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1 1 **Microscopic characteristics of partially saturated dense sand and their link to macroscopic responses**

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4 2 **under triaxial compression conditions**

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52 16 Keywords: partially saturated sand, triaxial compression test, X-ray micro tomography, liquid bridge, principal

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55 17 curvature of air-water interface, water retention state

1 18 **Abstract**

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4 19 This paper presents a set of triaxial compression tests on partially saturated dense sands to clarify the microscopic
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7 20 characteristics and their link to the macroscopic responses. Constant suction tests (CS tests) and constant water
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10 21 content tests (CW tests) are conducted under low confining pressure to observe microscopic and macroscopic
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13 22 behaviors of the sands associated with dilative shear bands. X-ray micro tomography and image analysis
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16 23 techniques are applied to investigate the continuity as a defined index to evaluate the morphology of the pore water,
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20 24 the number of liquid bridges and the principal curvature of the air-water interface. The number of liquid bridges
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23 25 decreases for both the CS and CW tests, particularly during the strain softening process, while it decreases greater
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26 26 in the CW test than in the CS test. The curvature of the air-water interface remains at almost the same value under
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29 27 the CS test, while it tends to decrease under the CW test. The tendency of decreasing curvature corresponds to that
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32 28 of decreasing suction in the CW test. The peak deviator stress is higher in the CS test than in the CW test when the
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35 29 pore water is initially discontinuous, whereas it is identical between the two tests when the pore water is initially
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39 30 continuous. The residual stress is lower in the CW test than in the CS test, independent of the initial water retention
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42 31 states. The macroscopic responses at the different initial water retention states are qualitatively identical between
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45 32 poorly-graded sand and well-graded sand.
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1 35 **1 Introduction**

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4 36 Water retention states in partially saturated soils are varied with degrees of saturation and are classified into
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7 37 pendular, funicular and capillary states [36]. In the pendular state, liquid bridges exist between soil particles due
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10 38 to the surface tension of water and the hydrophilic property of the soil, where suction (difference between pore air
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13 39 pressure and pore water pressure) enhances the shear stiffness and the shear strength of partially saturated soils
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16 40 (e.g., [8-12, 30, 37, 40]). In order to investigate mechanical characteristics of partially saturated soils, triaxial
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20 41 compression tests have been conducted under a drained condition for water (e.g., [4, 13, 31]) and an undrained
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23 42 condition for water (e.g., [2, 31, 35, 43]). These studies have revealed that higher suction levels provide the higher
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26 43 shear strength, while suction tends to decrease with shearing under the undrained condition for water, particularly
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29 44 in the case of higher initial levels of suction. It is also known that partially saturated soil exhibits more brittle
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32 45 modes of failure associated with clearer shear bands than fully saturated soil (e.g., [5, 16, 17]). The effect of suction
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35 46 on the mechanical behavior of partially saturated soil depends on the water retention states (e.g., [21, 26, 41]), and
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39 47 hence the effect changes due to the variation in water retention states as shearing progresses. These macroscopic
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42 48 responses of partially saturated soil under triaxial compression must be interpreted by three-phase microstructural
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45 49 changes; however, the microscopic interpretation has not been sufficiently done. This is because observation of
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48 50 the three-phase microstructural changes associated with the development of shear bands in partially saturated soil
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51 51 is difficult using conventional techniques.

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54 52 An X-ray tomography technique has been widely used to observe the density change in geomaterials since the
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58 53 pioneering study by [7]. In recent years, an X-ray micro tomography technique has enabled discussions on the
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1 54 water retention behavior and the three-phase structural changes of partially saturated granular materials from a
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4 55 microscopic viewpoint. It was revealed that, at low saturation degrees, a large number of water units having very
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7 56 small volume (liquid bridges) exist and the number of liquid bridges decreases with increasing water saturation
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10 57 [42]. Microscopic characteristics such as the air-water interfacial contact angle [29], pore water distribution
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13 58 changes during drying process [20], and capillary collapse due to increasing water content [3] were clearly
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16 59 observed. Specific air-water interfacial areas in partially saturated glass bead samples and their influence on the
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20 60 material strength were quantitatively discussed [39]. Further studies are still needed, however, in order to clarify
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23 61 the link of grain- and pore-scale characteristics to the macroscopic responses for various kinds of soils under partial
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26 62 saturation conditions.

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29 63 The objective of the present study is to clarify the microscopic characteristics of partially saturated sands under
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32 64 triaxial compression and their link to the macroscopic responses. Triaxial compression tests on partially saturated
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36 65 sands are conducted under drained conditions for air and water, and a drained condition for air and an undrained
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39 66 condition for water. Dilative behaviors of dense sand specimens under low confining pressure are investigated,
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42 67 considering partially saturated soils near the ground surface such as well-compacted embankments. Poorly-graded
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45 68 silica sand, which is advantageous for high spatial resolution image acquisition using an X-ray micro tomography,
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48 69 is used to investigate the three-phase microstructural changes associated with the development of shear bands.
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51 70 Preliminarily, research group of the authors revealed that void ratio increased and the degree of saturation
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55 71 decreased as shear bands develop [19, 22], and that the tendency for decreasing the number of grain contacts and
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58 72 water menisci was closely related to the macroscopic strain softening behavior [23]. In the present study, variations
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1 73 in the morphology of the pore water, the number of liquid bridges and the principal curvature of the air-water
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4 74 interface with the development of shear bands are investigated by image analysis techniques using the tomographic
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7 75 volumes. For the poorly-graded sand, it is difficult to measure the suction change during triaxial compression test
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10 76 because the change is quite small due to its low water retention capability. Hence, well-graded sand, which shows
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13 77 much higher water retention capability than the poorly-graded sand, is also tested and suction is measured during
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16 78 triaxial compression. Through comparison between the measured macroscopic suction changes and the variation
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20 79 in the microscopic principal curvature of the air-water interface, progressive changes of suction and morphological
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23 80 transitions of pore water during triaxial compression are demonstrated. The influence of the quantified microscopic
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26 81 characteristics on the deviator stress-axial strain relationship for partially saturated sand with different initial
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29 82 degrees of saturation, under triaxial compression with the different drained conditions for water, is discussed.
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36 84 **2 Experimental Setup**

39 85 **2.1 Material and specimen preparation**

42 86 **2.1.1 Poorly-graded sand**

45 87 Poorly-graded sand, whose particles have a diameter of several hundred μm , are generally used for the
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48 88 visualization of the microstructure using X-ray micro tomography due to the limitation of the spatial resolution.
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51 89 Kido and Higo (2017) [22] confirmed that sufficiently high accuracy of the image analyses was achieved when
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55 90 the diameter of sand particles corresponded to 27 voxels. For the experimental setup in the present study, the
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58 91 available voxel size is 12 μm . With this resolution, sand particles with a D_{50} larger than 324 μm are required for
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1 92 the subsequent image analyses. Thus, silica sand No. 5, whose D_{50} is 456 μm , was used in the present study.
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4 93 **Figures 1a** and **2** present a SEM image of the silica sand and its grain size distribution curve, respectively. The
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7 94 physical properties of silica sand include a soil density of 2.64 g/cm^3 , a maximum void ratio of 1.013, a minimum
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10 95 void ratio of 0.694, a uniformity coefficient of 1.3 and a fines content of 0.1%. The target relative density (D_r) of
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13 96 the triaxial specimens was 90% since earth structures, such as road embankments and river levees, are well
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16 97 compacted. **Figure 3** shows the water retention curve for the silica sand with a D_r of 90%. A water retention test
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20 98 was performed using the negative water column technique [38]. The suction level of less than 10 kPa gives the
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23 99 residual degree of saturation of about 20%, which means that the water retention capability of silica sand is rather
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26 100 low.

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29 101 An acrylic pedestal was equipped with a water-saturated ceramic disc. The air entry value (AEV) of the ceramic
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32 102 disc is 0.5 bar (50 kPa) which is much higher than the water retention capability of silica sand. Air-dried sand was
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35 103 poured directly from a certain height into the water that was filled in the mold with a diameter of 35 mm beforehand.
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39 104 Then, the sand was densely packed by tapping the mold. Once the packed sand had reached a height of 70 mm, a
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42 105 top cap was placed on the specimen and negative pressure of 10 kPa was applied inside the specimen through the
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45 106 cap to stabilize the specimen by a difference in pressure to remove the mold. A confining pressure of 10 kPa was
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48 107 applied in the triaxial cell and then the negative pressure of 10 kPa, applied inside the specimen through the cap,
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51 108 was removed. After that, the confining pressure was increased to 50 kPa. The top of the specimen was released
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54 109 through a water tank to atmospheric pressure, i.e., a pore air pressure of 0 kPa, with almost 100% humidity, as
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57 110 shown in **Fig. 4a**. As shown in this figure, suction was applied to the specimen by the water head difference
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1 111 between the top of the specimen and the water table in the burette connected to the bottom of the specimen in order
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4 112 to desaturate the specimen.
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10 114 **2.1.2 Well-graded sand**

13 115 The water retention capability of silica sand is quite low, as shown in **Fig. 3**; and hence, it is difficult to measure
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16 116 the variation in suction during triaxial compression. In order to investigate the contributions of suction to the
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20 117 macroscopic responses, triaxial compression tests using well-graded sand with a higher water retention capability
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23 118 were performed. The well-graded sand used in the present study is Yodogawa sand which is typical on-site soil
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26 119 and has been used to improve river levees in Japan. A SEM image of the Yodogawa sand is shown in **Fig. 1b**.

29 120 **Figure 2** shows a grain size distribution curve reflecting 73% sand, 15% silt and 12% clay (e.g., [25]). The physical
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32 121 properties of Yodogawa sand include a soil density of 2.66 g/cm³, an optimum water content of 13.7% and a
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36 122 maximum dry density of 1.86 g/cm³. The water retention curve for Yodogawa sand with a degree of compaction
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39 123 (D_c) of 90% is shown in **Fig. 3**. The water retention test was performed using both the negative water column
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42 124 technique with 5 kPa of suction and the pressure plate method [32] with over 5 kPa of suction. It is clearly seen
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45 125 that the water retention capability of Yodogawa sand is much higher than that of silica sand.
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48 126 Yodogawa sand was mixed with a water content less than the optimum water content of 13.0%. The mixed sand
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51 127 was statically compacted in a mold, 35 mm in diameter and 70 mm in height. The target D_c was 90% which is the
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55 128 standard for earth structures determined by the Japan Institute of Construction Engineering (2002). The static
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58 129 compaction was done in five equal layers, with each layer having a thickness of 14 mm. The interfaces between
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1 130 the different layers were carefully scarified.

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4 131 Each compacted specimen was placed on an acrylic pedestal which was equipped with a water-saturated ceramic
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7 132 disc with an air entry value (AEV) of 50 kPa, and a polyflon filter was set on top of the specimen to measure and
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10 133 control the pore air pressure and the pore water pressure independently. The specimen was set in the triaxial cell
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13 134 and a cell pressure of 20 kPa was firstly applied by air pressure. The cell pressure was increased up to 60 kPa,
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16 135 while the pore air pressure inside the specimen was increased up to 40 kPa under an undrained water condition.
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19 136 Before the triaxial compression tests, the drained water and consolidation processes were performed. Suction was
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22 137 applied by the axis translation technique (for details see, e.g., [38]); namely, the pore water pressure, lower than
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25 138 the pore air pressure of 40 kPa, was increased up to a certain value corresponding to a target level of suction. For
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28 139 example, a suction of 30 kPa is given by 40 kPa in pore air pressure and 10 kPa in pore water pressure. The cell
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31 140 pressure was increased up to 90 kPa, resulting in a confining pressure of 50 kPa. The process to drain the water
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34 141 was begun by opening the valve for the water path connected to the bottom of the specimen. The open marks of
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39 142 **Fig. 3** are the water retention states of the tested specimens. These plots are located outside of the area bounded
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42 143 by the drying and wetting curves because the wetting curve is not really a main wetting curve, but one of scanning
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45 144 curves.

51 146 **2.2 X-ray micro tomography**

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54 147 The X-ray micro tomography facility used in the present study is KYOTO-GEO μ XCT (TOSCANER-
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57 32250 μ hdk, TOSHIBA IT and Control Systems Corporation) [24], as shown in **Fig. 4b**. The cone-beam type of
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1 149 scanning collects the attenuation characteristic of X-rays for several horizontal cross sections in the vertical
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4 150 direction at the same time, and the three-dimensional CT volume is obtained by the image reconstruction. The X-
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7 151 ray tomography with two kinds of spatial resolution was performed at axial strain levels of 0%, 2%, 4%, 6%, 9%,
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10 152 12%, 15%, 18% and 21% for the silica sand, as shown in **Fig. 5**. One of them is global tomography, used to observe
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13 153 the entire specimen over a larger scan region. The other is local tomography, used to observe the local region of
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16 154 interest focusing on the shear band with higher spatial resolution, enough to identify individual soil particles and
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19 155 water retention states. Once the sand fails due to the shear band, the macroscopic responses of the sand seem to be
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23 156 predominant in the microstructural changes inside the shear band. Hence, the local region, including the shear
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26 157 band, was observed. The scan position of the local tomography, as shown in **Fig. 5**, is moved by triaxial
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29 158 compression. Thus, the scan position was adjusted at each axial strain, by which nearly the same region of interest
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32 159 was tracked until the end of the triaxial compression test. Voxel size of the global tomography was $77.7^2 \times 85.0 \mu\text{m}^3$,
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35 160 while that of the local tomography was $12.3^2 \times 14.0 \mu\text{m}^3$. Axial loading was suspended during both kinds of X-ray
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39 161 tomography for around two hours, after which it was resumed at the same strain rate. For Yodogawa sand, only
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42 162 the global tomography was performed at axial strains of 0%, 2%, 4%, 6%, 9%, 12%, 16%, 20%, 24% and 27% .
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48 164 **2.3 Test methods and conditions**

51 165 Triaxial compression tests were conducted for both materials under a low confining pressure of 50 kPa to
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54 166 investigate the strength property of partially saturated sand which exhibits dilation during shearing, assuming, for
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58 167 example, the near surface of embankments. The strain rate was 0.10%/min.
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1 168 In the present study, the triaxial compression tests on the partially saturated sand were conducted under two
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4 169 kinds of conditions. One is drained conditions for both air and water, in which the pore air pressure and the pore
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7 170 water pressure are constant during shearing; namely, the suction is constant. The other is a drained condition for
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10 171 air and an undrained condition for water, in which the pore water pressure is varied while maintaining the constant
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13 172 pore air pressure during shearing; namely, the suction is varied. The former condition assumes that partially
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16 173 saturated soil deforms without any excess pore water pressure due to slow shearing or the high permeability of the
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20 174 soil, while the latter condition assumes that partially saturated soil deforms with excess pore water pressure due to
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23 175 relatively fast shearing or low permeability, respectively. It should be noted that the room temperature was kept at
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26 176 20°C during the experiment, through which condensation and evaporation could mostly be prevented.

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29 177 The specimen conditions for silica sand are listed in **Table 1**, and those for Yodogawa sand are listed in **Table**
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32 178 **2**. The case names are CS for the triaxial compression test under drained conditions for both air and water (constant
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35 179 suction test), and CW for the triaxial compression test under a drained condition for air and an undrained condition
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39 180 for water (constant water content test), respectively. Four triaxial compression tests of Yodogawa sand were
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42 181 conducted under two different levels of initial suction, namely, 10 kPa and 30 kPa, respectively. Test case CS-s10
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45 182 represents the constant suction test with an initial suction of 10 kPa, whereas case CW-s30 represents the constant
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48 183 water content test with an initial suction of 30 kPa.

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52 185 **3 Image analysis methodologies**

53 54 55 186 **3.1 Trinarization**

1 187 One voxel mostly includes a gray value corresponding to one phase, while another voxel shares more than two
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4 188 phases. In this case, the gray value of the voxel is determined using the average density of their phases, which is
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7 189 called the partial volume effect. For partially saturated sand, there are three types of voxels due to the partial
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10 190 volume effect: the voxel sharing the soil particle phase and the pore water phase, the voxel sharing the pore water
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13 191 phase and the pore air phase and the voxel sharing the soil particle phase and the pore air phase, respectively. The
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16 192 gray value for the voxel sharing the soil particle phase and the pore air phase is similar to that for the pore water
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19 193 phase and, as such, the voxel is often misidentified as the pore water phase. Thus, the partial volume effect for
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23 194 partially saturated sand often causes an overestimation of the pore water volume [18].
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26 195 For the segmentation of the soil particle phase, the pore water phase and the pore air phase, the CT images have
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29 196 been trinarized using the region growing method (e.g., [17]). The trinarization technique used in the present study
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32 197 takes into account the voxels due to the partial volume effect separately from those of the pure soil, water and air
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35 198 phases when determining the tolerance of the region growing. This technique reasonably identifies the voxels due
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39 199 to the partial volume effect as the pore air phase; and therefore, the overestimation of the pore water volume is
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42 200 reduced as much as possible. The region growing technique is applied to the soil particle phase and the pore air
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45 201 phase, after which the pore water phase is given as the remaining voxels. The detail algorithm and the validity of
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48 202 the trinarization technique are described in [24].
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52 204 **3.2 Morphology analysis**

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55 205 The scope of this analysis is to reveal the morphological changes of the pore air and the pore water in the
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1 206 partially saturated sand under triaxial compression. In the present study, the pore air phase and the pore water
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4 207 phase extracted from the trinarized volumes are divided into some assemblies with individual continuity, referred
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7 208 to as clusters, using the 3D image analysis software Avizo9.4.0 (FEI). **Figure 6** shows an example of the trinarized
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10 209 image and the labeled image obtained by this image analysis. There exist pore water voxels due to the partial
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13 210 volume effect as well as the absorbed water surrounding the soil particles in the trinarized volume. The pore water
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16 211 voxels are removed using erosion and dilation (e.g., [15]) with two voxels in this order. The separated pore water
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20 212 is then labeled by assigning a unique number to all adjacent voxels that constitute a cluster, as shown in **Fig. 6b**.
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23 213 The cluster volumes and the number of clusters are quantified. The detail algorithm is described in [24].
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26 214 Here, an index “continuity” is defined as a ratio of a maximum cluster volume to a total cluster volume and the
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29 215 value range is from 0% to 100%. Based on this index, water retention states were identified as the following three
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32 216 states: 1) the continuity is almost 100% and 0% for water and air, respectively, 2) the continuity is almost 100%
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35 217 for both water and air and 3) the continuity is almost 0% and 100% for water and air, respectively [24]. It was also
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39 218 confirmed that the three states corresponded to insular-air saturation, funicular saturation and pendular saturation
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42 219 regimes [1] in this order. Therefore, the “continuity” is an indicator to identify the water retention states depending
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45 220 on saturation degrees. It should be noted, however, that the continuity does not strictly describe whether the pore
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48 221 water is continuous in reality but does just describe morphology of pore water, because the absorbed water (thin
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51 222 water film) around particles, via which the pore water is continuous in soils, is ignored. The voxels related to the
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55 223 absorbed water were removed during the erosion-dilation image processing.
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3.3 Curvature analysis

The air-water interface in partially saturated soil exhibits meniscus shape. The relationship between suction and the principal curvature of the air-water interface for ideal spheres is described by the following equation (e.g., [28]):

$$u_a - u_w = \frac{T}{r_1} - \frac{T}{r_2} \quad (1)$$

where u_a is the pore air pressure, u_w is the pore water pressure, $u_a - u_w$ is the suction, r_1 and r_2 are the radii of the principal curvature of the air-water interface and T is the surface tension of water which is 0.074 N/m at 20°C. Suction is positive for partially saturated soil; hence, r_1 is always smaller than r_2 . r_1 and r_2 are presented as follows [6]:

$$r_1 = R \left(\frac{1}{\cos \theta} - 1 \right) \quad r_2 = R \tan \theta - r_1$$
$$0^\circ \leq \theta \leq 53^\circ \quad r_1 < r_2, u_a - u_w > 0 \quad (2)$$
$$53^\circ \leq \theta \leq 90^\circ \quad r_1 \geq r_2, u_a - u_w \leq 0$$

where R is the radius of an ideal sphere, and θ is the filling angle of the pore water between ideal spheres. It is probable that particle shapes affect air-water interface curvatures, and hence the equation (1) may not be accurate for non-spherical particle. In the present study, since the particle shape of the silica sand, whose aspect ratio is roughly 1.5 from Fig. 1a, is relatively spherical among general soils, equation (1) is here assumed to be applicable.

Figure 7 shows the relationship between suction and principal curvature radii r_1 and r_2 based on equations (1) and (2) when R is 230 μm , which is almost equivalent to half of the D_{50} of each silica sand particle. This figure indicates that r_1 is two orders of magnitude smaller than r_2 with suction of 2.0 kPa which works in the triaxial

1 240 specimens of the silica sand. Additionally, the suction estimated by equation (1) monotonically increases as r_1
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4 241 decreases. This result confirms that the variation in suction can be sufficiently discussed by measuring r_1 . Hence,
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7 242 in the present study, principal curvature $1/r_1$ is calculated by an image analysis to evaluate suction.
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10 243 In the present study, the maximum value of principal curvature is calculated as the curvature of the air-water
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13 244 interface, $1/r_1$, using the 3D image processing software Avizo9.4.0 (FEI). Generally, the rate of change of the
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16 245 normal direction in given tangential directions is measured by the second fundamental form (e.g., [27]), i.e.,
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19 246 curvature tensor, which can be expressed by a symmetric two-by-two matrix. The eigenvalues of the matrix are
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23 247 the principal curvatures. In the image processing, curved surfaces are subdivided into triangular meshes, and then
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26 248 the second fundamental form on each triangular mesh is approximated to the computed normal deviations of the
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29 249 triangle mesh using a least-squares method. As a result, the maximum eigenvalue, namely, the maximum principal
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32 250 curvature, is calculated. Details of the algorithm, implemented in this software, are given in [33].
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36 251 The specific procedure of the curvature analysis is as follows. Firstly, erosion and dilation with two voxels are
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39 252 applied for the binary image of the pore water in order to remove the voxels for the pore water phase with respect
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42 253 to the partial volume effect and adsorbed water. The processed pore water phase is replaced with the original one,
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45 254 and then the air-water interface is extracted as shown in **Fig. 8a**. The interface is composed of numerous triangular
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48 255 meshes, as shown in **Fig. 8b**, whose size is almost homogeneous. Assuming that the target triangle for the
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51 256 calculation is the number zero, the triangles sharing the sides of the target triangle are assumed to be number one.
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54 257 Similarly, numbers are assigned to the triangles up to twenty, as shown in **Fig. 8b**. The triangles adjacent to the
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57 258 target triangle are considered as being in the neighborhood. The principal curvature for the target triangle is then
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1 259 approximated with the least squares method to minimize the errors (deviation) in the principal curvatures for all
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4 260 the considered triangles. This indicates that the principal curvature at a certain triangular mesh on the interface is
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7 261 approximated to be equivalent to those of the adjacent triangles. Subsequently, the median of the principal
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10 262 curvatures for all triangles is determined as the curvature of the air-water interface $1/r_1$ because a histogram of
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13 263 the curvatures obtained by the curvature analysis follows the log-normal distribution.

16 264 The validation of the curvature analysis was investigated. Firstly, the influence of the number of neighboring
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20 265 target triangles (see **Fig. 8b**) on the accuracy of the calculation of the principal curvature was investigated using a
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23 266 realistic virtual sphere with a diameter of 41 voxels whose voxel size is 12 μm on one side. It was found that the
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26 267 number of neighboring target triangles of twenty provided a more similar principal curvature to the theoretical one.
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29 268 Therefore, a limit value of twenty is used in the present study. Secondly, the validity of the curvature analysis for
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32 269 evaluating the principal curvature of the material was investigated. The curvature analysis was applied to realistic
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36 270 virtual spheres with five diameters, namely, 17 voxels, 19 voxels, 21 voxels, 31 voxels and 41 voxels, whose voxel
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39 271 size is 12 μm on one side. **Figure 9** shows a comparison of the calculated and the theoretical solutions, which
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42 272 indicates that both values have a linear relationship with the correlation coefficient of 0.971. Therefore, the
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45 273 curvature analysis supposed in the present study is an appropriate method for discussing the relative levels of the
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48 274 principal curvature of the air-water interface, i.e., suction.

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55 276 **4 Results**

58 277 **4.1 Macroscopic responses**

4.1.1 Deviator stress-axial strain relationships

Figure 10 shows the deviator stress-axial strain relationships and the volumetric strain-axial strain relationships.

The volume of the specimen at each axial strain was measured by counting the number of voxels corresponding to the specimens out of the global tomography volumes (see [16]). Several vertical short lines in the stress-strain relationships are due to the stress relaxation during scanning. The relaxation occurs because levels of the axial strain are kept constant during X-ray tomography. The previous study [16] confirmed that the stress relaxation did not have significant effect on the overall response of sand specimens under triaxial compression. It is found that the stress-strain curves for all cases exhibit strain hardening followed by strain softening, accompanied by volume expansion with an increasing axial strain.

As displayed in **Fig. 10a**, cases CS1 and CW1 show almost the same peak deviator stress at an axial strain of 6%. The stress-strain curves and the changes in volumetric strain until the residual stress are also similar to each other. The significant difference between the two cases is that, for case CW1, the deviator stress at an axial strain of 27%, hereinafter referred to as residual stress, is lower than that for case CS1. Similarly, almost the same peak deviator stress is obtained at an axial strain of 9% between cases CS-s10 and CW-s10, whereas the residual stress for case CW-s10 is lower than that for case CS-s10. These results confirm that the stress-strain curves for cases CS1 and CW1 are qualitatively identical to those for cases CS-s10 and CW-s10, respectively. As displayed in **Fig. 10b**, the difference in deviator stress is more significant; namely, the peak deviator stress and the residual stress for cases CW2 and CW-s30 are lower than those for cases CS2 and CS-s30, respectively. These results confirm that the stress-strain curves for cases CS2 and CW2 are qualitatively identical to those for cases CS-s30 and CW-

1 297 s30, respectively.
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4 298 **Figure 11** shows the amount of drained water during loading for cases CS-s10 and CS-s30. It can be seen that
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7 299 the water is drained with increasing axial strain, indicating that the water retention capability of the specimens
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10 300 becomes lower due to dilation.
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13 301 14 15 16 302 **4.1.2 Suction for Yodogawa sand** 17 18

19 303 **Figure 12** shows the suction-axial strain relationship, for which the suction is given by the difference between
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23 304 the pore air pressure of 40 kPa and the pore water pressure measured in the tests. The constant levels of suction of
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26 305 10 kPa and 30 kPa are maintained during the triaxial compression tests for cases CS-s10 and CS-s30, respectively.
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29 306 On the other hand, the suction decreases at the early stage of shearing for cases CW-s10 and CW-s30. For case
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32 307 CW-s10, the suction recovers to almost the initial level of suction of 10 kPa, after which it tends to slightly decrease
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35 308 as shearing progresses. In contrast, the suction for case CW-s30 greatly decreases with an increasing axial strain,
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39 309 although a slight increase in suction is observed at an axial strain of 3%. The same trend, namely, decreasing
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42 310 suction under an undrained condition for water during shearing, was observed in previous studies (e.g., [35]).
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45 311 46 47 48 312 **4.1.3 Macroscopic deformation observed by global tomography** 49 50

51 313 **Figure 13** shows representative vertical slices of the CT images obtained by global tomography and distributions
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54 314 of shear strain calculated by digital image correlation, namely, DIC (e.g., [14, 17, 34]). The displacement fields
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58 315 were obtained by DIC with a zero-normalized cross-correlation coefficient between the reference and deformed
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1 316 CT images. Shear strain levels are provided by the B matrix for the eight-node isoparametric finite elements and
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4 317 the second invariant of incremental deviatoric strain tensor is defined as the incremental shear strain (see [17] for
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7 318 details). In the CT images, the materials with high density, such as metal, are described by the whiter color, while
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10 319 those with low density, such as air and the dilated regions, are described by the black color. It is clearly seen from
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13 320 the vertical slices at an axial strain of 0% for cases CS1, CS2, CW1 and CW2 that the upper part of the specimens
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16 321 is blacker than the bottom part. This indicates that the distribution of the degrees of saturation in the specimens is
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19 322 heterogeneous at the initial condition. This is because a higher level of suction is imposed on the upper part of the
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22 323 specimens by the negative water column technique; thus, a greater amount of water is drained. The lower density
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25 324 region increases at the middle of the specimens with an increasing axial strain due to dilation. As shown by the
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28 325 vertical slices beyond an axial strain of 9%, the lower density regions develop locally in the upper right direction.
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31 326 This region apparently corresponds to where the larger shear strains locally occur, indicating that the lower density
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34 327 region developing locally in the specimen is a shear band. Clearer shear bands can be observed in the specimens
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37 328 for cases CS1 and CS2 than for cases CW1 and CW2 probably because a greater degree of volume change is
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40 329 allowed during shearing in the drained conditions for both air and water than in the drained condition only for air.
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45 330 It is found from Fig. 13 that, for the Yodogawa sand, the middle part of the specimens dilates significantly and
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48 331 then the blacker portions increase as shearing progresses. The shear bands are invisible from the CT images for
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51 332 cases CS-s10 and CS-s30; however, the DIC analysis provides the clear shear strain localization in the upper right
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54 333 direction. In contrast, for cases CW-s10 and CW-s30, no clearer shear bands are observed than those for cases CS-
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57 334 s10 and CS-s30. It is found from the above results that the failure mode of the specimens for the silica sand is
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1 335 qualitatively identical to that for the Yodogawa sand during shearing.

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7 337 4.2 Microscopic characteristics

10 338 4.2.1 Local void ratio and degree of saturation in trinarized volumes

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13 339 **Figure 14** shows the vertical slices of the local tomography images and their trinarized images at axial strain
14 340 levels of 0%, 9% and 21% for cases CS1, CS2, CW1 and CW2. The whitest portions in the CT images indicate

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20 341 inherent metal inclusions in the silica sand. The metal inclusions were regarded as the soil particle phase in the

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23 342 trinarized images. In the CT images, the lower density region in the upper right direction can be observed beyond

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26 343 an axial strain of 9%; namely, this region is spatially consistent with the shear band observed by global tomography.

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29 344 The trinarized slices demonstrate that the amount of pore water becomes smaller in the shear band than in the other
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32 345 portions, which suggests that the reduction in the water retention capability is mainly attributed to the development
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36 346 of dilative shear bands.

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39 347 The void ratios and the degrees of saturation in the trinarized volumes obtained by local tomography at each

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42 348 axial strain are calculated by counting the number of voxels for the soil particle phase, the pore water phase and

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45 349 the pore air phase. **Figure 15** shows the calculated local void ratio-axial strain relationship. The local void ratios

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48 350 at an axial strain of 0% for cases CS1 and CW1 are smaller than those for cases CS2 and CW2, since the specimens

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51 351 for cases CS1 and CW1 are denser. The local void ratios for all cases tend to be increased up to an axial strain of

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54 352 12%. The changes in the local void ratio become small beyond an axial strain of 15%, indicating that the sand is

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58 353 getting closer to the critical state. Reductions in the void ratio are observed at some of axial strain levels, although

1 354 the shearing progresses. This is probably attributed to a limitation of the local CT volume, which contains both

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4 355 dilated loose portions inside the shear bands and less-dilated or contracted dense portions outside the shear bands.

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7 356 Depending on the ratio of dense and loose portions, the calculated void ratio of the local volume possibly decreases.

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10 357 Nevertheless, the trend of increasing void ratio and approaching the critical state can be clearly seen in **Fig. 15**.

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13 358 **Figure 16** shows the calculated local degrees of saturation-axial strain relationship. At an axial strain of 0%, the

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16 359 local degrees of saturation for each case are different from those of the entire specimens listed in **Table 1**. One of

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19 360 the possible reasons is that the trinarized volumes contain a larger or smaller amount of pore water than the other

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22 361 portions of the specimens, namely, the heterogeneity of the degree of saturation. Another possible reason is that

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25 362 the voxels due to the partial volume effect are misidentified as the pore water phase, although the overestimation

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28 363 of the degree of saturation due to this phenomenon is reduced as much as possible. It can be seen from **Fig. 16** that

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31 364 the degrees of saturation apparently decrease with an increasing axial strain for each case.

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35 365 It is found from **Figs. 15** and **16** that, under the same confining pressure, the local void ratios at the critical state

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38 366 are almost identical between the drained condition for water and the undrained condition for water, and that the

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41 367 local degrees of saturation are also almost identical between the two conditions when the initial degrees of

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44 368 saturation are similar.

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51 370 **4.2.2 Water retention states in trinarized volumes evaluated by morphology analysis**

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54 371 Variations in the water retention states in the trinarized volumes under triaxial compression were investigated

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57 372 based on the continuity for pore water and pore air. **Figure 17** shows the variations in continuity for the pore water

1 373 and the pore air. The continuity for the pore air is kept at almost 100% for all cases. The continuity for the pore
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4 374 water decreases with an increasing axial strain for cases CS1 and CW1, whereas it remains at almost 0% for cases
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7 375 CS2 and CW2. **Figures 18a to 18d** show the labeled images of the pore water cluster at axial strain levels of 0%,
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10 376 4% and 15% for cases CS1, CS2, CW1 and CW2, respectively. It is clearly seen that, for cases CS1 and CW1 with
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13 377 high initial degrees of saturation, one continuous pore water cluster with large volume transits into discontinuous
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16 378 pore water clusters with increasing axial strain. For cases CS2 and CW2 with low initial degrees of saturation, the
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20 379 pore water clusters remain discontinuous state during loading. **Figure 18e** demonstrates examples of the
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23 380 discontinuous pore water clusters existing as liquid bridges between soil particles.

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26 381 **Figures 19a to 19d** show the cumulative volume of pore water clusters-cluster volume relationships for cases
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29 382 CS1, CS2, CW1 and CW2, respectively. It can be seen that one cluster with volume larger than 1.7×10^6 voxels
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32 383 (3.6 mm^3), showing 100% cumulative volume, exists at several levels of axial strain for cases CS1 and CW1. In
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35 384 the present study, the cluster with volume larger than 1.7×10^6 voxels is regarded as the continuous pore water
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38 385 cluster (a labeled image at an axial strain of 0% in, e.g., **Fig. 18a**), while the clusters with volume smaller than
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42 386 1.7×10^6 voxels are regarded as the discontinuous pore water clusters (a labeled image at an axial strain of 0% in
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45 387 e.g., **Fig. 18b**). In the present study, it is assumed that the discontinuous pore water clusters are identified as liquid
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48 388 bridges. Although some of the discontinuous pore water clusters may comprise a few liquid bridges even if the
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51 389 volume of such clusters is smaller than 1.7×10^6 voxels, the clusters play a role of bonding between particles and
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54 390 are obviously different from the continuous pore water which exists as a much larger volume cluster.

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58 391 **Figure 20** illustrates the number density of the discontinuous pore water clusters-axial strain relationship. The
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1 392 number density is calculated by dividing the number of the discontinuous pore water clusters by the total volume
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4 393 of the discontinuous pore water clusters. In other words, the number density means the number of liquid bridges
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7 394 per unit pore water volume. It can be seen that the number densities for cases CS1 and CW1 tend to decrease with
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10 395 increasing axial strain up to 12%, and then turn to increase slightly beyond an axial strain of 15%. The increase in
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13 396 the number density indicates that the continuous pore water transits into the discontinuous pore water with small
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16 397 volume due to dilation, as shown in **Figs. 18a** and **18c**. On the other hand, the number density for case CW2 is
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19 398 smaller and tends to decrease more greatly than that for case CS2 with increasing axial strain. At the initial state,
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23 399 there is a relatively big difference in the level of the number density between cases CS2 and CW2, even though
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26 400 the global void ratio and the global degree of saturation are similar between both the specimens. In the case of the
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29 401 lower degree of saturation, the water retention state and the volume distribution of pore water will vary by just a
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32 402 little difference in the local degree of saturation. It is therefore quite difficult to investigate a local part of each
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35 403 specimen where the water retention state and the volume distribution of pore water are the same using local
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39 404 tomography.

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42 405 It is found from **Figs. 10** and **20** that, except for the result at an axial strain of 4%, the number density for case
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45 406 CW1 is almost the same as that for case CS1 until an axial strain of 6% where the peak deviator stresses for the
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48 407 two cases are similar to each other. The smaller number density at the axial strain of 4% for case CW1 possibly
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51 408 results from a variation in the local void ratio. As seen in **Fig. 15**, the void ratio for case CW1 greatly increases
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54 409 from axial strains of 2% to 4%. In this case, grain contacts where liquid bridges exist may be greatly decreased,
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58 410 which reduces the number of liquid bridges. Investigating variation in grain contacts would be important to clarify

1 411 this point. On the other hand, the smaller number density is observed in case CW1 than in case CS1 at axial strains
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4 412 of 18% and 21%, corresponding to the lower residual stress in case CW2 than in CS2. It is also found that the
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7 413 number density for case CW2 is smaller than that in case CS2 at both an early stage of shearing and the residual
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10 414 state, corresponding to the lower peak deviator stress and the lower residual stress in case CW2 than in case CS2.

16 416 4.2.3 Variation in principal curvature of air-water interface

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20 417 **Figure 21** shows the principal curvature of the air-water interface-axial strain relationship in the trinarized
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23 418 volumes. The principal curvatures of the air-water interface at an axial strain of 0% for cases CS2, CW2, CW1
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26 419 and CS1 are higher in this order; they correspond to the magnitude correlation of the initial levels of suction for
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29 420 each case shown in **Table 1**. The principal curvature of the air-water interface for case CS1 is almost constant up
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32 421 to an axial strain of 21%, and comparable to the test condition under which the macroscopic suction remains
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35 422 constant during shearing. The principal curvature of the air-water interface for case CS2 tends to be also similar at
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39 423 each axial strain. On the other hand, the principal curvatures of the air-water interface for cases CW1 and CW2
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42 424 tend to gradually decrease with increasing axial strain and finally become lower than those for cases CS1 and CS2
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45 425 at an axial strain of 21%, respectively. These results confirm that the principal curvature of the air-water interface
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48 426 under an undrained condition for water during shearing decreases as the shearing progresses.

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51 427 The variations in suction and principal curvature of the air-water interface are compared to each other. As shown
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55 428 in **Figs. 12** and **21**, the constant levels of suction for cases CS-s10 and CS-s30 almost correspond to the constant
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58 429 principal curvatures of the air-water interface for cases CS1 and CS2, respectively. In addition, the decrease in the

1 430 levels of suction for cases CW-s10 and CW-s30 corresponds to the decrease in the principal curvatures of the air-
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4 431 water interface for cases CW1 and CW2, respectively. Therefore, the decrease in suction under the undrained
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7 432 condition for water during shearing can be interpreted as the decrease in the principal curvature of the air-water
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10 433 interface.

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16 435 **5 Discussions**

20 436 **5.1 Effects of drained conditions for water on microscopic characteristics of pore water**

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23 437 Increasing the void ratio due to dilation causes a decrease in the degree of saturation as the shearing progresses,
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26 438 and then the water retention state in partially saturated sand with higher initial degrees of saturation tends to transit
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29 439 from a continuous pore water cluster to discontinuous pore water clusters (see **Figs. 18a** and **18c**). These trends
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32 440 are independent of the drained conditions for water during shearing. In water retention test [24], a continuous pore
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36 441 water cluster with large volume transited into discontinuous pore water clusters with smaller volume during the
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39 442 drying process. Decreasing degree of saturation, either by dilation or by drying causes the transition of the pore
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42 443 water morphology from the continuous state to the discontinuous state. On the other hand, under the undrained
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45 444 condition for water during shearing, the number of liquid bridges decreases more greatly than under the drained
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48 445 condition for water during shearing and there exists the smaller number of liquid bridges (see **Fig. 20**). Additionally,
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51 446 the principal curvature of the air-water interface under the undrained condition for water during shearing tends to
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54 447 decrease, whereas that under the drained condition for water during shearing remains almost constant.

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58 448 The cause of the decrease in the number of liquid bridges under the undrained condition for water during
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1 449 shearing is interpreted in terms of the variation in the principal curvature of the air-water interface. Let us consider
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4 450 the liquid bridges between the ideally spherical grain contacts and the radius of the principal curvature of the air-
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7 451 water interface, r_1 , presented by equation (2). As shown in equation (2), the pore water can exist as liquid bridges
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10 452 and suction works with θ smaller than 53° . Assuming that R is $228 \mu\text{m}$ for the silica sand grains and the range in
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13 453 the principal curvatures of the air-water interface in the present study is from $0.009 \mu\text{m}^{-1}$ to $0.014 \mu\text{m}^{-1}$ (see **Fig.**
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16 454 **21**), the range in θ is between 40° and 48° ; and thus, liquid bridges can exist between the grains. The principal
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20 455 curvature of the air-water interface decreases in the undrained condition for water during shearing, which
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23 456 corresponds to the increase in r_1 in equation (2). Considering that R is constant, the increase in r_1 causes the
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26 457 increase in θ ; namely, along with the increase in θ , up to larger than 53° , comes the tendency for a decrease in
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29 458 liquid bridges at the grain contacts during shearing. Therefore, the decrease in liquid bridges under triaxial
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32 459 compression with the undrained condition for water is attributed to the decrease in the principal curvature of the
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35 460 air-water interface as well as the decrease in grain contacts. In contrast, the principal curvature of the air-water
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39 461 interface in the drained condition for water during shearing remains constant; namely, θ remains constant. This
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42 462 indicates that the decrease in liquid bridges is not attributed to the varying principal curvatures of the air-water
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45 463 interface, but mainly to the decrease in grain contacts under the drained condition for water during shearing. Thus,
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48 464 there is a larger decrease in the number of liquid bridges in the undrained condition for water.
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54 466 **5.2 Link of microscopic characteristics to macroscopic responses**

58 467 Firstly, the link of microscopic characteristics to the peak deviator stress is discussed. **Figure 10b** indicates that

1 468 the partially saturated silica sand with low initial degrees of saturation displays higher peak deviator stress under
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4 469 the CS test compared to that under the CW test. For both cases, a lot of liquid bridges (the discontinuous pore
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7 470 water with the continuity of almost 0% as shown in Fig. 17) initially exist as indicated by Figs. 18b, 18d and 18e.
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10 471 The grains are strongly bonded together due to suction working as inter-particle force through the liquid bridges.
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13 472 Therefore, the shear stiffness and shear strength of the partially saturated soil are enhanced. Even if the levels of
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16 473 initial suction are almost the same between the CS test and the CW test of the silica sand, the number of liquid
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20 474 bridges tends to decrease more greatly induced by the decrease in the principal curvature of the air-water interface
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23 475 under the CW test, resulting in the lower peak deviator stress than under the CS test. This idea is supported by the
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26 476 results obtained in the CW test with high level of initial suction for Yodogawa sand; the level of suction (the
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29 477 principal curvature of the air-water interface) is significantly decreased at an early stage of loading, resulting in
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32 478 the lower peak deviator stress than under the CS test.

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35 479 On the other hand, Fig. 10a indicates that the levels of peak deviator stress for the silica sand with high initial
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39 480 degrees of saturation under the CS test and the CW test are identical to each other; this trend is consistent with the
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42 481 levels of peak deviator stress for the Yodogawa sand with low initial suction. A continuous pore water exists in the
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45 482 partially saturated silica sand as indicated by labeled images at an axial strain of 0% in Figs. 18a and 18c. Assuming
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48 483 this water retention state corresponds to the “bulk water”, explained by, e.g., [21] and [41] (the pore water with
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51 484 the continuity larger than 30%, as shown in Fig. 17), the suction contributes to the isotropic effective stress. In the
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54 485 continuous water retention state, the effect of suction through liquid bridges on the bonding between grains is less
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58 486 significant than in the discontinuous water retention state. In cases CS-s10 and CW-s10 for the Yodogawa sand,
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1 487 the pore water is more continuous than that in the cases with high level of suction, and there is little difference in
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4 488 the measured suction between the two cases at the peak deviator stress with an axial strain of about 6% as shown
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7 489 in **Fig. 12**. Thus, the contribution of suction, mainly due to bulk water, to the shear stiffness and the shear strength,
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10 490 is similar between the CS test and the CW test, leading to the identical levels of peak deviator stress between them.
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13 491 Secondly, the link of microscopic characteristics to the residual stress is discussed. As indicated in **Fig. 10a**, the
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16 492 residual stress for the silica sand and the Yodogawa sand under the CW test is lower than that under the CS test.
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19 493 Even though the pore water is initially continuous, it becomes discontinuous as shearing progresses; namely, the
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23 494 proportion of liquid bridges increases. The tendency for a decrease in the number of liquid bridges during the post-
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26 495 peak softening process can be observed for both the CS test and the CW test. However, under the CW test the
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29 496 principal curvature of the air-water interface, i.e., the suction, decreases and becomes lower than that under the CS
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32 497 test during the post-peak softening process even if the levels of initial suction are the same. Additionally, the
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35 498 number of liquid bridges decreases by a greater degree than in the CS test. As a result, the contribution of the
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39 499 suction through liquid bridges to the shear stiffness and the shear strength of the partially saturated soil becomes
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42 500 lower under the CW test than under the CS test during the post-peak softening process. This provides the difference
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45 501 in residual stress between CW and CS tests. In residual stress state of Yodogawa sand, the deviator stress and the
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48 502 measured suction of the CW test are smaller than those of the CS test (see **Figs. 10b and 12**), and the difference in
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51 503 deviator stress between the two tests is clearer than that of the silica sand. The same trend of the Yodogawa sand
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54 504 as the silica sand confirms that the deviator stress at the critical state is smaller in the CW test than in the CS test.
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1 506 **6 Concluding remarks**

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4 507 Triaxial compression tests under drained conditions for both air and water (constant suction test, i.e., CS test)
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7 508 and those under a drained condition for air and an undrained condition for water (constant water content test, i.e.,
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10 509 CW test) were conducted on dense specimens with partially-water saturated conditions. Triaxial tests were carried
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13 510 out under low confining pressure to investigate soil behaviors associated with dilative shear bands. Silica sand was
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16 511 used as the poorly-graded sand in order to clarify the microscopic characteristics of pore water using X-ray micro
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19 512 tomography and image analysis techniques. Yodogawa sand, i.e., typical on-site soil, was used as the well-graded
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23 513 sand in order to measure the suction variation and to confirm the findings obtained using the silica sand.
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26 514 Trinarization quantified changes in the void ratio and degree of saturation. Morphology analysis classified the
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29 515 continuous water cluster, i.e., bulk water, and the discontinuous water clusters, i.e., liquid bridges. Curvature
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32 516 analysis of air-water interface provided the suction changes. Using these image analysis results, the link of the
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35 517 microscopic characteristics to the macroscopic responses of the partially saturated silica sand was revealed.
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39 518 Although the differences of the macroscopic stress-strain responses and the microscopic analyzed quantities such
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42 519 as the curvature between CS test and CW test seems small for the silica sand, these findings were successfully
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45 520 supported by the tests using the Yodogawa sand, in which the tendencies are the same as those of the silica sand
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48 521 and the differences of stress-strain behaviors and measured suction were sufficiently clear.
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51 522 Under the same confining pressure, the void ratios at the critical state were almost identical between the CS test
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54 523 and the CW test. The degrees of saturation were also almost identical between the two tests when the initial degrees
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57 524 of saturation were similar. However, the morphology of the pore water and curvatures of the air-water interface
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1 525 were clearly different depending on drained conditions during shearing and initial degrees of saturation. The
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4 526 relationships between the microscopic characteristics and the macroscopic responses of the partially saturated soil
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7 527 under different initial water retention states clarified in this study are summarized in **Fig. 22**.

10 528 In cases of high initial degrees of saturation (**Fig. 22a**), pore water exists as bulk water rather than liquid bridges
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13 529 and the suction mainly contributes to isotropic effective stress. Although suction slightly decreases in CW test,
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16 530 there are little difference in suction levels between CS test and CW test due to the low suction level. Hence, the
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20 531 effect of suction on shear stiffness and shear strength are similar between the two tests, resulting in almost the
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23 532 same peak deviator stress. As shearing progresses, the water retention state becomes discontinuous with dilation,
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26 533 and suction contributes to shear stiffness and shear strength through liquid bridges. The number of liquid bridges
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29 534 increases in the middle of the experiment, and then turns to decrease at the residual stress state. The curvature of
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33 535 the air-water interface, i.e., suction, gradually decreases in CW test, leading to a greater decrease in the number of
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36 536 liquid bridges than in CS test. Finally, the residual stress of CW test becomes lower than that of CS test.

39 537 The water retention state in cases of low initial degrees of saturation (**Fig. 22b**) is more discontinuous than
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42 538 that in the high saturation cases, and thus the suction contributes to shear stiffness and shear strength through liquid
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45 539 bridges rather than through bulk water. Since the suction level is higher than that of the high saturation cases,
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48 540 suction decreases more greatly in CW test, resulting in lower peak deviator stress. The number of liquid bridges
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51 541 decreases with dilation as shearing progresses. Similar to the high saturation cases, the decrease in the curvature
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55 542 of the air-water interface induces a greater decrease in the number of liquid bridges in CW test, which results in
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58 543 lower deviator stress in CW test than that in CS test at the residual stress state.

1 544 As future work, it will be important to clarify the cause of the decrease in the principal curvature of air-water
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4 545 interface under the CW test. It will also be important to investigate the microscopic characteristics of partially
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7 546 saturated contractive sand and their link to the macroscopic responses.
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8 **Tables**

9
 10 652 **Table 1** Specimen conditions for silica sand

Case	CS1	CS2	CW1	CW2
D_r (%)	94.97	90.41	95.35	91.69
Void ratio	0.710	0.725	0.709	0.722
Initial S_r (%)	63.60	47.60	60.52	46.20
Suction (kPa)	1.66	1.70	1.67	1.69

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 19 653
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 21 654 **Table 2** Specimen conditions for Yodogawa sand

Case	CS-s10	CS-s30	CW-s10	CW-s30
D_c (%)	89.39	89.55	89.94	89.85
Void ratio	0.585	0.586	0.589	0.588
Initial S_r (%)	49.11	44.50	47.84	44.22
Suction (kPa)	10	30	10	30

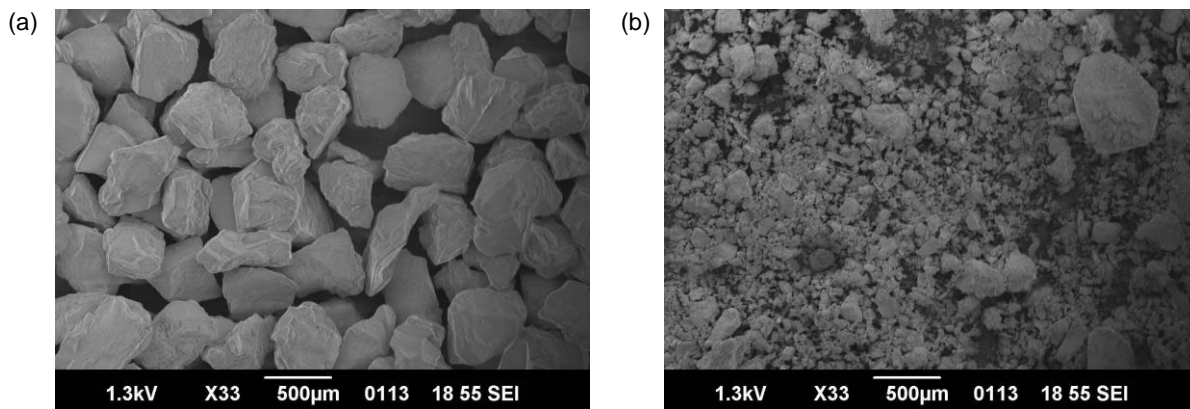


Fig. 1 SEM images: **a** silica sand (poorly-graded sand) and **b** Yodogawa sand (well-graded sand)

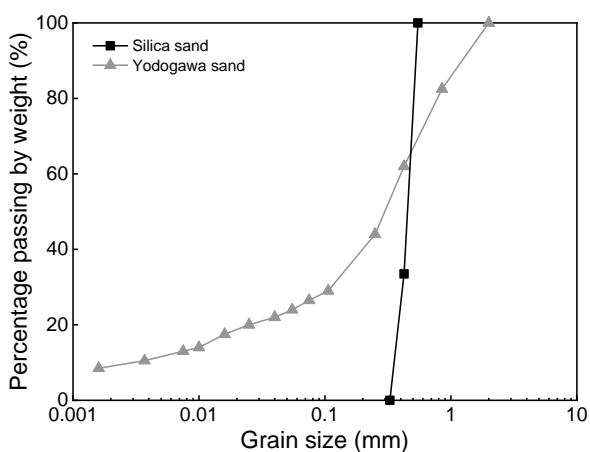


Fig. 2 Grain size distribution curves

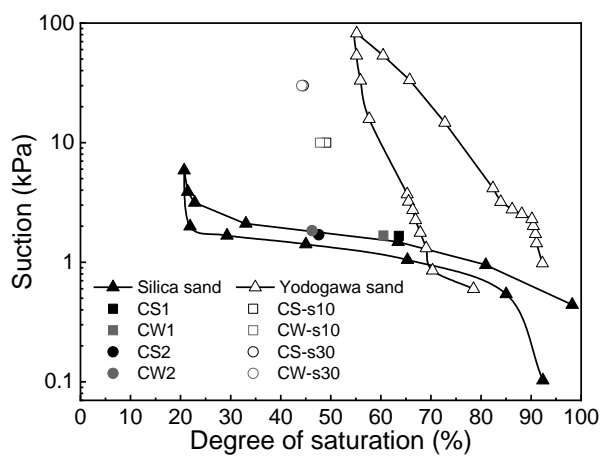


Fig. 3 Water retention curves

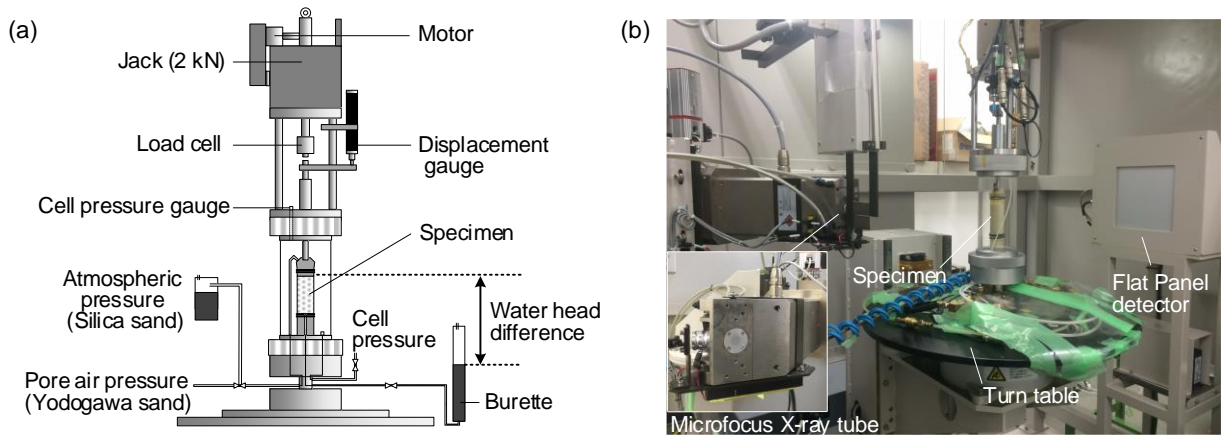


Fig. 4 **a** setup of triaxial compression test apparatus and **b** overview of X-ray micro tomography system

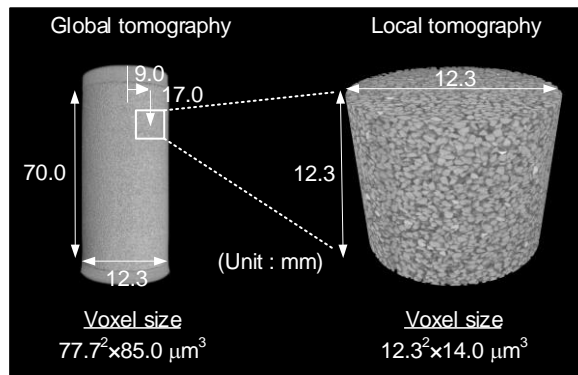


Fig. 5 Scan types with different spatial resolutions

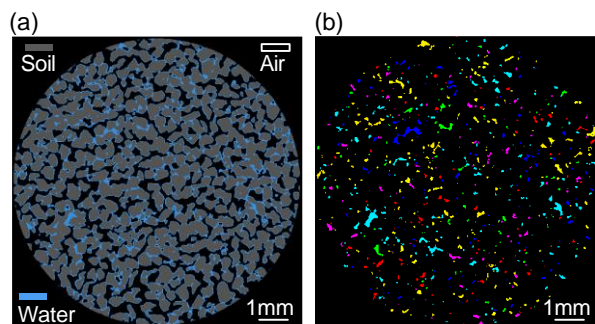


Fig. 6 Example of images: **a** trinarized image and **b** labeled image of pore water obtained by morphology analysis

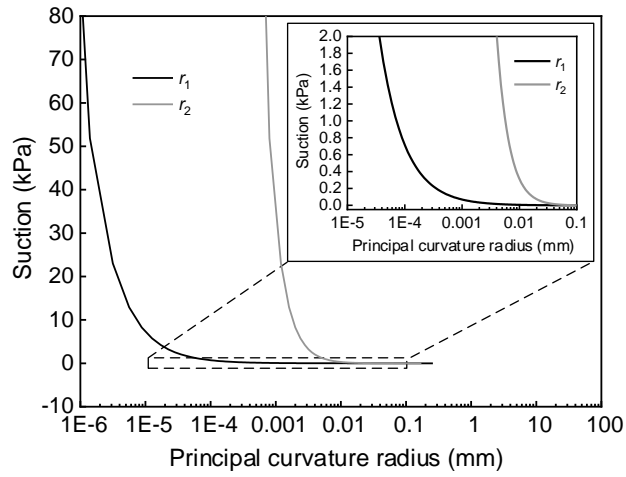


Fig. 7 Relationship between suction and principal curvature radii r_1 and r_2

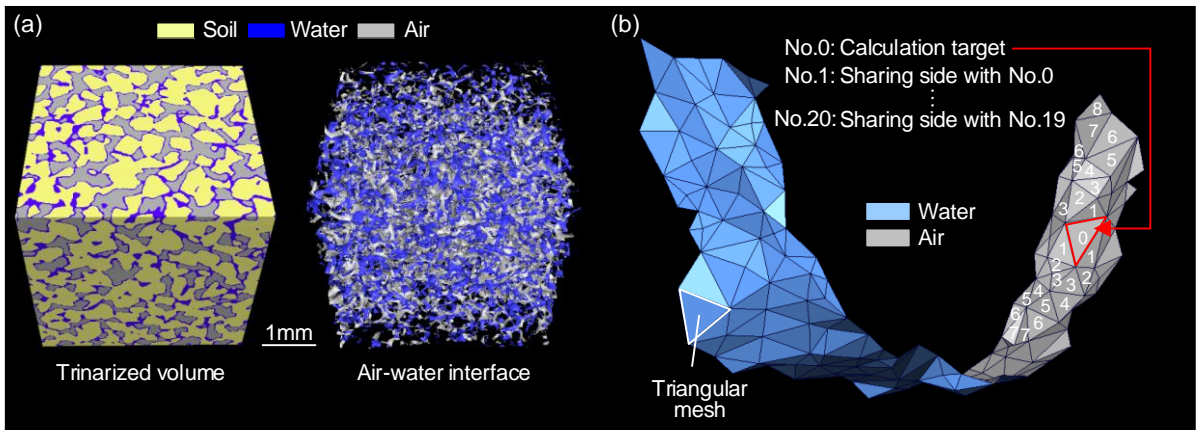


Fig. 8 **a** air-water interface extracted from trinarized volume and **b** triangular meshes and principle of considering neighborhood of calculation target

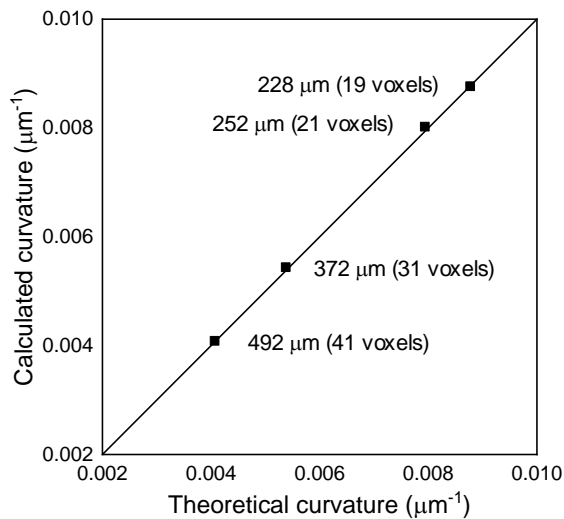


Fig. 9 Relationship between calculated curvature and theoretical curvature

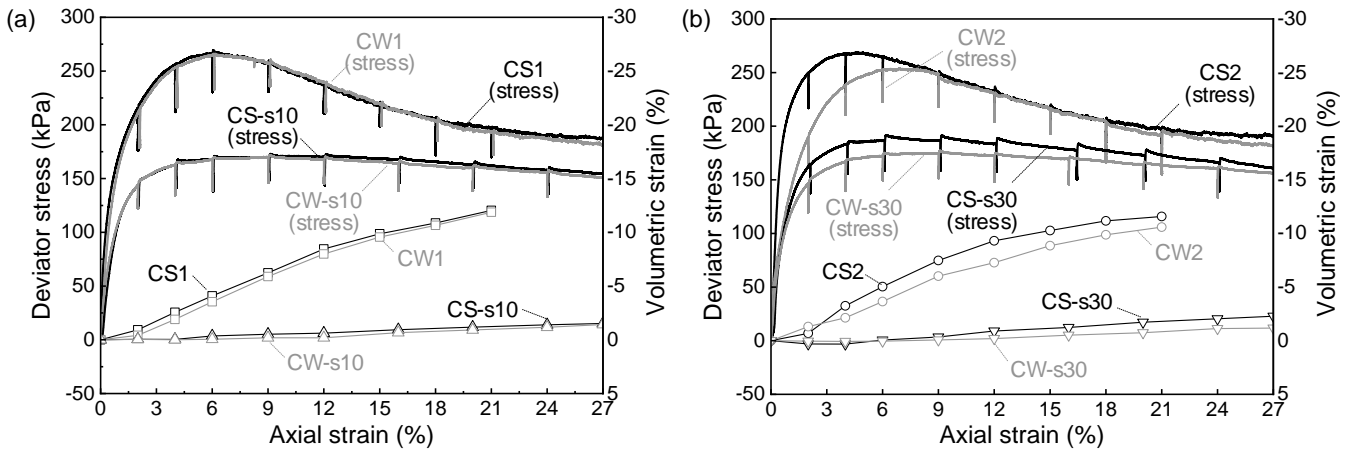


Fig. 10 Deviator stress-axial strain relationships and volumetric strain-axial strain relationships: **a** cases CS1, CW1, CS-s10 and CW-s10 and **b** cases CS2, CW2, CS-s30 and CW-s30

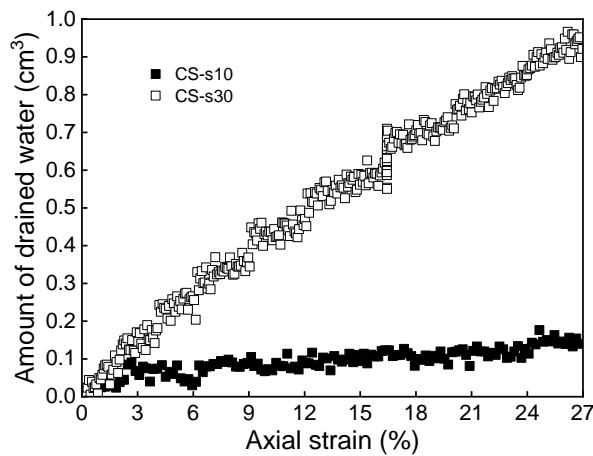


Fig. 11 Amount of drained water-axial strain relationship for cases CS-s10 and CS-s30

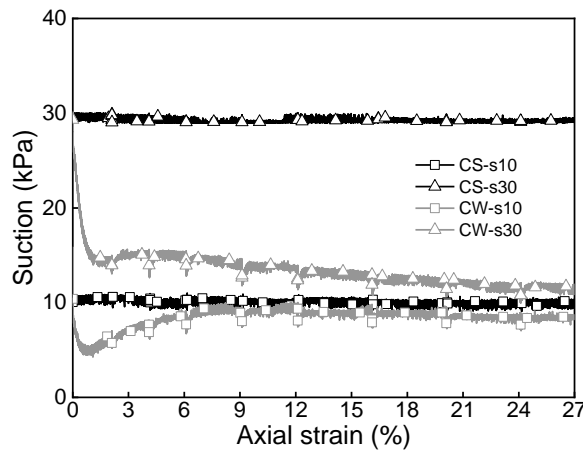


Fig. 12 Suction-axial strain relationship for Yodogawa sand

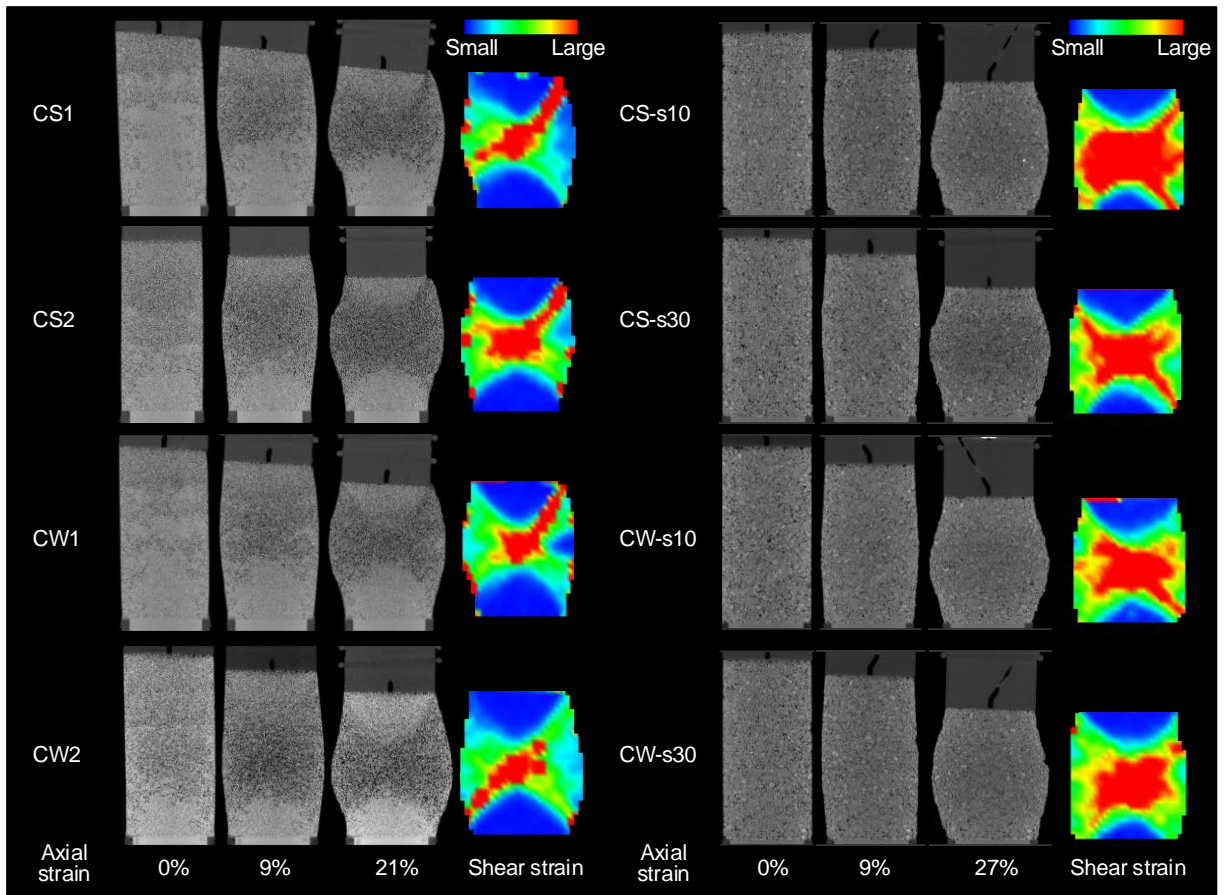


Fig. 13 Vertical slices of global tomography images and distributions of incremental shear strain

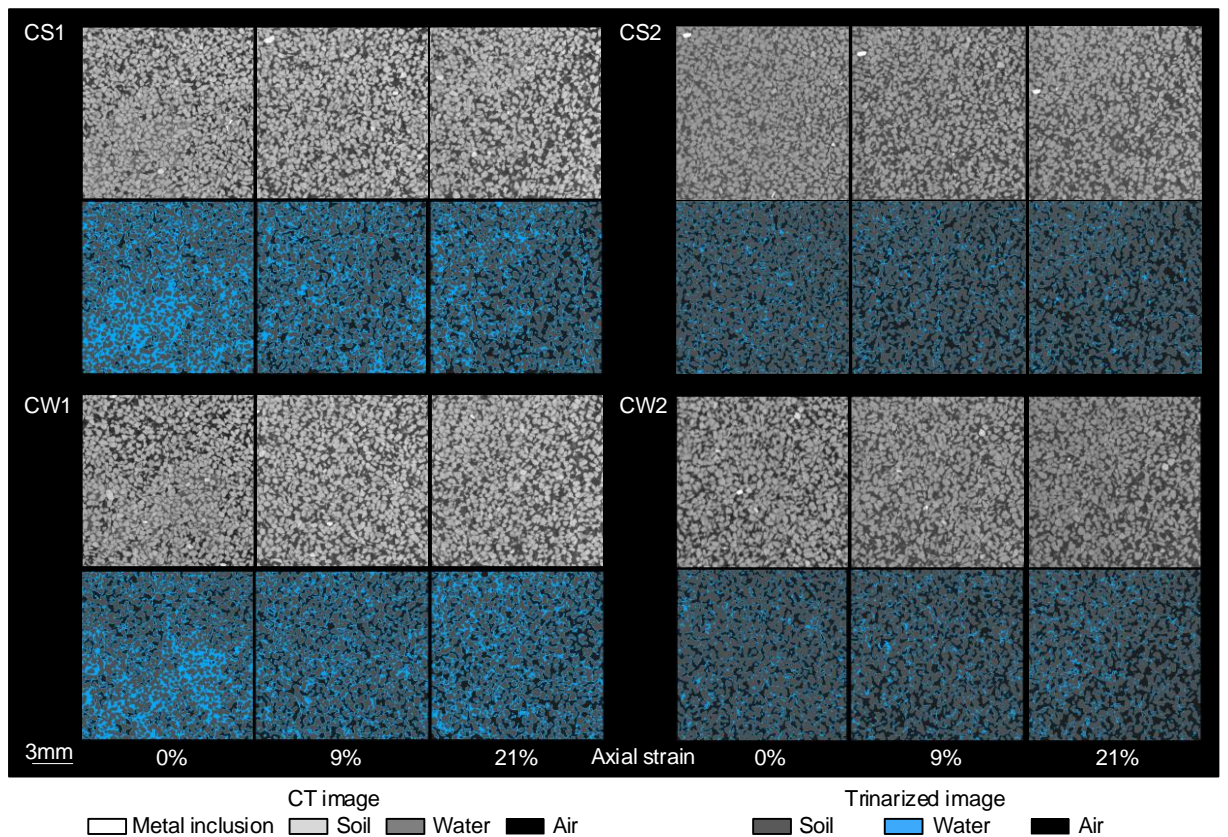


Fig. 14 Vertical slices of local tomography and trinarized images for silica sand

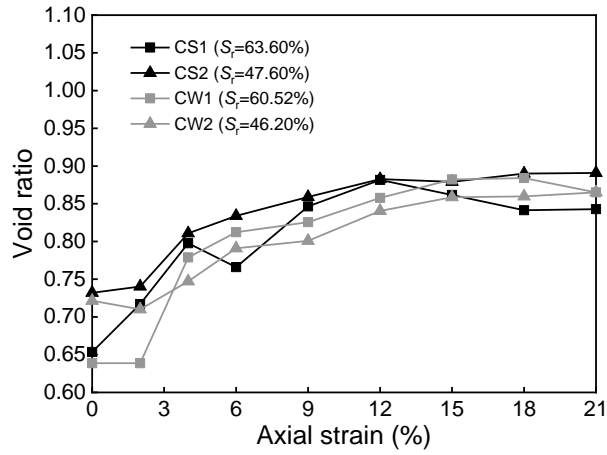


Fig. 15 Local void ratio-axial strain relationship for silica sand

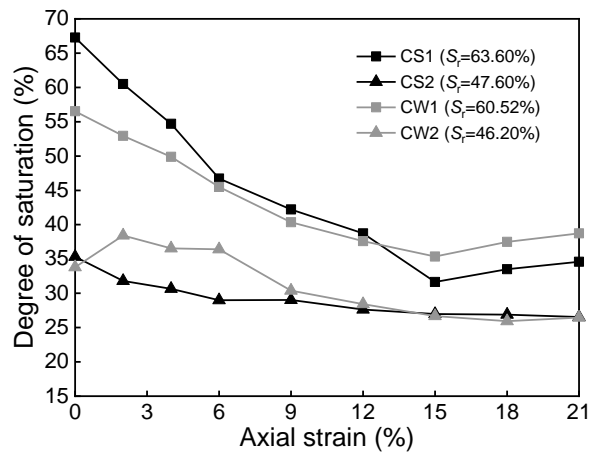


Fig. 16 Local degree of saturation-axial strain relationship for silica sand

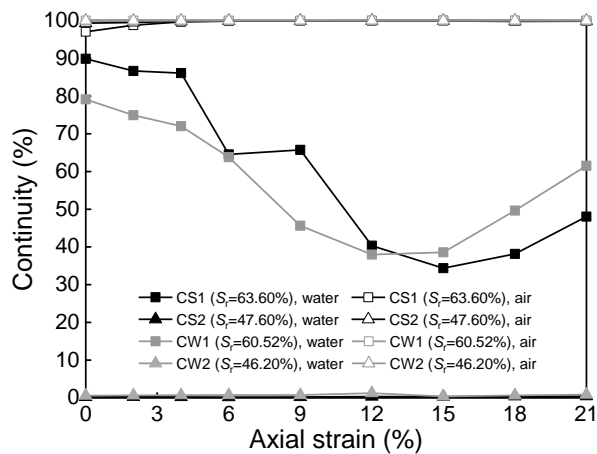


Fig. 17 Variations in continuity for pore water and pore air for silica sand

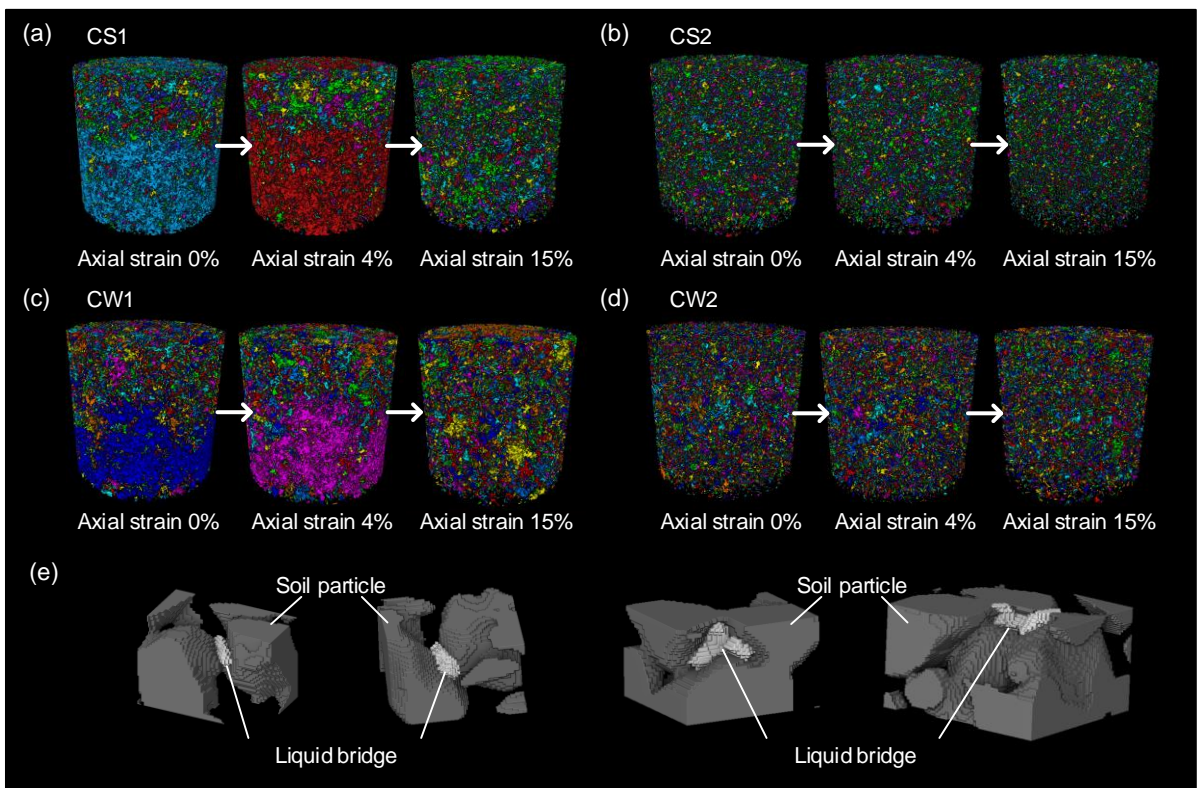


Fig. 18 Transitions of water retention states for silica sand: **a** case CS1, **b** case CS2, **c** case CW1, **d** case CW2 and **e** liquid bridges between soil particles

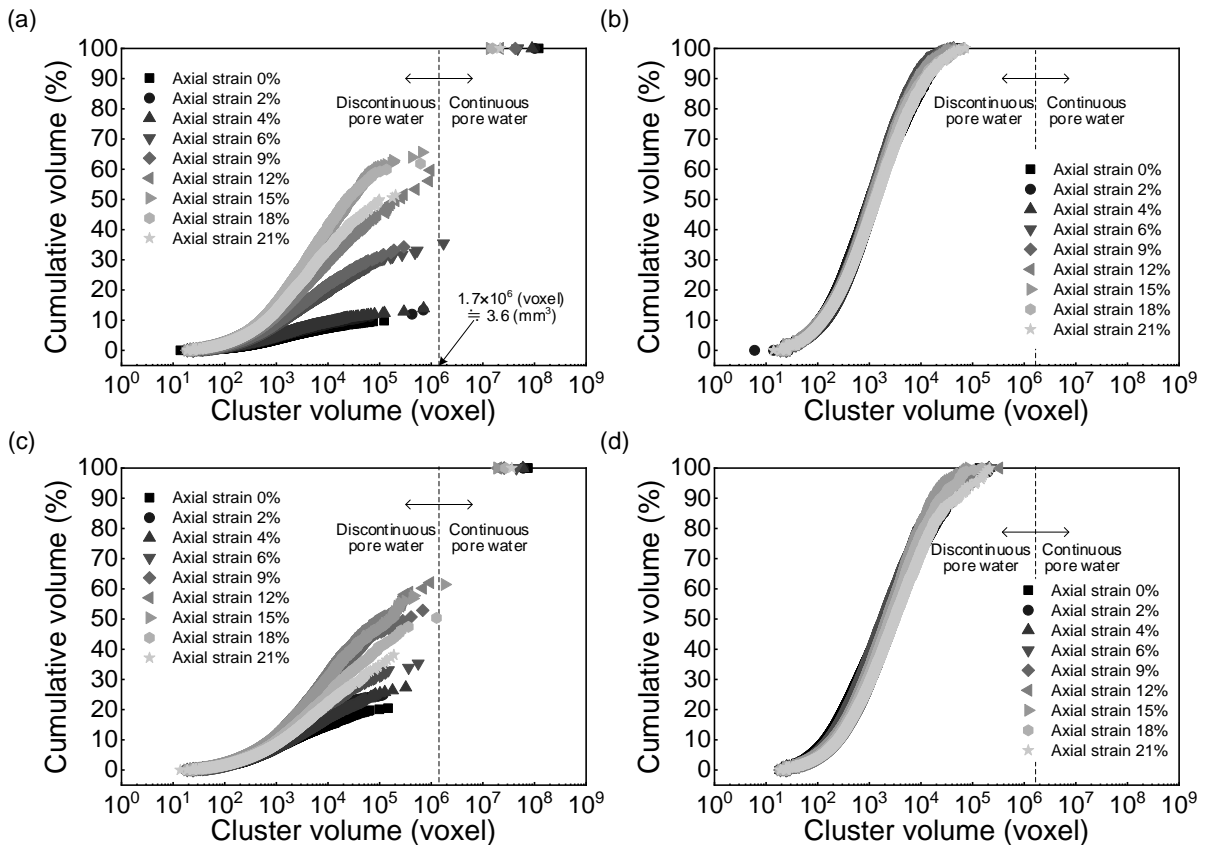


Fig. 19 Cumulative volume of pore water clusters-cluster volume relationships for silica sand: **a** case CS1, **b** case CS2 **c** case CW1 and **d** case CW2

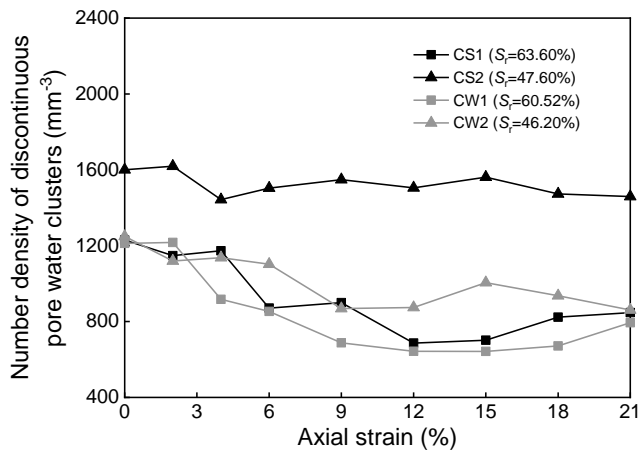


Fig. 20 Number density of discontinuous pore water clusters-axial strain relationship for silica sand

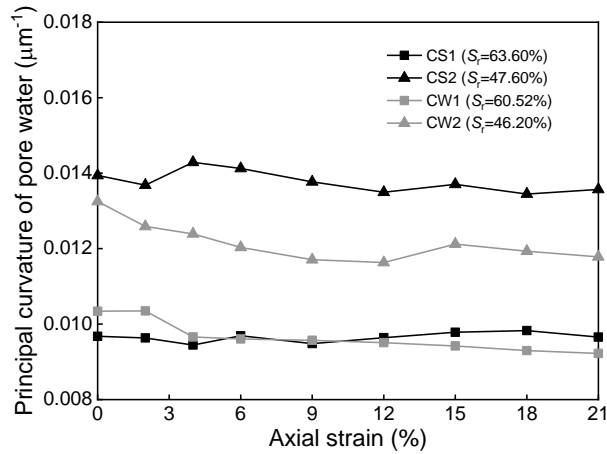


Fig. 21 Principal curvature of air-water interface-axial strain relationship for silica sand

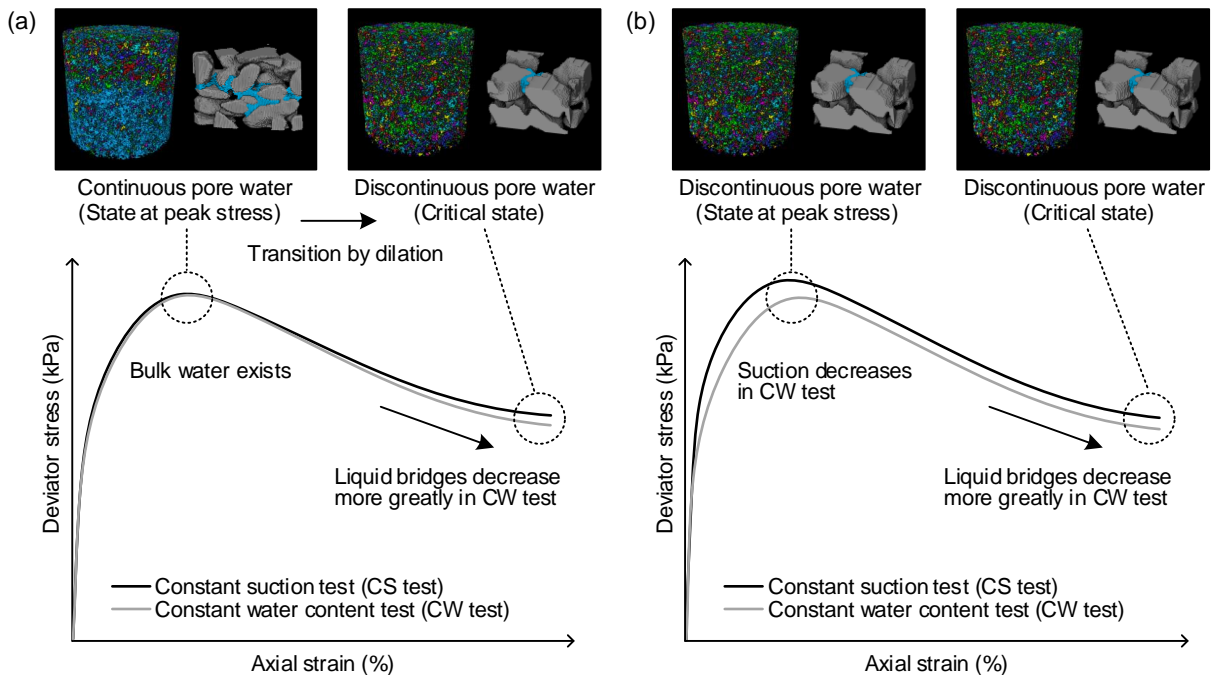


Fig. 22 Microscopic characteristics of partially saturated soil and their link to macroscopic responses: **a** under high initial degree of saturation and **b** under low initial degree of saturation