



Development and application of a ^3He Neutron Spin Filter at J-PARC

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ABSTRACT

We are developing a neutron polarizer with polarized ^3He gas, referred to as a ^3He spin filter, based on the Spin Exchange Optical Pumping (SEOP) for polarized neutron scattering experiments at Materials and Life Science Experimental Facility (MLF) of Japan Proton Accelerator Research Complex (J-PARC). A ^3He gas-filling station was constructed at J-PARC, and several ^3He cells with long spin relaxation times have been fabricated using the gas-filling station. A laboratory has been prepared in the MLF beam hall for polarizing ^3He cells, and compact pumping systems with laser powers of 30 W and 110 W, which can be installed onto a neutron beamline, have been developed. A ^3He polarization of 85% was achieved at a neutron beamline by using the pumping system with the 110 W laser. Recently, the first user experiment utilizing the ^3He spin filter was conducted, and there have been several more since then. The development and utilization of ^3He spin filters at MLF of J-PARC are reported.

1. Introduction

A ^3He spin filter is a neutron polarization device in which polarized ^3He gas is enclosed in a special glass cell that does not contain boron. ^3He spin filters are widely used to polarize neutrons and analyze neutron spins in neutron scattering experiments because of its capabilities to polarize neutrons in broad energy range and cover wide solid angle [1–6]. Since ^3He nuclei have a strongly spin-dependent neutron absorption cross section, neutrons passing through polarized ^3He gas are polarized. ^3He nuclei are polarized using laser techniques known as Spin Exchange Optical Pumping (SEOP) [7] or Metastability Optical Pumping (MEOP) [8]. In SEOP, ^3He gas, N_2 gas, and pure rubidium or rubidium and potassium are encapsulated within a glass cell. The Rb or Rb-K cells are heated to around 170 °C or 220 °C, respectively, to evaporate the alkali metals, and circularly polarized laser light irradiates the cell. Unpaired electrons of alkali atoms are polarized by optical pumping, and the spins of the electrons are exchanged with those of ^3He nuclei by hyperfine interaction. A diode laser array with output power on the order of 100 W and a wavelength of 794.7 nm, which

corresponds to the D_1 absorption line of Rb, is typically employed. The ^3He polarization can be improved by using Rb-K compared to the pure Rb case because of the high optical pumping efficiency of Rb-K [9]. When the laser irradiation is stopped, the ^3He polarization exponentially decays due to non-uniformity of the surrounding magnetic field, collisions of ^3He atom with each other, and collisions of ^3He atoms with impurities in the glass cell and a glass wall.

Ex-situ and *in-situ* methods are used for ^3He spin filters with SEOP. In the *ex-situ* method, a ^3He cell is polarized in a pumping station away from a neutron beamline, transferred in a magnetic cavity to keep the ^3He polarization, and installed to a neutron beamline. Since the ^3He polarization decays during a neutron experiment, the ^3He cell needs to be periodically exchanged with another polarized ^3He cell for long-running experiments. However, the advantage of using *ex-situ* method is that it can be used for a neutron beamline with limited space, and the installation of the instrument is straightforward. This is beneficial at spallation neutron sources, which generally have small experimental spaces. As there are many experiments with short

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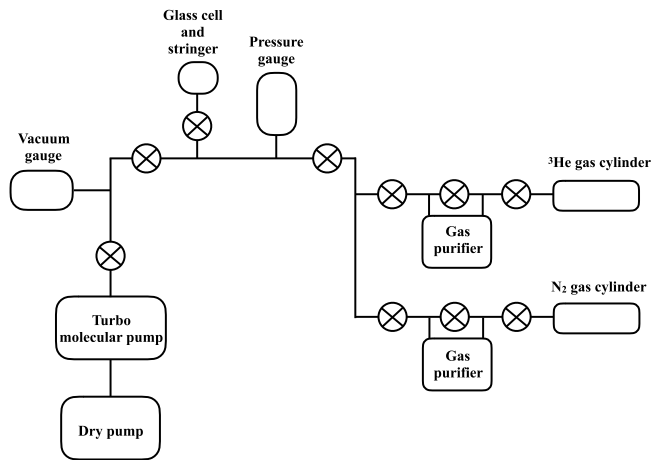


Fig. 1. The schematic diagram of the gas-filling station to fabricate ^3He cells.

measuring time at intense neutron beam facilities, those may be done without re-polarizing the ^3He cell. In order to keep high ^3He polarization during the experiment, a long relaxation time of ^3He polarization is required, which makes details of the fabrication process of the ^3He cells important.

On the other hand, in the *in-situ* method, the ^3He cell is installed along with a laser system onto a neutron beamline. This method is suitable for longer experiments because a stable ^3He polarization can be maintained for weeks. However, a larger space compared to the space required for the *ex-situ* method is needed, and the safety regulations involving the use of high powered lasers on a neutron beamline must be considered. For details of the technique for optically polarized ^3He and its application, please see a reviewed paper Ref. [10].

At MLF of J-PARC, world class intense pulsed neutron beams are provided to 23 neutron beamlines, and many neutron scattering experiments are conducted [11]. In Japan, the fundamental development of ^3He polarization technique based on SEOP began in the 1990s [12], and further development is still ongoing by Ino et al. at High Energy Accelerator Research Organization [13–17]. An *in-situ* SEOP system dedicated for a polarized neutron spectrometer POLANO [11,18] at beamline No. 23 has been developed, and now under commissioning [15,16]. In order to promote user experiments using ^3He spin filters on the other beamlines at MLF, we are carrying out the development of ^3He spin filters with both *in-situ* and *ex-situ* methods for their versatile uses on site at J-PARC. Experiments demonstrating their use have been performed on several beamlines. User experiments and fabrication of ^3He cells have recently begun. In this paper, we report development and utilization of ^3He spin filters at MLF of J-PARC.

2. Gas-filling station

As mentioned in the previous section, a clean gas-filling station and fabrication process without impurities are important for a ^3He cell with long relaxation time. The first gas-filling station was constructed in 2018 at J-PARC. The schematic diagram of the gas-filling station is shown in Fig. 1. All the gas lines are wrapped with heaters for vacuum bake-out. An ultimate pressure of the gas-filling station was 5×10^{-8} Pa after baking at 150 °C for a few days.

Boron-free aluminosilicate GE180 glass is widely used for the glass cell of a ^3He spin filter due to a small wall relaxation effect and low permeability of ^3He [4,19–21]. Our fabrication process of ^3He cells is similar to those of NIST, JCMS, and ORNL [19–21]. A cylindrical-shaped glass cell made of GE180 is attached to a glassware made of Pyrex glass, which is referred to as “string”, (Fig. 2). The glass cell and the string are rinsed with neutral detergent, pure water, acetone, and alcohol before connecting to the gas-filling station. Rubidium and potassium ampoules

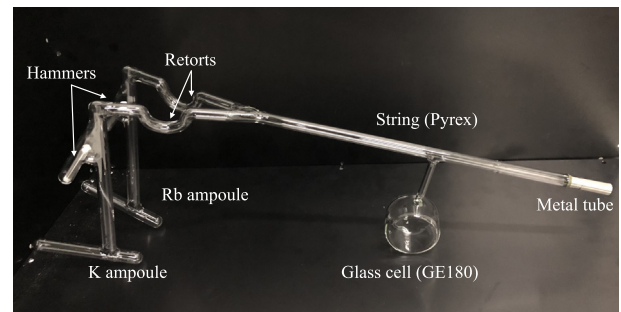


Fig. 2. The glassware used to fabricate a ^3He cell. The ampoules are put in the parts labeled as K ampoule and Rb ampoule in the figure. The glass cell made of GE180 is attached to the string made of Pyrex at the top of the glass cell. The string is connected to the pumping station using a Swagelok connector at the metal tube.



Fig. 3. ^3He cell (Sekichiku) fabricated at J-PARC.

Table 1

List of ^3He cells using rubidium and potassium with long relaxation times at J-PARC.

Name	Dimensions [mm]	^3He pressure [atm]	T_1 [h]
Chidori	$\phi 45 \times 75$	3.1	201
Karigane	$\phi 45 \times 75$	3.1	180
Sekichiku	$\phi 60 \times 60$	3.1	175
Hanabishi	$\phi 40 \times 90$	3.1	165

are also rinsed with acetone and alcohol. After connecting the string to the gas-filling station, the ampoules are put in the string without breaking the ampoules, and the string is sealed off with a gas torch. The ampoules are broken in a nitrogen atmosphere by dropping hammers made of an iron rod covered with glass. The string and the glass cell are baked out for a week at 200 °C and 400 °C, respectively. In the next step, rubidium and potassium are distilled to retorts at 200 °C and 220 °C, respectively, by using the heaters. After that, the glass tube parts containing ampoules and hammers are pulled off with a gas torch. The glass cell and the string are baked out again for one week, and then the alkali metals in the retorts are distilled to the glass cell. Finally, ^3He and N_2 gases purified using GC50 getters (SAES Getters) are filled to the glass cell, and the glass cell is sealed off by placing a gas torch at the Pyrex part. When encapsulating the gases above 1 atm, the glass cell is submerged in liquid nitrogen to keep the pressure inside the glass cell below 1 atm while sealing. So far, nine ^3He cells have been fabricated with ^3He pressures of 3.1 atm using the gas-filling station at J-PARC, and four of them have found to have spin relaxation times over 150 hours. A list of these ^3He cells and a photograph of Sekichiku cell are presented in Table 1 and Fig. 3, respectively.

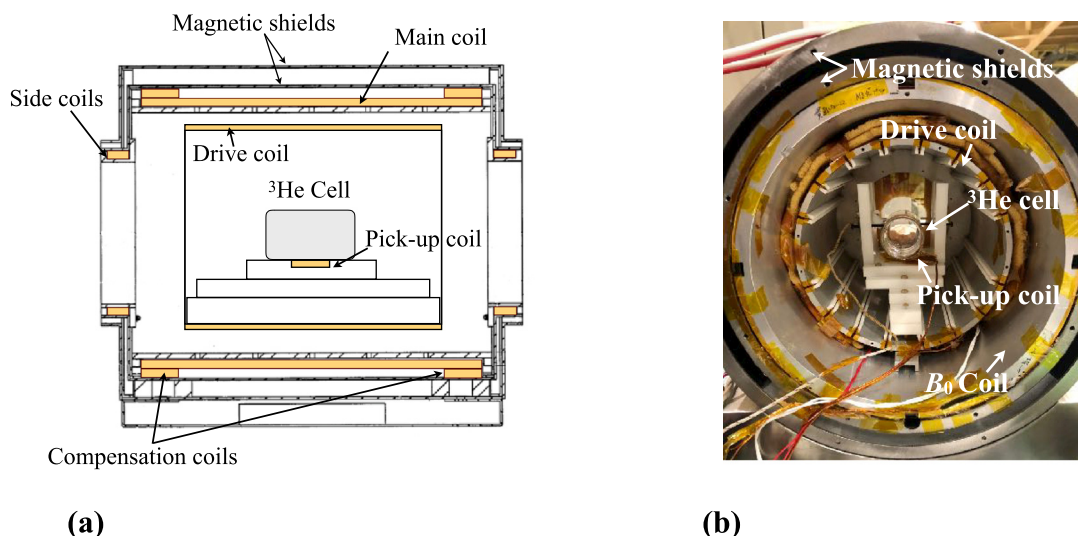


Fig. 4. (a) Cross-sectional view of the B_0 coil and the NMR coils. The magnetic shields are made from permalloy and its thicknesses are 2 mm. A static magnetic field is generated with the B_0 coil to keep ^3He polarization. The drive coil is used to generate an oscillating magnetic field to flip ^3He spins with AFP-NMR, and the spin flip signal is detected with the pick-up coil. (b) A photograph of the inside of the coil with the ^3He cell and the NMR coils.

3. NMR and pumping systems

3.1. B_0 coil and NMR system

To keep the ^3He polarization, a coil with double magnetic shields, referred to as B_0 coil, is mainly used for the laser pumping [22]. Fig. 4 shows a cross-sectional view of the B_0 coil and a photograph of its interior. The coil is consisting of a main solenoid, two compensation coils, and two side coils. The dimensions of the B_0 coil are 29.2 cm in diameter and 36.8 cm in length. The B_0 coil was designed using the finite element method to make the field gradient less than 5×10^{-4} /cm in a cylindrical region with a diameter of 10 cm and a length of 10 cm at the center of the coil. A magnetic field of 1.5 mT is produced to keep the ^3He polarization. Frequency-sweep Adiabatic Fast Passage (AFP)-NMR was employed to flip the ^3He spin as well as to evaluate the ^3He polarization [17]. A cosine winding is used as a drive coil to apply an oscillating magnetic field to the ^3He cell. The ^3He precession signal is detected by a pick-up coil located under the ^3He cell. The coils were installed inside the B_0 coil as shown in Fig. 4. The polarization loss due to spin flip was measured as 3.8×10^{-5} per flip. The ^3He cell is heated by electrical heating using a rubber heater with a power of 100 W wound around the drive coil during optical pumping.

3.2. Pumping station at MLF

We have set up a laboratory in the MLF beam hall to polarize ^3He cells. ^3He cells polarized at the laboratory can be provided to neutron beamlines in MLF. A pumping system using an air-cooling laser module with a laser power of 30 W [23] has been installed in the laboratory (Fig. 5). The dimensions of the pumping system are 60 cm \times 60 cm \times 40 cm, and it is used for both *in-situ* and *ex-situ* methods. Using a volume Bragg grating (VBG) element in the laser module, the center wavelength of the laser is tuned to 794.7 nm with a line width of 0.2 nm in Full Width at Half Maximum (FWHM). The laser light is circularly polarized using a quarter-wavelength plate installed in the module. The polarized laser light from the laser module is reflected by dielectric multi-layer coated quartz mirrors and irradiates the ^3He cell. The temperature at the cell surface is kept about 170 $^\circ\text{C}$ for a pure Rb cell and about 210 $^\circ\text{C}$ for a cell with rubidium and potassium during optical pumping. A ^3He polarization of $\sim 70\%$ is achieved with this laser system. A polarization measurement system using Electron Paramagnetic Resonance (EPR) [24] is also equipped.

3.3. Pumping system using a fiber laser

In order to achieve higher ^3He polarization, we have developed a separate pumping system using a fiber laser made by DILAS Diode Laser, Inc. (M1F4S22-794.7.[0,6]-100C-IS9), as shown in Fig. 6. The pumping system using this fiber laser has been constructed at another laboratory in the J-PARC research building, about 30 m away from the MLF experimental hall. This fiber laser is an air-cooled diode laser with a maximum output power of 110 W. The center wavelength is tuned to 794.7 nm with a line width of 0.4 nm in FWHM using a VBG element. An optical fiber with an internal diameter of 400 μm is connected to the laser shield box. An unpolarized laser beam from the fiber exit is split into two paths by a polarized beam splitter. The laser beams are circularly polarized with quarter-wavelength plates. The polarized laser beams irradiate both sides of a ^3He cell. The degree of circular polarization in each path was measured using a polarimeter as more than 98% at the ^3He cell position. The dimensions of the pumping system are 70 cm \times 70 cm \times 40 cm, and it can also be used for the *in-situ* method. The pumping system is currently used only for the *ex-situ* method, and will be installed in the laboratory at MLF in the near future after following a safety review.

4. Performance evaluation of the ^3He spin filter

We conducted an experiment at the NOBORU beamline [11,25], which is a test neutron beamline for general purpose, to evaluate the ^3He polarization for the Chidori cell, which has the longest T_1 as shown in Table 1. The ^3He cell was polarized using the pumping system with the fiber laser at J-PARC research building. Time evolution of the ^3He polarization $P_{\text{He}}(t)$ during optical pumping can be described as

$$P_{\text{He}}(t) = P_{\text{He}}^{\infty} (1 - \exp(-t/\tau)), \quad (1)$$

where t is the time from the start of optical pumping, P_{He}^{∞} is the ^3He polarization at $t \rightarrow \infty$, and τ is a polarization build-up time. The time evolution of the ^3He polarization in this experiment is shown in Fig. 7, and the build-up time τ was 3.5 hours. For reference, the typical build-up time is 6.5 hours when the pumping system with 30 W laser is used.

After the ^3He polarization was saturated, the ^3He cell, the B_0 coil, and the NMR system were transferred to the NOBORU beamline in MLF while applying a magnetic field to the ^3He cell using a battery to keep the ^3He polarization. Neutron transmission was measured using a Gas Electron Multiplier (GEM) detector [26] for the polarized and unpolarized ^3He cell as well as for an empty glass cell with the same

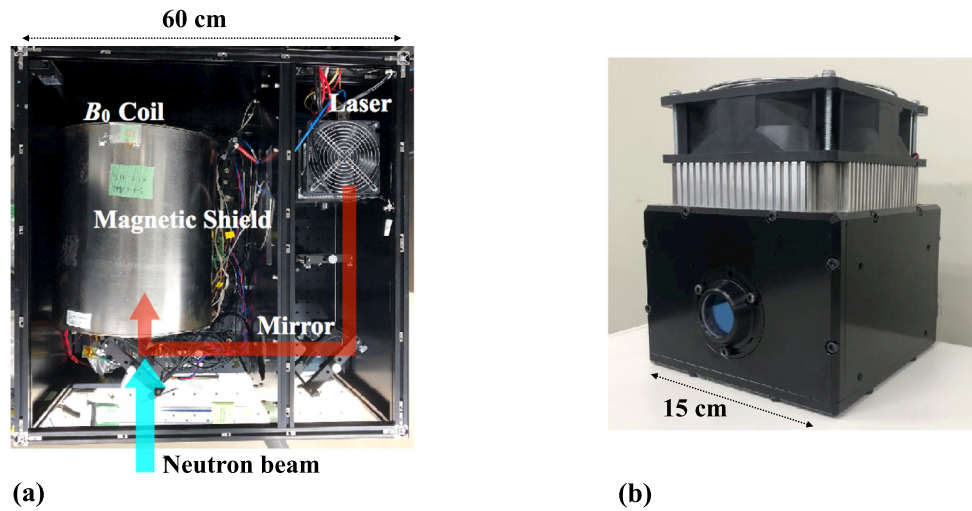


Fig. 5. (a) Compact pumping system which can be installed on a neutron beamline. The optical system has been constructed in the laser shield box. (b) Air cooling laser module. The laser power is 30 W.

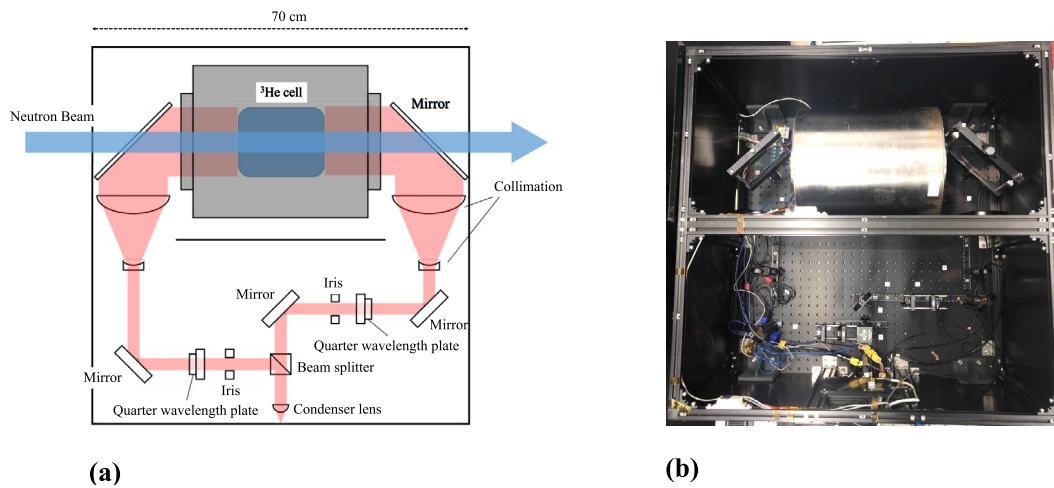


Fig. 6. (a) Schematic diagram of the pumping system using a fiber laser. (b) Photograph of the pumping system. The optical system was constructed in the laser shield box.

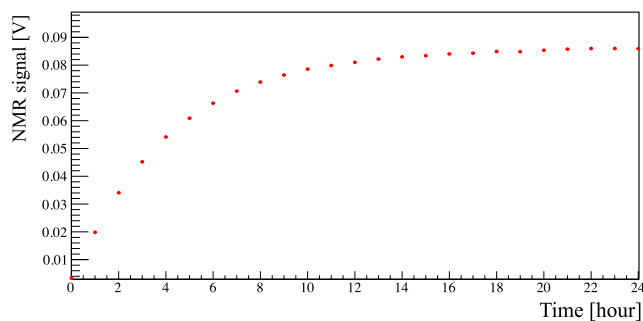


Fig. 7. Time evolution of the ^3He polarization measured by AFP-NMR during optical pumping. The ^3He polarization saturates at about 24 h after the start of the laser irradiation. The temperature at the surface of the cell was 205 °C during optical pumping.

dimensions. The number of the transmitted neutrons through the ^3He cell is described as [27]

$$N_0 T_{\text{cell}} \epsilon(\lambda) \exp(-\rho d \sigma(\lambda)) \cosh(P_{\text{He}} \rho d \sigma(\lambda)), \quad (2)$$

where, N_0 is the number of incident neutrons, T_{cell} is the transmission of the cell windows, $\epsilon(\lambda)$ is the detection efficiency of the GEM detector, P_{He} is ^3He polarization, ρ is the number density of ^3He nuclei, d is

the thickness of ^3He gas, $\sigma(\lambda)$ is the absorption cross section of the ^3He nucleus for unpolarized neutrons, and λ is the wavelength of incident neutrons, which is calculated from the time of flight. The ratio of transmitted neutrons for the polarized to unpolarized ^3He cell $N_{\text{pol}}/N_{\text{unpol}}$ is described as

$$\frac{N_{\text{pol}}}{N_{\text{unpol}}} = \cosh(P_{\text{He}} \rho d \sigma(\lambda)). \quad (3)$$

The product of the ^3He gas pressure and thickness ρd was obtained from the measurement of the ratio of transmitted neutrons for the unpolarized ^3He cell to the empty glass cell as 20.8 ± 0.1 atm-cm. The measured ratios of $N_{\text{pol}}/N_{\text{unpol}}$ are plotted against the neutron wavelength in Fig. 8 along with the curve of best fit using Eq. (3) with P_{He} as a free parameter. The ^3He polarization was measured to be $85.3 \pm 0.5\%$. The uncertainty comes mostly from the measured ρd .

The neutron wavelength dependences on the neutron polarization is obtained using the following equation:

$$P_n = \sqrt{1 - \frac{T_0^2}{T^2}}, \quad (4)$$

where P_n is the polarization of neutrons passed through the polarized ^3He cell, T is the transmission of polarized ^3He cell, and T_0 is the transmission of the unpolarized ^3He cell. The transmissions T_0 and T

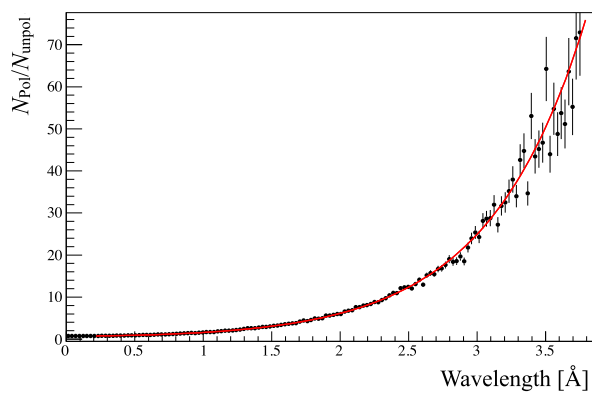


Fig. 8. Ratio of the numbers of transmitted neutrons for the polarized to the unpolarized ^3He cell. The red line shows the best fit. The error bars show the statistical uncertainty.

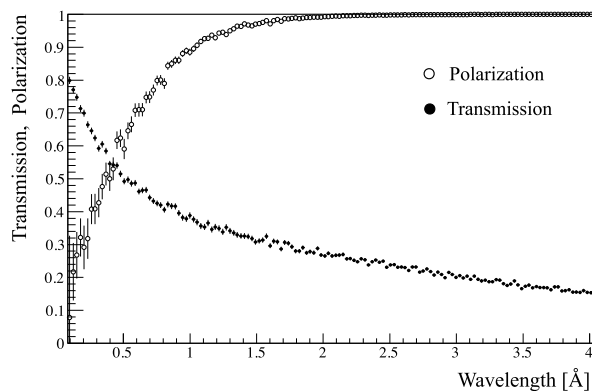


Fig. 9. Neutron wavelength dependences of neutron polarization (open circles) and neutron transmission (filled circles).

are described as

$$T_0 = \frac{N_{\text{unpol}}}{N_{\text{cell}}}, \quad T = \frac{N_{\text{pol}}}{N_{\text{cell}}}, \quad (5)$$

where N_{cell} is the transmitted neutrons for the empty glass cell. The neutron wavelength dependence of the neutron polarization and the neutron transmission are shown in Fig. 9. Each measurement took 15 min. The relaxation time T_1 on the beamline was 170 h, which was measured with AFP-NMR.

5. Trial and user experiments using ^3He spin filter

We have carried out several trial experiments with *in-situ* and *ex-situ* SEOP methods at neutron beamlines of MLF [28–31]. In 2017, the ^3He spin filter was first supplied to a user experiment. Although no user experiment was conducted in 2018, five user experiments were performed in 2019–2020. Neutron beamlines, where we conducted experiments using ^3He spin filters, are listed in Table 2. In this section, we report on experiments performed on typical beamlines at MLF using ^3He spin filters, including the results already reported. The experiments prior to 2020 were conducted using the pumping system with the 30 W laser, while the pumping system with the 110 W laser was used with *ex-situ* method in the experiments conducted in 2020. The pumping system with the 30 W laser used for the *in-situ* experiments described in this section will also be replaced by that with the 110 W laser in the near future.

5.1. BL10 NOBORU

A trial experiment for magnetic imaging was carried out [28] with *in-situ* and *ex-situ* systems using ^3He cell with pure Rb and the 30 W

Table 2

List of neutron beamlines, where experiments were conducted with ^3He spin filters.

No.	Name	Purpose of beamline	Type of ^3He spin filter
04	ANNRI	Nuclear reaction	<i>ex-situ</i>
05	NOP	Fundamental physics	<i>ex-situ</i>
06	VIN ROSE	Spin echo	<i>ex-situ</i>
10	NOBORU	General purpose	<i>ex-, in-situ</i>
15	TAIKAN	SANS	<i>ex-, in-situ</i>
17	SHARAKU	Reflectometer	<i>in-situ</i>
18	SENJU	Single crystal diffraction	<i>ex-situ</i>
22	RADEN	Pulsed neutron imaging	<i>in-situ</i>

laser as shown in Fig. 5. The *in-situ* system with a ^3He cell of 17 atm-cm and the *ex-situ* system with ^3He cell of 11 atm-cm were used as a neutron spin polarizer and analyzer, respectively. The dimensions of the ^3He cells were 35 mm in diameter and 55 mm in length for both. A coil with a 0.35 mm-thick ring-shaped magnetic steel core was placed between the polarizer and the analyzer as a sample. A two-dimensional RPMT neutron detector [32] was set after the analyzer. The magnetic field applied to the sample was controlled with the current sent to the coil. Two-dimensional images of the neutron polarization were obtained, and obvious differences in the spin rotations of neutrons were observed as the wavelength dependence of the neutron polarization in two current conditions, $I = 1.42$ A and $I = 0$ A. This result indicates that it is possible to perform an experiment for magnetic imaging with ^3He spin filters at J-PARC.

Furthermore, a study for magnetic atomic resolution holography with polarized neutrons using a ^3He spin filter is currently underway. Two user experiments for this study were conducted at NOBORU in 2019–2020.

5.2. BL04 ANNRI

BL04 ANNRI [11,33] is a neutron beamline used for studies of nuclear science such as nuclear data for nuclear technology, astrophysics, and quantitative analyses. A germanium detector assembly using 22 high-quality germanium crystals is installed on this beamline to measure the (n, γ) reaction. Measurements of the (n, γ) reaction for the study of fundamental symmetry violation in nuclei have been conducted as a user experiment using a ^3He spin filter [34–36]. A ^3He spin filter with the *ex-situ* system was used to polarize epithermal neutrons to measure the (n, γ) reaction with polarized neutrons at the 0.74 eV neutron resonance of ^{139}La . A lanthanum sample was placed at the center of the germanium detector assembly, and the emitted γ -rays from the sample were detected by the germanium detectors. The ^3He cell, solenoid, and compensation coils with a magnetic shield made of permalloy were installed between the detector and a neutron collimator (Fig. 10). The magnetic field used to keep the ^3He polarization was perpendicular to the neutron beam in order to obtain a vertically polarized neutron beam. The magnetic shield had two openings in the path of the neutron beam to avoid a decay of the neutron polarization. Neutron polarization was held by the magnetic field produced with the guide magnets to a lanthanum target. A ^3He cell using pure Rb with a pressure thickness of 19.3 atm-cm and dimensions of 50 mm diameter and 70 mm length was used. The ^3He cell was polarized with the 30 W laser system at MLF, and the ^3He polarization at the start of the measurement was typically 60%. The relaxation time of the ^3He polarization was 130 h. In 2019, as a result of three separate measurements using the ^3He spin filter, a significant spin-dependent angular distribution of γ -rays from the 0.74 eV neutron resonance of ^{139}La was observed [36].

Furthermore, a (n, γ) reaction measurement using polarized epithermal neutrons to obtain nuclear data for reactor science is being planned, and the development of a ^3He cell for epithermal neutrons is an ongoing process.

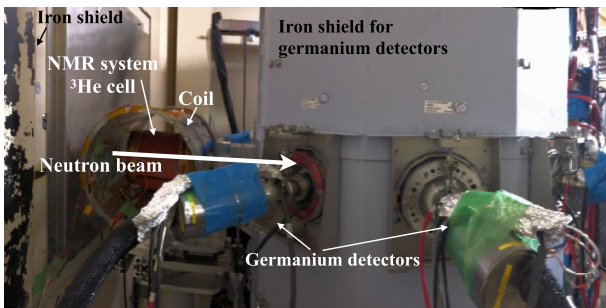


Fig. 10. Photograph of the *ex-situ* system installed at ANNRI beamline. The coil wrapped with permalloy sheets, NMR system, and ^3He cell were placed between the iron shields.

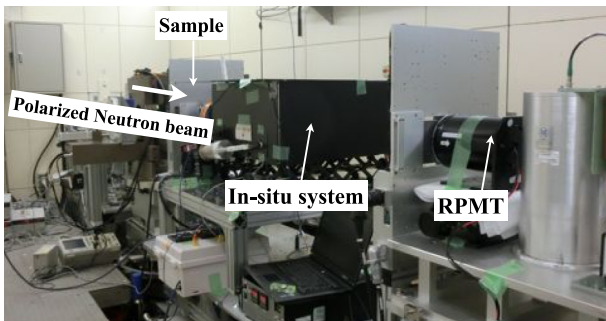


Fig. 11. Photograph of the *in-situ* system installed at SHARAKU beamline.

5.3. BL17 SHARAKU

A trial experiment using the *in-situ* system on a beamline for neutron reflectometry: BL17 SHARAKU [11,37] was performed [28]. The SHARAKU beamline has a neutron polarizer and an analyzer consisting of supermirrors. Some of the supermirrors are stacked vertically in order to cover a wide area, which analyze the scattered neutron spins with wavelengths longer than 0.2 nm. However, some neutrons, which transmitted through the supermirrors are scattered by the supermirrors, and spacial distribution of neutrons reflected from a sample is disturbed. A ^3He spin filter was used as the analyzer instead of the supermirrors to reduce uncertainty produced with the supermirrors (Fig. 11). The ^3He cell using pure Rb with a pressure thickness of 11.1 atm-cm, diameter of 35 mm, and length of 55 mm was used. A Fe-Cr multi-layered thin film was used as a sample, and off-specular measurements were performed using the *in-situ* system. The ^3He polarization was 60% during the experiment. The *in-situ* system successfully demonstrated and worked stably over 4 days of the experimental time. The sample had been measured previously at ISIS, and we obtained the identical results on SHARAKU. This satisfactory implies that our *in-situ* system worked as designed.

5.4. BL06 VIN ROSE

BL06 VIN ROSE [11,38] is equipped with a modulated intensity by zero effort (MIEZE) spectrometer to analyze the slow dynamics of condensed matter by measuring the intermediate scattering function. Experiments using a combination of a MIEZE spectrometer and the time of flight (TOF-MIEZE [39]) method have been carried out, and additionally, polarization analysis of scattered neutrons by a sample is planned in order to study spin dynamics by installing a spin analyzer. We installed a ^3He cell after the sample position as an analyzer to confirm that the ^3He spin filter can be used for the MIEZE spectrometer [29]. The MIEZE signal was measured with and without the ^3He cell, and no difference was observed in the signal contrast. This result

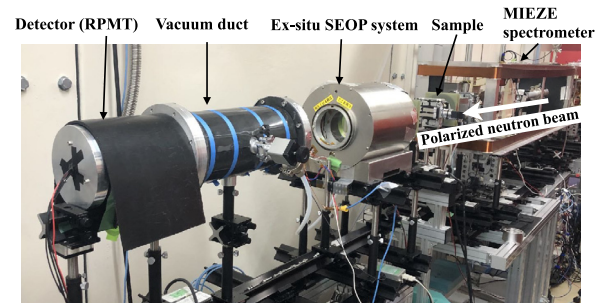


Fig. 12. Photograph of the *ex-situ* system installed on the VIN ROSE beamline. The polarized neutron beam was scattered by the sample, and the spin of the scattered neutron was analyzed with the ^3He spin filter. The vacuum duct was wrapped with boron sheets to reduce background neutrons.

implies that diffusion of ^3He gas does not affect the MIEZE signal and a ^3He spin filter can be used as a second analyzer in the MIEZE spectrometer. Furthermore, in 2020, an experiment to measure the spin dynamics of ferrofluid was conducted by using a ^3He spin filter as shown in Fig. 12. The ^3He polarization of 72% and relaxation time of 120 h was achieved on this beamline in the experiment, and the analysis is currently underway.

5.5. BL15 TAIKAN

BL15 TAIKAN [11,40] is a small and wide angle neutron scattering instrument. We conducted a trial experiment to measure coherent and incoherent scattering cross sections from hydrogen contained material by analyzing scattered neutrons using the *in-situ* system [30]. Supermirrors were installed upstream of the beamline as a neutron polarizer, and the ^3He spin filter was used as a spin analyzer. Silver behenate was used as the sample, and the ^3He cell with pure Rb was placed 120 mm downstream of the sample. The diameter of ^3He cell was 35 mm, the length was 55 mm, and the gas pressure was 3 atm. A stable ^3He polarization of 68% was kept during the experiment. Scattered neutrons reaching a part of the small-angle detector bank were spin analyzed. The coherent and incoherent scattering components were successfully separated by analyzing the scattered neutrons. Although the trial experiment was successful, the solid angle of the ^3He cell was inadequate and the pressure-thickness was not suitable for the neutron wavelength used in TAIKAN. In order to cover a larger solid angle, an *ex-situ* SEOP system using a larger ^3He cell, which is to be placed 10 mm downstream of the sample, is being prepared (Fig. 13). A new ^3He cell with rubidium and potassium has been fabricated with a diameter of 60 mm, a thickness of 40 mm, and a gas pressure of 1.5 atm for cold neutrons. More trial and user experiments will be carried out in 2020.

5.6. BL21 NOVA

Experiments using ^3He spin filters at BL21 NOVA [11,41], which is a total diffractometer, are also planned. NOVA is equipped with large solid angle neutron detectors. The detectors and a sample are installed in a large vacuum chamber. We are planning to install two *ex-situ* systems upstream of the beamline and around the sample to perform polarized neutron scattering experiments. The schematic view of the *ex-situ* systems at NOVA is shown in Fig. 14. The *ex-situ* system for the polarizer is installed at the upstream of the beamline, and the neutron polarization is held by the magnetic field produced with a guide coil to the sample. The scattered neutrons by the sample are analyzed with two ^3He cells. This *ex-situ* system for the analyzer which can be installed to the vacuum chamber is developed as shown in Fig. 15. A sample is suspended from the top of the vacuum flange. The neutrons scattered by the sample are analyzed with two ^3He cells placed at a distance of 20 mm from the sample. The stage for the ^3He cells can be rotated

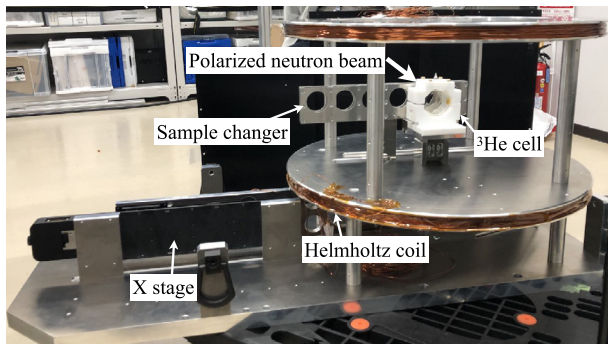


Fig. 13. Photograph of an *ex-situ* SEOP system for a SANS experiment at TAIKAN beamline. The *ex-situ* system is equipped with a sample changer which can be used to place up to six samples. The sample position is moved with the X-stage. The ^3He polarization is kept by the Helmholtz coil. The ^3He cell with the dimensions of $\phi 60 \times 40$ mm is placed at 10 mm downstream of the sample to cover a large solid angle.

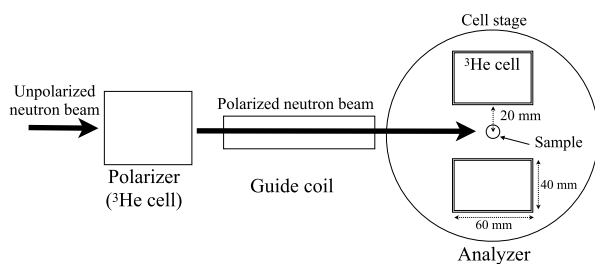


Fig. 14. Schematic view from the top of the *ex-situ* systems at NOVA. A neutron beam is polarized by the ^3He cell installed upstream of the beamline, and the neutron polarization is held by the magnetic field produced with the guide coil to the sample. The scattered neutrons by the sample are analyzed with two ^3He cells.

so that the scattered neutrons of any angle can be analyzed. The ^3He polarization is kept using the Helmholtz coil that produces 40 W of heat. The heat from the coil flows to the flange through the 30 mm thick aluminum plates connected to the coil, and the flange is air-cooled. A thermal simulation of the *ex-situ* system in a vacuum was performed. As a result of the simulation, the maximum temperature was determined to be 76 °C at the bottom of the coil, which does not cause a problem. A test experiment using the *ex-situ* system will be carried out in 2020.

6. Conclusion

Development of ^3He spin filters for user experiments are being carried out at MLF at J-PARC. The *in-situ* system using a 30 W laser has been developed, and several trial experiments using ^3He cells with pure Rb have been conducted. Recently, a gas-filling station has been constructed at J-PARC and several high quality ^3He cells using rubidium and potassium were fabricated. Additionally, the pumping system using a 110 W fiber laser was developed. High ^3He polarization of 85% was achieved on a neutron beamline using the ^3He cell with rubidium and potassium and the pumping system using the fiber laser. The first user experiment utilizing a ^3He spin filter took place in 2017, and five user experiments were conducted in 2019–2020. These experiments are beginning to yield scientific results. Preparations are underway for further user experiments using ^3He spin filters.

CRediT authorship contribution statement

T. Okudaira: Methodology, Conceptualization, Investigation, Software, Formal analysis, Writing - original draft, Funding acquisition. **T. Oku:** Methodology, Conceptualization, Investigation, Supervision, Writing - review & editing, Funding acquisition. **T. Ino:** Methodology,

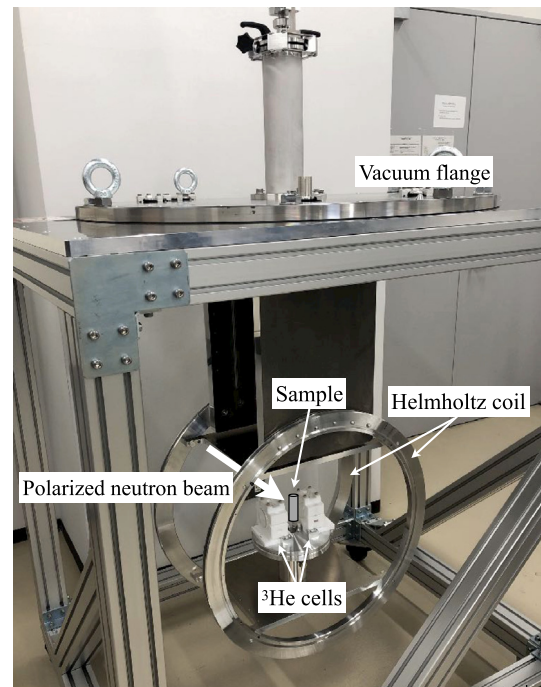


Fig. 15. Photograph of the *ex-situ* system for the analyzer. The sample, the ^3He cells, and the Helmholtz coil are suspended from the vacuum flange. The ^3He polarization is kept with the Helmholtz coil.

Conceptualization, Investigation, Resources, Writing - review & editing. **H. Hayashida:** Methodology, Investigation. **H. Kira:** Methodology, Investigation. **K. Sakai:** Methodology, Investigation. **K. Hiroi:** Investigation, Resources. **S. Takahashi:** Investigation. **K. Aizawa:** Supervision, Writing - review & editing. **H. Endo:** Investigation, Resources. **S. Endo:** Investigation, Resources. **M. Hino:** Investigation, Resources. **K. Hirota:** Methodology, Investigation. **T. Honda:** Methodology, Investigation, Resources. **K. Ikeda:** Methodology, Resources. **K. Kakurai:** Conceptualization. **W. Kambara:** Investigation. **M. Kitaguchi:** Methodology, Investigation. **T. Oda:** Methodology, Investigation, Resources. **H. Ohshita:** Methodology, Investigation, Resources. **T. Otomo:** Methodology, Resources. **H.M. Shimizu:** Methodology. **T. Shinohara:** Methodology. **J. Suzuki:** Methodology. **T. Yamamoto:** Methodology, Investigation.

Declaration of competing interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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References

- [1] E. Babcock, S. Zahir, T. Tobias, S. Denis, S. Johann, F. Artem, M. Stefan, P. Patrick, R. Aurel, I. Alexander, *J. Phys. Conf. Ser.* 711 (2016) 012008.
- [2] I. Dhiman, R. Ziesche, T. Wang, H. Bilheux, L. Santodonato, X. Tong, C.Y. Jiang, I. Manke, W. Treimer, T. Chatterji, N. Kardjilov, *Rev. Sci. Instrum.* 88 (9) (2017) 1–7, <http://dx.doi.org/10.1063/1.5001525>, <http://dx.doi.org/10.1063/1.5001525>.
- [3] W.C. Chen, J.G. Barker, R. Jones, K.L. Krycka, S.M. Watson, C. Gagnon, T. Perevozchivoka, P. Butler, T.R. Gentile, *J. Phys. Conf. Ser.* 862 (1) (2017) <http://dx.doi.org/10.1088/1742-6596/862/1/012004>.
- [4] E. Lelievre-Berna, *Phys. B* 397 (1–2) (2007) 162–167.
- [5] W.T.H. Lee, T. D'Adam, *Neutron News* 27 (2) (2016) 35–38, <http://dx.doi.org/10.1080/10448632.2016.1166912>.
- [6] C.J. Beecham, S. Boag, C.D. Frost, T.J. McKetterick, J.R. Stewart, K.H. Andersen, P.M. Bentley, D. Jullien, *Physica B* 406 (12) (2011) 2429–2432, <http://dx.doi.org/10.1016/j.physb.2010.11.054>.
- [7] M. Bouchiat, T.R. Carver, C.M. Varnum, *Phys. Rev. Lett.* 5 (1960) 373.
- [8] F.D. Colegrove, L.D. Scheerer, G.K. Walters, *Phys. Rev.* 132 (1963) 2561.
- [9] E. Babcock, I. Nelson, S. Kadlecck, B. Driehuis, L.W. Anderson, F.W. Hersman, T.G. Walker, *Phys. Rev. Lett.* 91 (2003) 123003.
- [10] T.R. Gentile, P.J. Nacher, B. Saam, T.G. Walker, *Rev. Modern Phys.* 89 (2017) 045004.
- [11] K. Nakajima, Y. Kawakita, S. Itoh, J. Abe, K. Aizawa, H. Aoki, H. Endo, M. Fujita, K. Funakoshi, W. Gong, M. Harada, S. Harjo, T. Hattori, M. Hino, T. Honda, A. Hoshikawa, K. Ikeda, T. Ino, T. Ishigaki, Y. Ishikawa, H. Iwase, T. Kai, R. Kajimoto, T. Kamiyama, N. Kaneko, D. Kawana, S.O. Kawamura, T. Kawasaki, A. Kimura, R. Kiyonagi, K. Kojima, K. Kusaka, S. Lee, S. Machida, T. Masuda, K. Mishima, K. Mitamura, M. Nakamura, S. Nakamura, A. Nakao, T. Oda, T. Ohhara, K. Ohishi, H. Ohshita, K. Oikawa, T. Otomo, A.S. Furukawa, K. Shibata, T. Shinohara, K. Soyama, J. Suzuki, K. Suzuya, A. Takahara, S. Takata, M. Takeda, Y. Toh, S. Torii, N. Torikai, N.L. Yamada, T. Yamada, D. Yamazaki, T. Yokoo, M. Yonemura, H. Yoshizawa, *Quantum Beam Sci.* 1 (2017) 9.
- [12] H. Sato, Y. Masuda, K. Sakai, M. Doi, K. Asahi, Z.J. Zheng, Y. Mori, P.P.J. Delheij, Y. Matsuda, M. Iinuma, H.M. Shimizu, A. Masaike, *Hyperfine Interact.* 84 (1994) 205–209.
- [13] T. Ino, S. Muto, Y. Masuda, G. Kim, V.R. Skoy, *Physica B* 356 (2005) 109.
- [14] T. Ino, M. Nakamura, T. Oku, T. Shinohara, J. Suzuki, K. Ohoyama, H. Hiraka, *Physica B* 404 (2009) 2667.
- [15] T. Ino, M. Ohkawara, K. Ohoyama, T. Yokoo, S. Itoh, Y. Nambu, M. Fujita, H. Kira, H. Hayashida, K. Hiroi, K. Sakai, T. Oku, K. Kakurai, *J. Phys. Conf. Ser.* 711 (2016) 012011.
- [16] T. Ino, K. Ohoyama, T. Yokoo, S. Itoh, M. Ohkawara, H.H. H. Kira, K. Sakai, K. Hiroi, T. Oku, K. Kakurai, L.J. Chang, *J. Phys. Conf. Ser.* 862 (2017) 012011.
- [17] T. Ino, *JPS Conf. Proc.* 22 (2018) 01101.
- [18] R. Kajimoto, T. Yokoo, M. Nakamura, Y. Kawakita, M. Matsuura, H. Endo, H. Seto, S. Itoh, K. Nakajima, S. Ohira-Kawamura, *Physica B* 562 (2019) 148–154.
- [19] C. Jiang, X. Tong, D. Brown, W. Lee, H. Ambaye, J. Craig, L. Crow, H. Culbertson, R. Goyette, M. Graves-Brook, M. Hagen, B. Kadron, V. Lauter, L. McCollum, J. Robertson, A. B. Winn, Vandegrift, *Physics Procedia* 42 (2013) 191–199.
- [20] W.C. Chen, T.R. Gentile, C.B. Fu, S. Watson, G.L. Jones, J.W. Mclver, D.R. Rich, *J. Phys. Conf. Ser.* 294 (2011) 012003.
- [21] Z. Salhi, E. Babcock, P. Pistel, A. Ioffe, *J. Phys. Conf. Ser.* 528 (2014) 012015.
- [22] H. Kira, Y. Sakaguchi, J. Suzuki, T. Oku, M. Nakamura, M. Arai, K. Kakurai, Y. Endo, Y. Arimoto, T. Ino, H. Shimizu, T. Kamiyama, K. Ohoyama, H. Hiraka, K. Tsutsumi, K. Yamada, L.-J. Chang, *Physics Procedia* 42 (2013) 200.
- [23] T. Oku, H. Hayashida, H. Kira, K. Sakai, K. Hiroi, T. Shinohara, Y. Sakaguchi, T. Ino, K. Ohoyama, L.-J. Chang, M. Nakamura, J. Suzuki, K. Aizawa, M. Arai, Y. Endoh, K. Kakurai, *JPS Conf. Proc.* 8 (2015) 036009.
- [24] M.V. Romalis, G.D. Cates, *Phys. Rev. A* 58 (1998) 3004.
- [25] K. Oikawa, F. Maekawa, M. Harada, T. Kai, S. Meigo, Y. Kasugai, M. Ooi, K. Sakai, M. Teshigawara, S. Hasegawa, M. Futakawa, Y. Ikeda, N. Watanabe, *Nucl. Instrum. Methods Phys. Res. A* 589 (2008) 310–31.
- [26] S. Uno, T. Uchida, M. Sekimoto, T. Murakami, K. Miyama, M. Shoji, E. Nakano, T. Koike, K. Morita, H. Satoh, T. Kamiyama, Y. Kiyonagi, *Physics Procedia* 26 (2012) 142–152.
- [27] K.P. Coulter, T.E. Chupp, A.B. McDonald, C.D. Bowman, J.D. Bowman, J.J. Szymanski, V. Yuan, G.D. Cates, D.R. Benton, E.D. Earle, *Nucl. Instrum. Methods Phys. Res. A* 288 (1990) 463–466.
- [28] H. Hayashida, T. Oku, H. Kira, K. Sakai, K. Hiroi, T. Ino, T. Shinohara, T. Imagawa, M. Ohkawara, K. Ohoyama, K. Kakurai, M. Takeda, D. Yamazaki, K. Oikawa, M. Harada, N. Miyata, K. Akutsu, M. Mizusawa, J.D. Parker, Y. Matsumoto, S. Zhang, J. Suzuki, K. Soyama, K. Aizawa, M. Arai, *J. Phys. Conf. Ser.* 711 (2016) 012007.
- [29] H. Hayashida, M. Hino, H. Endo, T. Oku, T. Okudaira, K. Sakai, T. Oda, *J. Phys. Conf. Ser.* 1316 (2019) 012013.
- [30] H. Kira, H. Hayashida, H. Iwase, K. Ohishi, J. Suzuki, T. Oku, K. Sakai, K. Hiroi, S. Takata, T. Ino, K. Ohoyama, M. Ohkawara, T. Shinohara, K. Kakurai, K. Aizawa, M. Arai, *JPS Conf. Proc.* 8 (2015) 036008.
- [31] K. Sakai, T. Oku, H. Hayashida, H. Kira, T. Shinohara, K. Oikawa, M. Harada, K. Kakurai, K. Aizawa, M. Arai, Y. Sakaguchi, J. Suzuki, T. Ino, K. Ohoyama, *J. Phys. Conf. Ser.* 528 (2014) 012016.
- [32] K. Hirota, T. Shinohara, K. Ikeda, K. Mishima, T. Adachi, T. Morishima, S. Satoh, T. Oku, S. Yamada, H. Sasao, J. ichi Suzuki, H.M. Shimizu, *Phys. Chem. Chem. Phys.* 7 (2008) 1836–1838.
- [33] M. Igashira, Y. Kiyonagi, M. Oshima, *Nucl. Instrum. Methods Phys. Res. A* 600 (2009) 332.
- [34] T. Yamamoto, H.M. Shimizu, M. Kitaguchi, K. Hirota, T. Okudaira, C.C. Haddock, N. Oi, I. Ito, S. Endo, S. Takada, J. Koga, T. Yoshioka, T. Ino, K. Asahi, T. Momose, T. Iwata, K. Sakai, T. Oku, A. Kimura, M. Hino, T. Shima, Y. Yamagata, *JPS Conf. Proc.* 22 (2018) 011018.
- [35] T. Yamamoto, H.M. Shimizu, M. Kitaguchi, K. Hirota, C.C. Haddock, S. Endo, K. Ishizaki, T. Sato, S. Takada, J. Koga, T. Yoshioka, T. Ino, K. Asahi, T. Momose, T. Iwata, K. Sakai, T. Okudaira, T. Oku, A. Kimura, M. Hino, T. Shima, Y. Yamagata, *EPJ Web Conf.* 219 (2019) 09002.
- [36] T. Yamamoto, T. Okudaira, S. Endo, H. Fujioka, K. Hirota, T. Ino, K. Ishizaki, A. Kimura, M. Kitaguchi, J. Koga, S. Makise, Y. Niinomi, T. Oku, K. Sakai, T. Shima, H.M. Shimizu, S. Takada, Y. Tani, H. Yoshikawa, T. Yoshioka, *Phys. Rev. C* 101 (2020) 064624.
- [37] M. Takeda, D. Yamazaki, K. Soyama, R. Maruyama, H. Hayashida, H. Asaoka, T. Yamazaki, M. Kubota, K. Aizawa, M. Arai, Y. Inamura, T. Itoh, K. Kaneko, T. Nakamura, T. Nakatani, K. Oikawa, T. Ohara, Y. Sakaguchi, K. Sakasai, T. Shinohara, J. Suzuki, K. Suzuya, I. Tamura, K. Toh, H. Yamagishi, N. Yoshida, T. Hirano, *Chinese J. Phys.* 50 (2012) 161–170.
- [38] M. Hino, T. Oda, N.L. Yamada, H. Endo, H. Seto, M. Kitaguchi, M. Harada, Y. Kawabata, *J. Nucl. Sci. Technol.* 54 (2017) 1223–1232.
- [39] T. Oda, M. Hino, M. Kitaguchi, P. Geltenbort, Y. Kawabata, *Rev. Sci. Instrum.* 87 (2016) 105124.
- [40] T. Shinohara, S. Takata, J. Suzuki, T. Oku, K. Suzuya, K. Aizawa, M. Arai, T. Otomo, M. Sugiyama, *Nucl. Instrum. Methods Phys. Res. A* 600 (2009) 111.
- [41] T. Otomo, K. Suzuya, M. Misawa, N. Kaneko, H. Ohshita, K. Ikeda, M. Tsubota, T. Seya, T. Fukunaga, K. Itoh, M. Sugiyama, K. Mori, Y. Kameda, T. Yamaguchi, K. Yoshida, K. Maruyama, Y. Kawakita, S. Shamoto, K. Kodama, S. Takata, S. Satoh, S. Muto, T. Ino, H.M. Shimizu, T. Kamiyama, S. Ikeda, S. Itoh, Y. Yasu, K. Nakayoshi, H. Sendai, S. Uno, M. Tanaka, K. Ueno, *KENS Report* 17, 2009–2010, p. 28.