311–317.
- [21] Ramírez, J.A., Rider, B. and Virág, B., Beta ensembles, stochastic Airy spectrum, and a diffusion, *Journal of the American Mathematical Society*, **24** (2011), 919–944.
- [22] Ruelle, D., Superstable interactions in classical statistical mechanics, *Commun. Math. Phys.* **18** (1970), 127–159.
- [23] Tsai, Li-Cheng, Infinite dimensional stochastic differential equations for Dyson's model, Probability Theory and Related Fields, DOI 10.1007/s00440-015-0672-2.
- [24] Valkó, B. and Virág, B., Continuum limits of random matrices and the Brownian carousel, *Inventions* **177** (2009), 463–508.