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**Occupied Electronic States of Li in Li, LiO, and LiO  
Analyzed by Soft X-ray Emission Spectroscopy**

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# Occupied Electronic States of Li in Li, Li<sub>2</sub>O<sub>2</sub>, and Li<sub>2</sub>O Analyzed by Soft X-ray Emission Spectroscopy

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## ABSTRACT

Lithium metal and lithium oxides are components of lithium–oxygen (Li–O<sub>2</sub>) batteries. In order to accurately identify Li-compounds and understand the degradation mechanism, fundamental knowledge on electron structures of constituent elements is vital. However, experimentally-derived occupied states of Li has been missing due to their intrinsic difficulties in detection. Herein, using soft X-ray emission spectroscopy, ultrahigh-energy-resolution spectra of Li–K were collected for three critical Li-compounds; Li, Li<sub>2</sub>O<sub>2</sub>, and Li<sub>2</sub>O. Large chemical shifts to lower energies and peak broadening were observed in compound specific Li–K and O–K spectra. Theoretical calculations confirm that these changes derive from the characteristic electronic configurations of 1s and 2p states with core level shifts in <sup>+</sup>Li. The large chemical shift (~4.6 eV) between the Li and Li<sub>2</sub>O peaks was utilized to visualize the chemical state mapping of the Li metal/oxide phase, facilitating the identifications of chemical phases in Li compounds.

## Introduction

Li–O<sub>2</sub> air battery has attracted significant attentions as alternative means of electrochemical storage because of their high specific energy densities.<sup>1</sup> In Li–O<sub>2</sub> cells, Li metal anode is oxidized and release Li<sup>+</sup> ions into electrolyte on discharge; during charging, the process is reversed. Once conducted through the electrolyte, Li<sup>+</sup> reacts with O<sub>2</sub><sup>2-</sup> to form the final discharge product, lithium peroxides (Li<sub>2</sub>O<sub>2</sub>), at cathode. The discharge product of lithium oxide Li<sub>2</sub>O is also possible which increases the specific energy stored, but reverse charging is hampered due to its insulating property.<sup>2,3</sup> In addition

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6 to electrochemical storage, Li, Li-alloys, and lithium oxides can serve as an intense neutron source<sup>4,5</sup>  
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9 and a tritium breeder in fusion reactors.<sup>6,7</sup> Li metal and oxides are chemically reactive in ambient  
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12 atmosphere to form impurity phases and corrosion products that can degrade performance and raise  
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15 safety concerns.<sup>8,9</sup> Chemical phase analysis of Li-containing products contributes to our  
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18 understandings of chemical reactions in the system. Chemical state distribution of Li compounds were  
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21 previously reported using electron energy loss spectroscopy and Auger electron microscopy.<sup>10,11</sup>  
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24 Analysis of Li compounds is often hindered by (i) low atomic number of Li (i.e. poor X-ray scattering  
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27 power and extremely low photon energies at the *K*-edge), (ii) its high affinity to moisture resulting in  
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30 compositional changes,<sup>12,13</sup> and (iii) X-ray/electron irradiation-induced sample damage that occurs  
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33 during characterizations. X-ray photoelectron spectroscopy (XPS) is a common technique to obtain  
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36 quantitative insights of core electrons, including Li 1s electrons.<sup>3,6,14</sup> But, even with high-resolution  
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39 XPS, Li 1s peaks overlap as the shift of the binding energy is small compared to the energy resolution  
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42 of XPS; a shift of less than 1.5 eV is observed for the Li 1s core levels between Li metal and Li<sub>2</sub>O.<sup>14</sup>  
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45 Thus, additional peak decompositions by fitting are required. Léon *et al.* investigated the valence  
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48 electron structures of Li<sub>2</sub>O and Li<sub>2</sub>O<sub>2</sub> at the O *K*-edge using soft X-ray emission spectroscopy (SXES)  
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51 at a synchrotron facility; the resulting spectra describe the valence electron structure of O.<sup>15</sup> We note  
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54 that Li valence electron structures must be determined experimentally for understanding nature of  
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57 chemical bonds and identifying oxide phases with Li. High resolution *K*-edge spectral data can be  
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6 exploited to visualize chemical state mapping on a microscopic scale, as demonstrated in our previous  
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9 studies for Be ( $Z = 4$ ) and B ( $Z = 5$ ) compounds.<sup>16,17</sup> Thus, in this study, the ultrahigh energy resolution  
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12 ( $\sim 0.2$  eV) valence electron structures of Li in Li, Li<sub>2</sub>O<sub>2</sub>, and Li<sub>2</sub>O were experimentally analyzed. We  
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15 employ SXES with a high sensitive X-ray charge-coupled device attached to an electron probe micro-  
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18 analyzer (EPMA).<sup>18,19</sup> SXES is a bulk-sensitive measurement that covers ultra-soft X-rays in the  
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21 energy range of 50–170 eV including Li- $K$  emission at 54.1 eV (i.e.  $2p \rightarrow 1s$  transition in Li). As the  
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24 samples are metastable and may decompose during sample preparation or characterization, the  
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27 obtained spectra are directly compared with density of states (DOS) calculated using density functional  
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30 theory (DFT) with all-electron configurations for Li ( $1s^2 2s^1$ ). Due to large errors in band calculations  
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33 in oxides,<sup>20</sup> not only a conventional functional based on the generalized gradient approximation  
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36 (GGA) but also a hybrid functional are utilized in this study. Finally, we show chemical state mappings  
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39 of Li and Li<sub>2</sub>O based on the knowledge of electron structures at the  $K$ -edge.  
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## 41 **Experimental and Computational Methods**

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43 Lithium metal rods (>99.9%, Mitsuwa Chemical Co., Ltd.), lithium peroxide (>99%, Kojundo  
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46 Chemical Laboratory Co., Ltd.) and lithium oxide powders (>99%, Kojundo Chemical Laboratory Co.,  
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49 Ltd.) were analyzed. Li metal rod oxidation was carried out in air at 25 °C for 10 min, thereafter the  
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52 sample was covered with oxidized layer. Part of the surface layer was physically peeled off and  
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55 interface between metal and oxide phases was analyzed by an area measurement using SXES. X-ray  
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58 powder diffraction (XRD) was carried out in air using Co- $K\alpha$  (RINT TTR-III, Rigaku) in the range  
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of  $2\theta = 10\text{--}80^\circ$  at intervals of  $0.02^\circ$  step. XRD data were analyzed using the RIETAN-FP program.<sup>21</sup>

The samples were analyzed by a JXA-8500F field emission EPMA by JEOL equipped with SXES (SS-94000SXES by JEOL). Sample powders were pressed on to carbon tape mounted on the sample holder. The measurement time and the acceleration voltages of the electron beam were 60 s and 2.0 kV.

DFT calculations were performed using the Perdew–Burke–Ernzerhof (PBE) functional and HSE06 functional implemented in the Vienna ab initio simulation package (VASP).<sup>22–25</sup> The electron configurations  $1s^2 2s^1$ ,  $2s^2 2p^4$  were used for the Li and O atoms, respectively. The  $k$ -point meshes of  $6\times 6\times 6$ ,  $8\times 8\times 4$ , and  $6\times 6\times 6$  were generated using the Monkhorst-Pack scheme for  $2\times 2\times 2$  supercell of Li (cubic;  $Im\bar{3}m$ ),  $2\times 2\times 1$  supercell of  $\text{Li}_2\text{O}_2$  (hexagonal;  $P6_3mmc$ ), and  $2\times 2\times 2$  supercell of  $\text{Li}_2\text{O}$  (cubic;  $Fm\bar{3}m$ ). A plane-wave cutoff energy of 400 eV was used for calculations. Self-consistency was achieved with a tolerance for a total energy of  $10^{-5}$  eV, and the atomic positions were relaxed until the force was less than  $0.02$  eV  $\text{\AA}^{-1}$ . The energy of DOS was referenced to the Li 1s band centroid,  $E_{\text{Li-1s}}$ .

$E_{\text{Li-1s}}$  was calculated as follows;

$$E_{\text{Li-1s}} = \frac{\int_{-\infty}^{E_{\text{max}}} E g_{\text{Li-1s}}(E) dE}{\int_{-\infty}^{E_{\text{max}}} g_{\text{Li-1s}}(E) dE} \quad (1)$$

where  $g_{\text{Li-1s}}(E)$  and  $E_{\text{max}}$  are the density of states of Li 1s orbitals at  $E$  and maximum energy of the Li 1s orbitals, respectively.

## Results and Discussion

Li<sub>2</sub>O<sub>2</sub> and Li<sub>2</sub>O commercial powder samples were stored in an Ar-filled groove box and were analyzed by XRD. Based on previous structural assessments on Li<sub>2</sub>O<sub>2</sub>,<sup>26,27</sup> Rietveld refinements were carried out using the structure proposed by Föppl et al.<sup>28</sup> The results of multi-phase Rietveld analysis (Figure S1 in the Supporting Information (SI)) showed minor peaks for LiOH (*P4/nmm*) in addition to the Li<sub>2</sub>O<sub>2</sub> (*P6<sub>3</sub>mmc*) and Li<sub>2</sub>O (*Fm $\bar{3}$ m*) phases. The Li<sub>2</sub>O<sub>2</sub> and Li<sub>2</sub>O powders contained ~6.3 and ~11.7 w.t.% LiOH as impurity phases, whereas no other impurity phase was found. Figure 1 shows Li DOS in Li, Li<sub>2</sub>O<sub>2</sub> and Li<sub>2</sub>O using the PBE exchange-correlation functional and Heyd-Scuseria-Ernzerhof (HSE) hybrid functional. In Figure 1a, energetic difference in HSE between Li 1s centroid ( $E_{\text{Li-1s}}$ ) and  $E_{\text{F}}$  in Li metal was 50.1 eV and had a significantly smaller error than that in PBE (46.4 eV). The value was nevertheless underestimated compared to the experimental (54.1 eV). PDOS calculations for Li<sub>2</sub>O<sub>2</sub> (Figure 1b) indicated valence electron states of Li<sup>+</sup> ions composing of  $\sigma_{\text{g}}$  (bonding orbitals:  $p_z$  from -6.8 to -3.5 eV),  $\pi_{\text{u}}$  (bonding orbitals:  $p_x$  and  $p_y$  from -6.8 to -3.2 eV) and  $\pi_{\text{g}}^*$  (antibonding:  $p_x$  and  $p_y$  from -2.5 to 0 eV),<sup>29</sup> with a short gap approximately at -3.2 eV (Figure 1b and Figure S2 in SI). The Li<sub>2</sub>O band gap in PBE (4.9 eV) in Figure 1c was smaller than that in HSE (5.9 eV) and the experimental value (7.99 eV).<sup>30</sup> Energetic differences between  $E_{\text{Li-1s}}$  and valence band maximum (VBM) in Li<sub>2</sub>O<sub>2</sub> and Li<sub>2</sub>O were -46.6 and -46.0 eV in HSE (Figure 1 and Figure S3 in SI). These indicate core level shifts as large as 3.5 eV (Li<sub>2</sub>O<sub>2</sub>) and 4.1 eV (Li<sub>2</sub>O) for Li<sup>+</sup>. These shifts are at least an order of magnitude higher than the SXES energy resolutions, indicating their applicability

for chemical state mapping. Core electron state in  $\text{Li}_2\text{O}_2$  were at least two times wider than those in  $\text{Li}$  and  $\text{Li}_2\text{O}$ , due to difference of electron structures at two Li atom sites in  $\text{Li}_2\text{O}_2$  (Figure S3 in SI). This can be a cause of broad  $\text{Li-K}$  spectrum for  $\text{Li}_2\text{O}_2$ .

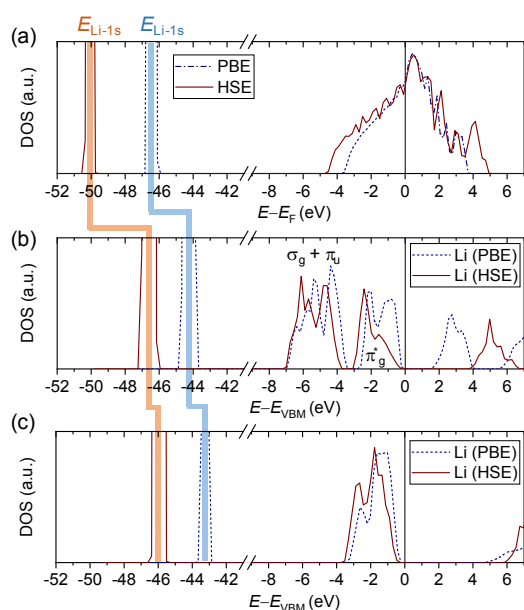


Figure 1 Electronic configurations of Li in Li (a),  $\text{Li}_2\text{O}_2$  (b), and  $\text{Li}_2\text{O}$  (c) obtained by DFT using PBE and HSE. Orange and blue solid lines represent positions of Li 1s centroid ( $E_{\text{Li-1s}}$ ) and indicate core level shifts of  $\text{Li}^+$ .

Figure 2 shows the SXES spectra of the Li metal,  $\text{Li}_2\text{O}_2$ , and  $\text{Li}_2\text{O}$  samples. In the obtained spectra,  $n^{\text{th}}$ -order peaks for Li ( $n = 1$ ) and O ( $n = 3-10$ ) were observed; peaks of C ( $n = 2$  and 3) were also observed due to the use of carbon conduction tape. The  $\text{Li-K}$  spectrum of Li metal showed an asymmetric Doniach-Šunjić line shape in the range of 52–54.5 eV (Fig. 2b). The  $\text{Li-K}$  peak intensities for  $\text{Li}_2\text{O}_2$  and  $\text{Li}_2\text{O}$  were significantly lower than that of Li. Ionization to  $\text{Li}^+$  causes numerous Li valence electrons to be lost, and results in poor signal to noise (S/N) ratios. In Figure 3, the SXES

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Li-*K* spectra are directly compared with the DOS and PDOS of Li using HSE. The energy was referenced to  $E_{\text{Li-1s}}$ . The calculated Li 2p occupied states were convoluted with Gaussian functions with full width at half-maxima (FWHM) of 0.6 eV (Li and Li<sub>2</sub>O) and 0.9 eV (Li<sub>2</sub>O<sub>2</sub>) because of the difference in 1s band width (Figure S3 in SI). The shapes of Li-*K* spectra agreed well with the convoluted 2p state. The calculated energies using HSE were underestimated as large as 4.0 eV for Li metal and 5.3 eV for the oxides. A comparison between the experimental spectra and DOS using PBE (Figure S3 in SI) supports that the obtained spectral shapes describe Li valence electronic states. The Li-*K* spectrum for Li<sub>2</sub>O<sub>2</sub> mainly derived from antibonding  $\pi_g^*$  orbitals appeared at 49–52 eV, while the peak of bonding orbitals ( $\sigma_g$  and  $\pi_u$ ) were below the SXES low-energy limit. But, the contributions of both bonding and antibonding orbitals were seen in the broad peak of O-*K*(4) for Li<sub>2</sub>O<sub>2</sub> in Figure 2c. The O-*K*(4) FWHM of the Li<sub>2</sub>O<sub>2</sub> sample was 2.26 times wider than that from the Li<sub>2</sub>O sample (Figure 2c), which was consistent with the valence band width ratio of 2.12 calculated using HSE (Figure 1b,c). The wide valence band was attributed to the anisotropic 2p states in Li<sub>2</sub>O<sub>2</sub> (Figure S2 in SI). Although the unidentified O-*K* peak in Li<sub>2</sub>O was observed at 528.9 eV in the previous SXES results,<sup>15</sup> no peak was found at the corresponding position of 132.2 eV for O-*K*(4) in the Li<sub>2</sub>O spectrum in the present work.



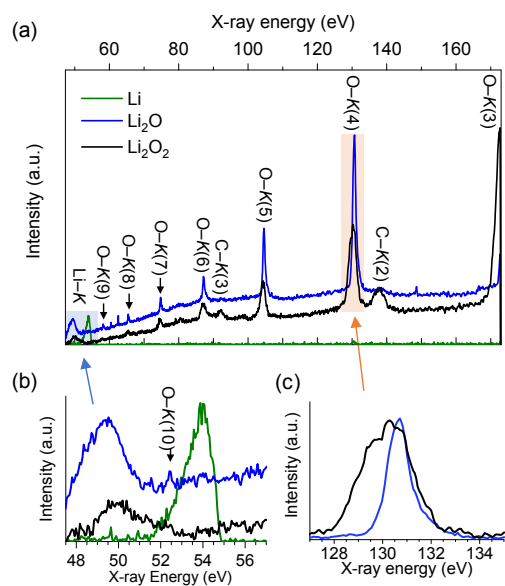


Fig. 2 SXES spectra from the Li metal, Li<sub>2</sub>O<sub>2</sub>, and Li<sub>2</sub>O samples in the energy ranges of 47.5–173 eV (a), 47.5–57 eV for Li–K (b), and 127–135 eV for O–K(4) (c), in which number in bracket denotes the X-ray diffraction order.

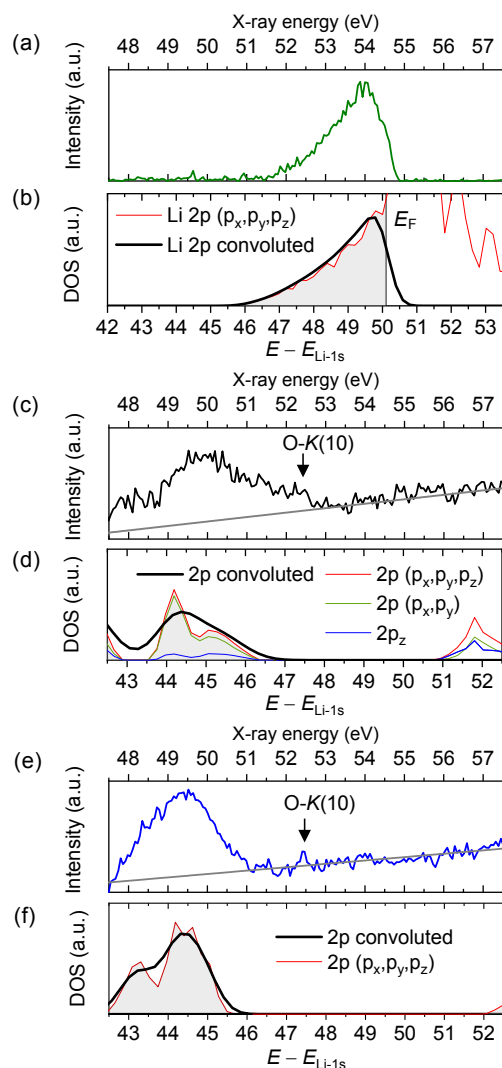


Figure 3 The experimental Li- $K$  spectra for Li (a),  $\text{Li}_2\text{O}_2$ (c), and  $\text{Li}_2\text{O}$  (e) using the SXES compared to the convoluted DOSs of occupied Li 2p states in Li metal (b),  $\text{Li}_2\text{O}_2$ (d), and  $\text{Li}_2\text{O}$  (f) using HSE. Gray solid lines in panel (c) and (e) denotes linear background. Calculated energy is referenced to Li 1s centroid ( $E_{\text{Li-1s}}$ ) where occupied Li 2p states below the Fermi energy and the VBMs are colored by gray. A comparison between the spectra and DOS using PBE is shown in Figure S3 in SI.

The chemical state and element mappings on a Li metal sample that has been naturally oxidized in air with a partial metallic surface are shown in Figure 4. The SXES spectra were collected at  $40 \times 32$  measuring points (1280 points), with an acquisition time of 30 s per pixel (pixel size:  $1 \mu\text{m} \times 1 \mu\text{m}$ ). From the naturally oxidized Li metal sample, the Li- $K$  spectra of  $\text{Li}_2\text{O}$  ( $E < 51$  eV) and Li ( $> 51$

eV) were observed without peak overlap. No  $\text{Li}_2\text{O}_2$  peak was found from the collected data, as predicted by the MALT thermodynamic database (Table S1 in SI). Chemical state mappings were constructed by calculating integrated intensities in the following energy ranges:  $\text{Li}_2\text{O}$  (47.5–51.0 eV) and Li metal (51.0–55.0 eV). Chemical state mappings (Figure 4a,b) visualize Li metal/oxide phase distribution on the sample; the metallic Li phase was distributed on the right side of the observed area, while the  $\text{Li}_2\text{O}$  phase lies on the left. Together with O element mapping (Figure 4c), it is possible to distinguish the  $\text{Li}_2\text{O}$  phase from other oxide phases with lacking Li. To overcome Li-K weak emissions from oxide samples, an identification whether  $\text{Li}_2\text{O}_2$  or  $\text{Li}_2\text{O}$  phase can be accomplished by integrating their antibonding orbitals appeared in the range of 128–130 eV for O-K(4) (Figure 2c).

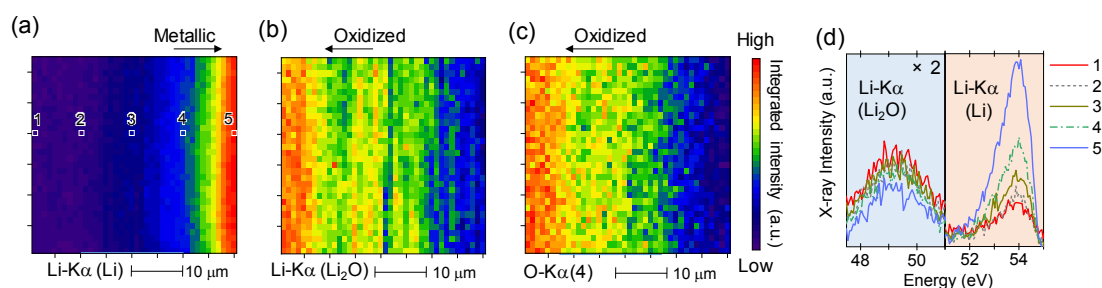


Fig.4 Chemical state mappings of Li metal ( $E = 51.0\text{--}55.0$  eV) (a) and  $\text{Li}_2\text{O}$  ( $E = 47.5\text{--}51.0$  eV) (b), element mapping of O using 4<sup>th</sup> order O-K ( $E = 129\text{--}132$  eV) (c), and Li-K spectra at five measuring area (d). The positions of the five measuring area are shown in panel (a).

## Conclusion

The electron structures in Li metal,  $\text{Li}_2\text{O}_2$ , and  $\text{Li}_2\text{O}$  were experimentally analyzed using soft X-ray emission spectroscopy. The chemical shifts of Li-K and peak broadening of  $\text{Li}_2\text{O}_2$  O-K arose from

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6 Li core level shifts and characteristic 2p configurations as confirmed by theoretical calculations using  
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9 the hybrid functional. A large chemical shift (~4.6 eV) without peak overlap was utilized to visualize  
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12 the chemical state mappings of Li and lithium oxide. Insights of Li electron structures are a basis for  
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15 investigating bond nature of Li and analyzing chemical state which could help further studies for  
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18 controlling degradation processes.  
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## 23 ASSOCIATED CONTENT

### 24 Supporting Information

25 The supporting information is available free of charge on the ACS Publications website.

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27 Rietveld patterns of the X-ray diffraction data, DOS of the Li and lithium oxides using HSE, core electron states using  
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29 HSE and crystal structure of  $\text{Li}_2\text{O}_2$ , comparison between the experimental spectra and DOS using PBE, phase change  
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31 of Li metal in air by thermodynamic calculations (Figures S1–S4 and Table S1).

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### 44 Notes

45 The authors declare no competing financial interest.  
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## TOC GRAPHICS

