# Topologically Representation of Cantor Cube Model for Geometric Patterns

### Shousuke Ohmori

Department of Physics, Waseda University, 3-4-1 Okubo, Shinjuku-ku, Tokyo 169-8555, Japan

### Abstract

The characteristic geometric structures found in disordered matters are discussed from the viewpoint of general topology. Cantor cube which is a topological space consisting of the infinite product space of 0 and 1 provides specific decomposition spaces representing topologically geometric patterns of matters such as graphs, clusterized structures, dendrites.

### 1 Introduction: Cantor cube model

Geometrically patterns of matters in solid and liquid states have been hugely studied from the viewpoint of disordered physics[1, 2]. In particular, fascinating characteristic topology structures of matters (the topology structures of matters is abbreviated as TSM hereafter), such as the graphic structure of polymers[3], the clusterized structure of molecular liquids[4, 5], or the dendritic structure in solidifications[6], have been found. To characterize such TSM, several mathematical methods based on topological concept have been studied. For instance, persistent homology is the method available for classification of geometric structures of amorphous materials[7, 8]. Note that this persistent homological analysis is mathematically based on a technique of algebraic topology.

The mathematical approach by using general topology to TSM has been successfully studied[9, 10, 11, 12]. In these studies, TSM are investigated based on the mathematical formation of a set of equivalence classes for a specific topological space X, independently of the detailed properties of each matter. That is, geometric patterns for TSM are discussed by connecting them with the decomposition spaces of X (For the details of this approach, see [12]). In the present article, we focus on characterizations of TSM by means of the Cantor cube model; in this model the specific topological space X is taken as a Cantor cube  $(\{0, 1\}^{\Lambda}, \tau_0^{\Lambda})$  and geometric patterns are discussed through the decomposition spaces of  $(\{0, 1\}^{\Lambda}, \tau_0^{\Lambda})$ , where  $(\{0, 1\}^{\Lambda}, \tau_0^{\Lambda})$  is the  $\Lambda$ -product space of  $(\{0, 1\}, \tau_0)$  with an index set  $\Lambda$  of Card  $\Lambda \geq \aleph_0$  and  $\tau_0$  is a discrete topology for  $\{0, 1\}^{\Lambda}, \tau_0^{\Lambda}$ ) homeomorphic to a compact metric space.

It is very known that any compact metric space is represented homeomorphically as a quotient space of fibers [13] of 0-dim, perfect, compact Hausdorff-space[14]. Indeed, the representation of a compact metric space can be obtained systematically as follows. Let X denote a Cantor cube  $(\{0,1\}^{\Lambda}, \tau_0^{\Lambda})$  with Card  $\Lambda \geq \aleph_0$ , and let Y be a compact metric space. Then, there exists a closed cover  $\{Y_1, \ldots, Y_n\}$   $(n < \infty)$  of Y, each diameter of which is less than 1/2. To this cover there corresponds a partition  $\{X_1, \ldots, X_n\}$  of  $\{0, 1\}^{\Lambda}$  such that

$$\begin{cases} X_1 = \{0\}_{\lambda_1} \times \{0,1\}^{\Lambda-\{\lambda_1\}}, \\ X_i = \{1\}_{\lambda_1} \times \dots \times \{1\}_{\lambda_{i-1}} \times \{0\}_{\lambda_i} \times \{0,1\}^{\Lambda-\{\lambda_1,\dots,\lambda_i\}} \ (i=2,3,\dots,n-1), \\ X_n = \{1\}_{\lambda_1} \times \dots \times \{1\}_{\lambda_{n-2}} \times \{1\}_{\lambda_{n-1}} \times \{0,1\}^{\Lambda-\{\lambda_1,\dots,\lambda_{n-1}\}}, \end{cases}$$
(1)

where  $\lambda_i$  is arbitrarily element of  $\Lambda$ , (i = 1, ..., n - 1) and each  $\{k_1\}_{\lambda_1} \times \cdots \times \{k_i\}_{\lambda_i} \times \{0, 1\}^{\Lambda - \{\lambda_1, ..., \lambda_i\}} = \{x : \Lambda \to \{0, 1\}, x(\lambda_l) = k_l \in \{0, 1\}, l = 1, ..., i\}$  stands for a cone. Let  $s_1 : X \to \Im(Y) - \{\emptyset\}$  be a map defined by  $s_1(x) = Y_i$  if  $x \in X_i$  for each i, where  $\Im(Y)$  is the collection of closed sets of Y. Note that  $Y = \bigcup_{x \in X} s_1(x)$ . Since  $Y_i$  is a compact metric space,

for each *i*, we have a closed cover  $\{Y_{i_1}, \ldots, Y_{i_{n_i}}\}$  of  $Y_i$ , each diameter of  $Y_{i_j}$  being less than  $1/2^2$ . Also,  $X_i$  has a partition  $\{X_{i_1}, \ldots, X_{i_{n_i}}\}$  composed of cones such that

$$\begin{cases} X_{i_{1}} = \{1\}_{\lambda_{1}} \times \dots \times \{1\}_{\lambda_{i-1}} \times \{0\}_{\lambda_{i}} \times \{0\}_{\mu_{1}} \times \{0,1\}^{\Lambda - (\{\lambda_{1},\dots,\lambda_{i}\} \cup \{\mu_{1}\})}, \\ X_{i_{j}} = \{1\}_{\lambda_{1}} \times \dots \times \{1\}_{\lambda_{i-1}} \times \{0\}_{\lambda_{i}} \times \{1\}_{\mu_{1}} \times \dots \times \{1\}_{\mu_{j-1}} \times \{0\}_{\mu_{j}} \times \{0,1\}^{\Lambda - (\{\lambda_{1},\dots,\lambda_{i}\}) \cup (\{\mu_{1},\dots,\mu_{j}\})} \\ (j = 2,3,\dots,n_{i} - 1), \\ X_{i_{n_{i}}} = \{1\}_{\lambda_{1}} \times \dots \times \{1\}_{\lambda_{i-1}} \times \{0\}_{\lambda_{i}} \times \{1\}_{\mu_{1}} \times \dots \times \{1\}_{\mu_{n_{i}-2}} \times \{1\}_{\mu_{n_{i}-1}} \times \{0,1\}^{\Lambda - (\{\lambda_{1},\dots,\lambda_{i}\}) \cup (\{\mu_{1},\dots,\mu_{j}\})} \end{cases}$$

$$(2)$$

where  $\mu_i$  is arbitrarily element of  $\Lambda - (\{\lambda_1, \ldots, \lambda_i\})$ . Let  $s_2 : X \to \Im(Y) - \{\emptyset\}$  be defined by  $s_2(x) = Y_{i_j}$  for  $x \in X_{i_j}$ . Then,  $Y = \bigcup_{x \in X} s_2(x)$  and  $s_2(x) \subset s_1(x)$  for all x. Continuing the procedure we have a sequence of functions  $\{s_n\}$  such that for each x and for each n, (i)  $s_n$  is upper semi-continuous, (ii)  $s_{n+1}(x) \subset s_n(x)$ , (iii)  $Y = \bigcup_{x \in X} s_n(x)$ , and (iv) dia  $s_n(x) \to 0$  as  $n \to \infty$ , where dia stands for diameter of a set. Thus, we obtain a continuous map f from X onto  $Y, x \mapsto \bigcap_n g_n(x)$  and the decomposition space  $(\mathcal{D}_f, \tau(\mathcal{D}_f))$  of X relative to f homeomorphic to Y, where  $\mathcal{D}_f = \{f^{-1}(y); y \in Y\}$  and  $\tau(\mathcal{D}_f) = \{\mathcal{U} \subset \mathcal{D}_f; \bigcup \mathcal{U} \in \tau_0^\Lambda\}$  is a decomposition topology. Through the homeomorphism, each point y of Y can be associated with an unique point  $f^{-1}(y)$  of  $\mathcal{D}_f$ . For instance, the decomposition space representing [0, 1] is obtained practically as the following two cases; letting  $M \equiv \{l/2^n; n = 1, 2, \ldots$  and  $l = 1, \ldots, 2^n - 1\}$ , then (i) for  $y = \sum_{i=1}^{\infty} a_i/2^i \notin M$ 

$$f^{-1}(y) = \{a_1\}_{\lambda_1} \times \{a_2\}_{\lambda_2} \times \dots \times \{0, 1\}^{\Lambda - \{a_1, a_2, \dots\}},\tag{3}$$

and (ii) for  $y = l/2^n \in M$ 

$$f^{-1}(y) = \left[ \{a_1\}_{\lambda_1} \times \{a_2\}_{\lambda_2} \times \dots \times \{a_{n-1}\}_{\lambda_{n-1}} \times \{0\}_{\lambda_n} \times \{1\}_{\lambda_{n+1}} \times \{1\}_{\lambda_{n+2}} \times \dots \times \{0,1\}^{\Lambda - \{\lambda_1,\lambda_2,\dots\}} \right] \\ \cup \left[ \{a_1\}_{\lambda_1} \times \{a_2\}_{\lambda_2} \times \dots \times \{a_{n-1}\}_{\lambda_{n-1}} \times \{1\}_{\lambda_n} \times \{0\}_{\lambda_{n+1}} \times \{0\}_{\lambda_{n+2}} \times \dots \times \{0,1\}^{\Lambda - \{\lambda_1,\lambda_2,\dots\}} \right]$$
(4)

for some  $a_1, \ldots, a_{n-1}$ . Here,  $f^{-1}(0) = \{0\}_{\lambda_1} \times \{0\}_{\lambda_2} \times \cdots \times \{0, 1\}^{\Lambda - \{\lambda_1, \lambda_2, \cdots\}}$  and  $f^{-1}(1) = \{1\}_{\lambda_1} \times \{1\}_{\lambda_2} \times \cdots \times \{0, 1\}^{\Lambda - \{\lambda_1, \lambda_2, \cdots\}}$ . Note that the decomposition space constructed in the above process is not unique.

### 2 Topologically representation for TSM

Here, we focus on several decomposition spaces of the Cantor cube  $X = (\{0, 1\}^{\Lambda}, \tau_0^{\Lambda})$  which represent geometric models for TSM.

First let us consider two network patterns  $Y_1$  and  $Y_g$  shown in (a) and (b) of Fig. 1;  $Y_1$  is a figure composed of three nodes  $e_1, e_2, a$  and two bonds  $E_1$  and  $E_2$  connecting  $e_1$  with a and  $e_2$  with a, respectively.  $Y_g$  is a finite graph[15]. Since  $Y_1$  is regarded as an arc, the construction of the decomposition space stated in the previous section for Y = [0, 1] can be directly applied to  $Y_1$ . Indeed, letting h be a homeomorphism from  $Y_1$  onto [0, 1], each point x of  $Y_1$  can be represented as the point of a decomposition space  $\mathcal{D}_1$  of  $\{0, 1\}^{\Lambda}$  by the following two types; (i) if  $h(x) \notin M (\equiv \{l/2^n; n = 1, 2, ... \text{ and } l = 1, ..., 2^n - 1\})$ , then

$$x \doteq \{k_1\}_{\lambda_1} \times \{k_2\}_{\lambda_2} \times \dots \times \{0,1\}^{\Lambda - \{\lambda_1, \lambda_2, \dots\}},\tag{5}$$

where  $k_1, k_2, \ldots$  are points in  $\{0, 1\}$  satisfying  $h(x) = \sum_{i=1}^{\infty} k_i/2^i$ , and  $\doteq$  is the sign of identification of x with a corresponding point  $f^{-1}(x)$  of  $\mathcal{D}_1$ , and (ii) if  $h(x) \in M$ , then

$$x \doteq \left[ \{k_1\}_{\lambda_1} \times \{k_2\}_{\lambda_2} \times \dots \times \{k_m\}_{\lambda_m} \times \{0\}_{\lambda_{m+1}} \times \{1\}_{\lambda_{m+2}} \times \{1\}_{\lambda_{m+3}} \times \dots \times \{0,1\}^{\Lambda - \{\lambda_1,\lambda_2,\dots\}} \right] \\ \cup \left[ \{k_1\}_{\lambda_1} \times \{k_2\}_{\lambda_2} \times \dots \times \{k_m\}_{\lambda_m} \times \{1\}_{\lambda_{m+1}} \times \{0\}_{\lambda_{m+2}} \times \{0\}_{\lambda_{m+3}} \times \dots \times \{0,1\}^{\Lambda - \{\lambda_1,\lambda_2,\dots\}} \right]$$
(6)

77

for some m, where  $k_1, \ldots, k_m$  are points in  $\{0, 1\}$  giving  $h(x) \in M$ . If we introduce a sign  $S_x$  defined by

$$S_x \equiv \begin{cases} (5), & h(x) \notin M\\ (6), & h(x) \in M, \end{cases}$$
(7)

then

$$x \doteq S_x \tag{8}$$

for  $x \in Y_1$ . Note that assuming  $h(e_1) = 0$  and  $h(e_2) = 1$ , the end points  $e_1$  and  $e_2$  form

$$e_1 \doteq \{0\}_{\lambda_1} \times \{0\}_{\lambda_2} \times \dots \times \{0,1\}^{\Lambda - \{\lambda_1, \lambda_2, \dots\}}, \ e_2 \doteq \{1\}_{\lambda_1} \times \{1\}_{\lambda_2} \times \dots \times \{0,1\}^{\Lambda - \{\lambda_1, \lambda_2, \dots\}}.$$
 (9)

The relation (8) shows that the geometric feature of  $Y_1$  is completely characterized in the decomposition space  $\mathcal{D}_1$  of X. For a finite graph  $Y_g$ , we denote the arcs composing of  $Y_g$  by  $E_1, \ldots, E_r(r < \infty)$ . There exists a partition  $\{X^1, \ldots, X^r\}$  of X corresponding to these arcs where each  $X^i$  is defined as well as that in (1) with indexes  $\mu_1, \ldots, \mu_{r-1} \in \Lambda$ . It is confirmed that a decomposition space  $\mathcal{D}_g$  of X represents  $Y_g$ ; the representations for a node x with bonds  $E_{t_1}, \ldots, E_{t_g}$  and a point y in a bond  $E_i$  are obtained as

$$x \doteq \bigcup_{j=1}^{q} (X^{t_j} \cap S_x^{t_j}), \ y \doteq X^i \cap S_y^i, \tag{10}$$

respectively. As a practical example of materials with the graphic structure we can consider a tree of a dendritic crystal[16]. A tree is a graph that does not contains a space homeomorphic to a unit sphere shown in (c) of Fig. 1. In this case, the representation for a tree by a decomposition space  $\mathcal{D}_t$  is the same as the relation (10).



Figure 1: Schematic explanation of three types of geometric patterns with network configuration. (a) geometric model  $Y_1$ ; two nodes  $e_1$  and  $e_2$  are connected by edges  $E_1$  and  $E_2$  thorough a node a. (b) a finite graph  $Y_g$ . (c) a tree  $Y_t$ .

Next, we focus on a cluster pattern  $Y_c$  for which each cluster is a finite graph, shown in Fig. 2 (a). Then,  $Y_c$  may be defined to be a topological space  $(\bigoplus_{i=1}^s C_i, \bigoplus_{i=1}^s \tau_i)$  where  $(\bigoplus_{i=1}^s C_i, \bigoplus_{i=1}^s \tau_i)$  is a disjoint union of a collection of finite graphs  $\{(C_i, \tau_i), i = 1, \ldots, s\}$ . To disjoint clusters  $C_1, \ldots, C_s$ , there corresponds a partition  $\{J_1, \ldots, J_s\}$  of X using new elements  $\xi_1, \ldots, \xi_{s-1} \in \Lambda$  such that

$$\begin{cases} J_1 = \{0\}_{\xi_1} \times \{0,1\}^{\Lambda-\{\xi_1\}}, \\ J_j = \{1\}_{\xi_1} \times \cdots \times \{1\}_{\xi_{j-1}} \times \{0\}_{\xi_j} \times \{0,1\}^{\Lambda-\{\xi_1,\cdots,\xi_{j-1}\}} \ (j=2,\ldots,s-1), \\ J_s = \{1\}_{\xi_1} \times \cdots \times \{1\}_{\xi_{s-2}} \times \{1\}_{\xi_{s-1}} \times \{0,1\}^{\Lambda-\{\xi_1,\cdots,\xi_{s-1}\}}. \end{cases}$$
(11)

By applying the relation (10) to each finite graph  $C_i$  for corresponding cone  $J_i, i = 1, ..., s$ , the representation of whole space  $Y_c$  by a decomposition space  $\mathcal{D}_c$  of X can be obtained as follows; for  $x \in Y_c$  contained in a cluster  $C_{i_0}$ ,

$$x \doteq J_{i_0} \cap \begin{cases} \cup_{j=1}^q (X^{t_j} \cap S_x^{t_j}), \\ X^i \cap S_x^i. \end{cases}$$
(12)

This topologically representation by a decomposition space  $\mathcal{D}_c$  for a clusterized structure can be applied to the tiling issue in material science that a polycrystal can be filled with an arbitrary finite number of single crystals characterized by a specific geometric structure, i.e., dendritic, or self-similar structure (according to the mathematical setting and discussion for the issue based on general topology, see [11]). Figure 2 (b) shows the roughly sketch of situation for this issue. Here, we consider dendritic crystals as single crystals. Then, the situation can be identified with the clusterized geometrical pattern in which each cluster is dendritic. Actually, we regard each dendritic crystal composing of the polycrystal as a cluster and then the geometric structure of the polycrystal is described by a kind of clusterized structure. Based on the representation (12) of the clusterized structure, we can obtain the following decomposition space  $\mathcal{D}_c$  of X representing the geometric structure of the polycrystal:

$$\mathcal{D}_c = \bigcup_{i=1}^n \mathcal{D}_i,\tag{13}$$

where

$$\mathcal{D}_i = \left\{ y \doteq J_i \cap \bigcup_{j=1}^q (X^{t_j} \cap S_y^{t_j}); y \in Y_t^i \right\} \cup \left\{ y \doteq J_i \cap X^j \cap S_y^j; y \in Y_t^i \right\}.$$
(14)

Note that  $\mathcal{D}_i$ ,  $i = 1, \dots, i$  are mutually disjoint each other. (13) and (14) show the relationship between each single dendritic crystal  $\mathcal{D}_i$  and a whole polycrystal  $\mathcal{D}_c$  for the tiling issue. The representation of decomposition spaces for the clusterized structure stated in this section can be widely applicable to discuss geometric aggregation structures of matters such as noncrystalline and amorphous as well as this tiling issue for a polycrystal.



Figure 2: Geometric models of (a) a clusterized structure  $Y_c$  where the number of clusters s = 3, and (b) schematic explanations of a polycrystal Z filled with dendritic decomposition spaces  $\mathcal{D}_i$ .

Finally, we comment that in this short article a Cantor cube is introduced as a conceptional model to obtain several topologically representations of TSM, e.g., the graphic and clusterized structures. Indeed, each character of these geometric structures can be connected with decomposition spaces of a Cantor cube. Therefore, by analyzing a mathematical property of a Cantor cube model even more, new universal properties of the geometric structures of matters might be revealed.

# 3 Conclusion

The mathematical method to characterize geometric patterns for TSM universally based on a Cantor cube  $(\{0,1\}^{\Lambda}, \tau_0^{\Lambda})$  has been shown. Typical geometric patterns such as graphic and

clusterized structures are focused on and their representations by decomposition spaces of the Cantor cube are investigated. A practical form of a decomposition space of a polycrystal in the tiling issue that a polycrystal filled with an arbitrary finite number of dendritic crystals is shown by handling it as a special case of the decomposition representation for the clusterized structural geometric model. (More details are shown in [12].)

#### acknowledgment

The authors are grateful to Prof. Y. Yamazaki, Prof. T. Yamamoto, and Prof. Emeritus A. Kitada at Waseda university for useful suggestions and encouragements.

# References

- [1] J. M. Ziman, Models of Disorder (Cambridge University Press, 1979).
- [2] N.E. Cusack, The Physics of Structurally Disordered Matter: An Introduction (Univ. Sussex Press, 1987).
- [3] Y. Tezuka and H. Oike, J. Am. Chem. Soc. **123**, 11570 (2001).
- [4] Y. Katayama, T. Mizutani, W. Utsumi, O. Shimomura, M. Yamakata, and K. Funakoshi, Nature (London). 403, 170 (2000).
- [5] T. Morishita, Phys. Rev. Lett. 87, 105701 (2001).
- [6] W. Kurz and D. J. Fisher, *Fundamentals of Solidification* (Trans Tech Publication Ltd, Switzerland 1998).
- [7] A. Hirata, K. Matsue, and M.Chen, Structural Analysis of Metallic Glasses with Computational Homology (Springer, 2016).
- [8] Y. Hiraoka, T. Nakamura, A. Hirata, E. G. Escolar, K. Matsue, and Y. Nishiura, PNAS. 113, 7035 (2016).
- [9] A. Kitada and Y. Ogasawara, Chaos, Solitons & Fractals. 24, 785 (2005); Erratum Chaos, Solitons & Fractals. 25, 1273 (2005).
- [10] A. Kitada, Y. Ogasawara, and T. Yamamoto, Chaos, Solitons & Fractals. 34, 1732 (2007).
- [11] A. Kitada, S. Ohmori and T. Yamamoto, J. Phys. Soc. Jpn. 85, 045001 (2016).
- [12] S. Ohmori, Y. Yamazaki, T. Yamamoto, and A. Kitada, Phys. Scr. 94, 105213 (2019).
- [13] Each decomposition space handled in the present paper is represented as a quotient space of fibers of a continuous map.
- [14] For example, see P. Alexandorff, Math. Ann. 96, 555 (1927), F. Hausdorff, Mengenlehre (Berlin, 1927).
- [15] S. B. Nadler Jr., Continuum Theory (Marcel Dekker, 1992).
- [16] A dendrite is defined as a topological space which is a Peano continuum (namely, a connected, locally connected, compact, metric space) containing no simple closed curve. Mathematically, a tree is a dendrite but it is not necessarily that a dendrite implies a tree. However, the dendritic structure composed of infinitely many branches may not be observed in diffusion process such as solidification in practice. Therefore, we consider a dendrite as a tree here. According to the detailed discussions for the dendritic structures in solidification, see [6].