Irreducible modules for W_3 -algebra

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1 Introduction

Recently a W_3 -algebra of central charge 6/5 was studied as an orbifold of an order three automorphism of the lattice vertex operator algebra $V_{\sqrt{2}A_2}$ by Dong, Lam, Tanabe, Yokoyama and the author [5]. The paper consists of four parts: (1) realization of a W_3 -algebra \mathcal{W} of central charge 6/5, (2) classification of irreducible \mathcal{W} -modules, (3) C_2 -cofiniteness and rationality of \mathcal{W} , and (4) determination of the character of each irreducible \mathcal{W} -module. In this note we will give an outline of the second part of the paper, namely, classification of irreducible \mathcal{W} -modules. Details can be found in [5]. Basic references to \mathcal{W}_3 -algebras are [1] and [7].

2 The W_3 -algebra W

For definitions of the materials discussed here we refer to [2, 8]. We also use certain properties of the vertex operator algebra $V_{\sqrt{2}A_2}$ (cf. [12]).

Let α_1, α_2 be the simple roots of type A_2 and set $\alpha_0 = -(\alpha_1 + \alpha_2)$. We denote the inner product by $\langle \cdot, \cdot \rangle$. Then $\langle \alpha_i, \alpha_i \rangle = 2$ and $\langle \alpha_i, \alpha_j \rangle = -1$ if $i \neq j$. Set $\beta_i = \sqrt{2}\alpha_i$ and let $L = \mathbb{Z}\beta_1 + \mathbb{Z}\beta_2 = \sqrt{2}A_2$ be the lattice spanned by β_1 and β_2 . We follow Sections 2 and 3 of [2] with $L = \sqrt{2}A_2$, p = 3, and q = 6. In our case $\langle \alpha, \beta \rangle \in 2\mathbb{Z}$ for all $\alpha, \beta \in L$, so that the alternating \mathbb{Z} -bilinear map $c_0 : L \times L \to \mathbb{Z}/6\mathbb{Z}$ defined by [2, (2.9)] is trivial. Thus the central extension

$$1 \longrightarrow \langle \kappa_6 \rangle \longrightarrow \widehat{L} \stackrel{-}{\longrightarrow} L \longrightarrow 1 \tag{2.1}$$

determined by the commutator condition $aba^{-1}b^{-1} = \kappa_6^{c_0(\bar{a},\bar{b})}$ splits. Then for each $\alpha \in L$, we can choose an element e^{α} of \hat{L} so that $e^{\alpha}e^{\beta} = e^{\alpha+\beta}$. The twisted group algebra $\mathbb{C}\{L\}$ is isomorphic to the ordinary group algebra $\mathbb{C}[L]$.

We adopt the same notation as in [9] to denote cosets of L in the dual lattice $L^{\perp} = \{\alpha \in \mathbb{Q} \otimes_{\mathbb{Z}} L \mid \langle \alpha, L \rangle \subset \mathbb{Z} \}$, namely,

$$L^0 = L, \quad L^1 = \frac{-\beta_1 + \beta_2}{3} + L, \quad L^2 = \frac{\beta_1 - \beta_2}{3} + L,$$

$$L_0 = L, \quad L_a = \frac{\beta_2}{2} + L, \quad L_b = \frac{\beta_0}{2} + L, \quad L_c = \frac{\beta_1}{2} + L,$$

$$L^{(i,j)} = L_i + L^j$$

for i=0,a,b,c and j=0,1,2, where $\{0,a,b,c\}\cong\mathbb{Z}_2\times\mathbb{Z}_2$. Then, $L^{(i,j)}, i\in\{0,a,b,c\}, j\in\{0,1,2\}$ are all the cosets of L in L^\perp and $L^\perp/L\cong\mathbb{Z}_2\times\mathbb{Z}_2\times\mathbb{Z}_3$.

Our notation for the vertex operator algebra $(V_L, Y(\cdot, z))$ associated with L is standard [8]. In particular, $\mathfrak{h} = \mathbb{C} \otimes_{\mathbb{Z}} L$ is an abelian Lie algebra, $\hat{\mathfrak{h}} = \mathfrak{h} \otimes \mathbb{C}[t, t^{-1}] \oplus \mathbb{C}c$ is the corresponding affine Lie algebra, $M(1) = \mathbb{C}[\alpha(n); \alpha \in \mathfrak{h}, n < 0]$, where $\alpha(n) = \alpha \otimes t^n$, is the unique irreducible $\hat{\mathfrak{h}}$ -module such that $\alpha(n)1 = 0$ for all $\alpha \in \mathfrak{h}$ and n > 0, and c = 1. As a vector space $V_L = M(1) \otimes \mathbb{C}[L]$ and for each $v \in V_L$, a vertex operator $Y(v,z) = \sum_{n \in \mathbb{Z}} v_n z^{-n-1} \in \operatorname{End}(V_L)[[z,z^{-1}]]$ is defined. The vector $\mathbf{1} = 1 \otimes 1$ is called the vacuum vector.

There are exactly 12 isomorphism classes of irreducible V_L -modules, which are represented by $V_{L^{(i,j)}}$, i=0,a,b,c and j=0,1,2. We use the symbol e^{α} , $\alpha \in L^{\perp}$ to denote a basis of $\mathbb{C}\{L^{\perp}\}$.

Let

$$x(\alpha) = e^{\sqrt{2}\alpha} + e^{-\sqrt{2}\alpha}, \qquad y(\alpha) = e^{\sqrt{2}\alpha} - e^{-\sqrt{2}\alpha}, \qquad w(\alpha) = \frac{1}{2}\alpha(-1)^2 - x(\alpha)$$

for $\alpha \in \{\pm \alpha_0, \pm \alpha_1, \pm \alpha_2\}$ and set

$$\omega=rac{1}{5}ig(w(lpha_1)+w(lpha_2)+w(lpha_0)ig), \ \widetilde{\omega}=rac{1}{6}ig(lpha_1(-1)^2+lpha_2(-1)^2+lpha_0(-1)^2ig), \ \omega^1=rac{1}{4}w(lpha_1), \qquad \omega^2=\omega-\omega^1, \qquad \omega^3=\widetilde{\omega}-\omega.$$

Then $\widetilde{\omega}$ is the Virasoro element of V_L and $\omega^1, \omega^2, \omega^3$ are mutually orthogonal conformal vectors of central charge 1/2, 7/10, 4/5 respectively (cf.[4]). The subalgebra $Vir(\omega^i)$ generated by ω^i is isomorphic to the Virasoro vertex operator algebra of given central charge.

We study certain subalgebras, and also submodules for them in V_{L_i} , i = 0, a, b, c and in V_{L_j} , j = 0, 1, 2. Set

$$\begin{split} M_k^i &= \{v \in V_{L_i} \mid (\omega^3)_1 v = 0\}, \\ W_k^i &= \{v \in V_{L_i} \mid (\omega^3)_1 v = \frac{2}{5} v\}, \qquad \text{for } i = 0, a, b, c, \\ \\ M_t^j &= \{v \in V_{L^j} \mid (\omega^1)_1 v = (\omega^2)_1 v = 0\}, \\ W_t^j &= \{v \in V_{L^j} \mid (\omega^1)_1 v = 0, \quad (\omega^2)_1 v = \frac{3}{5} v\}, \qquad \text{for } j = 0, 1, 2. \end{split}$$

Then M_k^0 and M_t^0 are simple vertex operator algebras. Furthermore, $\{M_k^i, W_k^i, i = 0, a, b, c\}$ and $\{M_t^j, W_t^j, j = 0, 1, 2\}$ are the sets of all inequivalent irreducible modules for M_k^0 and M_t^0 , respectively [9, 11, 12]. We have

$$M_k^0 \cong L(\frac{1}{2},0) \otimes L(\frac{7}{10},0) \oplus L(\frac{1}{2},\frac{1}{2}) \otimes L(\frac{7}{10},\frac{3}{2}), \qquad M_t^0 \cong L(\frac{4}{5},0) \oplus L(\frac{4}{5},3).$$

We also note that

$$V_{L^{(i,j)}} \cong \left(M_k^i \otimes M_t^j\right) \oplus \left(W_k^i \otimes W_t^j\right)$$

as an $M_k^0 \otimes M_t^0$ -module.

We consider the following three isometries of $(L, \langle \cdot, \cdot \rangle)$:

$$\tau: \beta_1 \to \beta_2 \to \beta_0 \to \beta_1,$$

$$\sigma: \beta_1 \to \beta_2, \qquad \beta_2 \to \beta_1,$$

$$\theta: \beta_i \to -\beta_i, \quad i = 1, 2.$$

Set $M=M_k^0$, which is invariant under τ and σ . Moreover, θ acts as the identity on M. We are interested in the subalgebra M^{τ} of the fixed points of τ in M. The weight 2 subspace of M^{τ} is spanned by ω , which is the Virasoro element of M with central charge 6/5. There are nontrivial relations among $w(\alpha_i)_0 w(\alpha_j)$, $i, j \in \{0, 1, 2\}$. For example,

$$w(\alpha_1)_0 w(\alpha_2) - w(\alpha_2)_0 w(\alpha_1)$$

= $w(\alpha_2)_0 w(\alpha_0) - w(\alpha_0)_0 w(\alpha_2)$
= $w(\alpha_0)_0 w(\alpha_1) - w(\alpha_1)_0 w(\alpha_0)$.

Set $J = w(\alpha_1)_0 w(\alpha_2) - w(\alpha_2)_0 w(\alpha_1)$. Then $\tau J = J$, $\sigma J = -J$ and $\theta J = J$. The weight 3 subspace of M^{τ} is of dimension 2 and it is spanned by $\omega_0 \omega$ and J. We have $\omega_1 J = 3J$ and $\omega_n J = 0$ for $n \geq 2$. Thus J is a highest weight vector for $Vir(\omega)$. Let $L(n) = \omega_{n+1}$ and $J(n) = J_{n+2}$. By a direct calculation, we have

$$[L(m), L(n)] = (m-n)L(m+n) + \frac{m^3 - m}{12} \cdot \frac{6}{5} \cdot \delta_{m+n,0}, \tag{2.2}$$

$$[L(m), J(n)] = (2m - n)J(m + n), \tag{2.3}$$

$$[J(m), J(n)] = (m-n) \Big(22(m+n+2)(m+n+3) + 35(m+2)(n+2) \Big) L(m+n)$$

$$-120(m-n) \Big(\sum_{k \le -2} L(k)L(m+n-k) + \sum_{k \ge -1} L(m+n-k)L(k) \Big)$$

$$-\frac{7}{10} m(m^2-1)(m^2-4)\delta_{m+n,0}.$$

$$(2.4)$$

Let $L_n = L(n)$ and $W_n = \sqrt{-1/210}J(n)$. Then the commutation relations in the above theorem coincide with the commutation relations (2.1) and (2.2) of [1]. Thus

Theorem 2.1 W is a W_3 algebra of central charge 6/5.

3 Irreducible modules for \mathcal{W}

First, we consider irreducible modules for M^{τ} . Let (U, Y_U) be one of the 8 irreducible M-modules. Following [3], we consider a new M-module $(U \circ \tau, Y_{U \circ \tau})$ such that $U \circ \tau = U$ as vector spaces and $Y_{U \circ \tau}(v, z) = Y_U(\tau v, z)$ for $v \in M$. Then $U \mapsto U \circ \tau$ induces a permutation on the set of irreducible M-modules. We can easily verify that

Lemma 3.1 (1) $M_k^0 \circ \tau = M_k^0$ and $W_k^0 \circ \tau = W_k^0$.

- (2) $M_{k}^{a} \circ \tau = M_{k}^{c}$, $M_{k}^{c} \circ \tau = M_{k}^{b}$, and $M_{k}^{b} \circ \tau = M_{k}^{a}$. (3) $W_{k}^{a} \circ \tau = W_{k}^{c}$, $W_{k}^{c} \circ \tau = W_{k}^{b}$, and $W_{k}^{b} \circ \tau = W_{k}^{a}$.

For any au-invariant space U, set $U(\epsilon) = \{u \in U \mid \tau u = \xi^{\epsilon}u\}, \ \epsilon = 0, 1, 2$, where $\xi = \exp(2\pi\sqrt{-1}/3)$. Thus $U(0) = U^{\tau}$ and $M(\epsilon) = \{v \in M_k^0 \mid \tau v = \xi^{\epsilon}v\}$. Likewise, set $W(\epsilon) = \{v \in W_k^0 \mid \tau v = \xi^{\epsilon}v\}$. From Lemma 3.1 and [6, Theorem 6.14], we see that $M(\epsilon)$ and $W(\epsilon)$ are inequivalent irreducible M^{τ} -modules for $\epsilon = 0, 1, 2$. Moreover, $M_k^i, i = a, b, c$ are equivalent irreducible M^{τ} -modules and that $W_k^i, i = a, b, c$ are also equivalent irreducible M^{τ} -modules by [6, Theorem 6.14]. Hence we obtain 8 inequivalent irreducible M^{τ} -modules. Those irreducible modules with their top levels and the action of L(0) and J(0) are collected in Table 1.

Table 1: irreducible M^{τ} -modules in M_{i}^{i} and W_{i}^{i}

| Table 1. inteducible 191 modules in 191 _k and 77 _k | | | |
|--|---|----------------|----------------|
| irred. module | top level | L(0) | J(0) |
| M(0) | C1 | 0 | 0 |
| M(1) | $\mathbb{C}u^1$ | 2 | $-12\sqrt{-3}$ |
| M(2) | $\mathbb{C}u^2$ | 2 | $12\sqrt{-3}$ |
| W(0) | $ig \mathbb{C}(y(lpha_1) + y(lpha_2) + y(lpha_0))$! | <u>8</u> 5 | 0 |
| W(1) | $\mathbb{C}(lpha_1(-1) - \xilpha_2(-1))$ | <u>3</u> | $2\sqrt{-3}$ |
| W(2) | $\mathbb{C}(\alpha_1(-1) - \xi^2 \alpha_2(-1))$ | യ ന്തവിതവിത | $-2\sqrt{-3}$ |
| M_k^c | $\mathbb{C}(e^{eta_1/2}-e^{-eta_1/2})$ | $\frac{1}{2}$ | 0 |
| W_k^c | $\mathbb{C}(e^{\beta_1/2}+e^{-\beta_1/2})$ | $\frac{1}{10}$ | 0 |

We now study irreducible τ -twisted (resp. τ^2 -twisted) M-modules. The argument here is similar to that in [10, Section 6]. Basic references to twisted modules for lattice vertex operator algebras are [2, 13]. We follow [2] with $L = \sqrt{2}A_2$, p = 3, q = 6, and $\nu = \tau$. Let $\mathfrak{h} = \mathbb{C} \otimes_{\mathbb{Z}} L$ and extend the \mathbb{Z} -bilinear form $\langle \cdot, \cdot \rangle$ on L to \mathfrak{h} linearly. Set

$$h_1 = \frac{1}{3}(\beta_1 + \xi^2 \beta_2 + \xi \beta_0), \qquad h_2 = \frac{1}{3}(\beta_1 + \xi \beta_2 + \xi^2 \beta_0).$$

For $n \in \mathbb{Z}$, set $\mathfrak{h}_{(n)} = \{\alpha \in \mathfrak{h} \mid \tau \alpha = \xi^n \alpha\}$. Since τ is fixed-point-free on L, it follows that $\mathfrak{h}_{(0)}=0$. Furthermore, $\mathfrak{h}_{(1)}=\mathbb{C}h_1$ and $\mathfrak{h}_{(2)}=\mathbb{C}h_2$. For $\alpha\in\mathfrak{h}$, we denote by $\alpha_{(n)}$ the component of α in $\mathfrak{h}_{(n)}$. Thus $(\beta_i)_{(1)} = \xi^{i-1}h_1$ and $(\beta_i)_{(2)} = \xi^{2(i-1)}h_2$ for i = 0, 1, 2.

Define the τ -twisted affine Lie algebra to be

$$\hat{\mathfrak{h}}[au] = \Big(\bigoplus_{n \in \mathbb{Z}} \mathfrak{h}_{(n)} \otimes t^{n/3}\Big) \oplus \mathbb{C}c$$

with the bracket

$$[x \otimes t^m, y \otimes t^n] = m\langle x, y \rangle \delta_{m+n,0} c$$

for $x \in \mathfrak{h}_{(3m)}$, $y \in \mathfrak{h}_{(3n)}$, $m, n \in (1/3)\mathbb{Z}$, and $[c, \hat{\mathfrak{h}}[g]] = 0$. The isometry τ acts on $\hat{\mathfrak{h}}[\tau]$ by $\tau(x \otimes t^m) = \xi^m x \otimes t^m$ and $\tau(c) = c$. Set

$$\hat{\mathfrak{h}}[\tau]^+ = \bigoplus_{n>0} \mathfrak{h}_{(n)} \otimes t^{n/3}, \quad \hat{\mathfrak{h}}[\tau]^- = \bigoplus_{n<0} \mathfrak{h}_{(n)} \otimes t^{n/3}, \quad \text{and} \quad \hat{\mathfrak{h}}[\tau]^0 = \mathbb{C}c$$

and consider the $\hat{\mathfrak{h}}[\tau]$ -module

$$S[au] = U(\hat{\mathfrak{h}}[au]) \otimes_{U(\hat{\mathfrak{h}}[au]^+ \oplus \hat{\mathfrak{h}}[au]^0)} \mathbb{C}$$

induced from the $\hat{\mathfrak{h}}[\tau]^+ \oplus \hat{\mathfrak{h}}[\tau]^0$ -module \mathbb{C} , where $\hat{\mathfrak{h}}[\tau]^+$ acts trivially on \mathbb{C} and c acts as 1 on \mathbb{C} . We define the weight in $S[\tau]$ by

$$\operatorname{wt}(x\otimes t^n)=-n \quad \text{and} \quad \operatorname{wt} 1=\frac{1}{9},$$

where $n \in (1/3)\mathbb{Z}$ and $x \in \mathfrak{h}_{(3n)}$ (cf. [2, (4.6), (4.10)]).

For $\alpha \in \mathfrak{h}$ and $n \in (1/3)\mathbb{Z}$, denote by $\alpha(n)$ the operator on $S[\tau]$ induced by $\alpha_{(3n)} \otimes t^n$. Then, as a vector space $S[\tau]$ can be identified with a polynomial algebra with variables $h_1(1/3+n)$ and $h_2(2/3+n)$, $n \in \mathbb{Z}$. The weight of the operator $h_j(j/3+n)$ is -j/3-n. The alternating \mathbb{Z} -bilinear map $c_0^{\tau}: L \times L \to \mathbb{Z}/6\mathbb{Z}$ defined by [2, (2.10)] is such that

$$c_0^{\tau}(\alpha,\beta) = \sum_{r=0}^{2} (3+2r)\langle \tau^r \alpha, \beta \rangle + 6\mathbb{Z}.$$

In our case $\sum_{r=0}^{2} \tau^{r} \alpha = 0$, since τ is fixed-point-free on L. Moreover,

$$\sum_{r=0}^{2} r \langle \tau^{r} \beta_{i}, \beta_{j} \rangle = \begin{cases} \pm 6 & \text{if } \tau \beta_{i} \neq \beta_{j} \\ 0 & \text{if } \tau \beta_{i} = \beta_{j}. \end{cases}$$

Hence $c_0^{\tau}(\alpha,\beta)=0$ for all $\alpha,\beta\in L$. This means that the central extension

$$1 \longrightarrow \langle \kappa_6 \rangle \longrightarrow \hat{L}_\tau \stackrel{-}{\longrightarrow} L \longrightarrow 1 \tag{3.1}$$

determined by the commutator condition $aba^{-1}b^{-1} = \kappa_6^{c_0^{\tau}(\bar{a},\bar{b})}$ splits.

We consider the relation between two central extensions \hat{L} of (2.1) and \hat{L}_{τ} of (3.1). Since both of \hat{L} and \hat{L}_{τ} are split extensions, we use the same symbol e^{α} to denote both of an element in \hat{L} and an element in \hat{L}_{τ} which correspond naturally to $\alpha \in L$. Actually, in Section 2 we choose $e^{\alpha} \in \hat{L}$ so that the multiplication in \hat{L} is $e^{\alpha} \times e^{\beta} = e^{\alpha+\beta}$. Also we can choose $e^{\alpha} \in \hat{L}_{\tau}$ such that the multiplication $e^{\alpha} \times_{\tau} e^{\beta}$ in \hat{L}_{τ} is related to the multiplication in \hat{L} by (cf. [2, (2.4)])

$$e^{\alpha} \times e^{\beta} = \kappa_6^{\epsilon_0(\alpha,\beta)} e^{\alpha} \times_{\tau} e^{\beta},$$
 (3.2)

where the \mathbb{Z} -linear map $\varepsilon_0: L \times L \to \mathbb{Z}/6\mathbb{Z}$ is defined by [2, (2.13)]. In our case

$$\varepsilon_0(\alpha, \beta) = -\langle \tau^{-1}\alpha, \beta \rangle + 6\mathbb{Z}.$$
(3.3)

As in Section 2, we usually write $e^{\alpha}e^{\beta}=e^{\alpha+\beta}$ to denote the product of e^{α} and e^{β} in \hat{L} . The automorphism τ of L lifts to an automorphism $\hat{\tau}$ of \hat{L} such that $\hat{\tau}(e^{\alpha})=e^{\tau\alpha}$ and $\hat{\tau}(\kappa_6)=\kappa_6$. Since ε_0 is τ -invariant, we can also think $\hat{\tau}$ to be an automorphism of \hat{L}_{τ} in a similar way. By abuse of notation we denote $\hat{\tau}$ by simply τ also.

We have $(1-\tau)L = \operatorname{span}_{\mathbb{Z}}\{\beta_1 - \beta_2, \beta_1 + 2\beta_2\}$. The quotient group $L/(1-\tau)L$ is of order 3 and generated by $\beta_1 + (1-\tau)L$. Now $K = \{a^{-1}\tau(a) \mid a \in \hat{L}_{\tau}\}$ is a central subgroup of \hat{L}_{τ} with $\overline{K} = (1-\tau)L$ and $K \cap \langle \kappa_6 \rangle = 1$. Here note that a^{-1} is the inverse of a in \hat{L}_{τ} and $a^{-1}\tau(a)$ is the product $a^{-1} \times_{\tau} \tau(a)$ in \hat{L}_{τ} . In \hat{L}_{τ} we can verify that

$$e^{3\beta_1} = (e^{\beta_0 - \beta_1})^{-1} \times_{\tau} \tau(e^{\beta_0 - \beta_1}) \in K.$$

Since $\kappa_3 e^{\beta_1} \times_{\tau} \kappa_3 e^{\beta_1} \times_{\tau} \kappa_3 e^{\beta_1} = e^{3\beta_1}$ and $\kappa_3 e^{\beta_1} \times_{\tau} \kappa_3 e^{-\beta_1} = 1$, it follows that

$$\hat{L}_{\tau}/K = \{K, \kappa_3 e^{\beta_1} K, \kappa_3 e^{-\beta_1} K\} \times \langle \kappa_6 \rangle K/K \cong \mathbb{Z}_3 \times \mathbb{Z}_6.$$

For j = 0, 1, 2, define a linear character $\chi_j : \hat{L}_\tau/K \to \mathbb{C}^\times$ by

$$\chi_j(\kappa_6) = \xi_6, \quad \chi_j(\kappa_3 e^{\beta_1} K) = \xi^j, \quad ext{and} \quad \chi_j(\kappa_3 e^{-\beta_1} K) = \xi^{-j},$$

where $\xi_6 = \exp(2\pi\sqrt{-1}/6)$. Let T_{χ_j} be the one dimensional \hat{L}_{τ}/K -module affording the character χ_j . As an \hat{L}_{τ} -module, K acts trivially on T_{χ_j} . Since $\sum_{r=0}^2 \tau^r \alpha = 0$ for $\alpha \in L$, those T_{χ_j} , j=0,1,2, are the irreducible \hat{L}_{τ} -modules constructed in [13, Section 6]. Let

$$V_L^{T_{m{\chi}_j}} = V_L^{T_{m{\chi}_j}}(au) = S[au] \otimes T_{m{\chi}_j}$$

and define the τ -twisted vertex operator $Y^{\tau}(\cdot,z):V_L\to \operatorname{End}(V_L^{T_{\chi_j}}(\tau))\{z\}$ as in [2]. We extend the action of τ to $V_L^{T_{\chi_j}}(\tau)$ so that τ is the identity on T_{χ_j} . The weight of every element in T_{χ_j} is defined to be 0.

By [2, Theorem 7.1], $(V_L^{T_{\chi_j}}(\tau), Y^{\tau}(\cdot, z))$, j=0,1,2 are inequivalent irreducible τ -twisted V_L -modules. Now among the 12 irreducible V_L -modules $V_{L^{(i,j)}}$, $i\in\{0,a,b,c\}$ and $j\in\{0,1,2\}$, the τ -stable irreducible modules are $V_{L^{(0,j)}}$, $j\in\{0,1,2\}$. Hence by [3, Theorem 10.2], $(V_L^{T_{\chi_j}}(\tau), Y^{\tau}(\cdot, z))$, j=0,1,2, are all the inequivalent irreducible τ -twisted V_L -modules. The isometry θ of $(L, \langle \cdot, \cdot \rangle)$ induces a permutation on $V_L^{T_{\chi_j}}(\tau)$, j=0,1,2. In fact, the permutation leaves $V_L^{T_{\chi_0}}(\tau)$ invariant and interchanges $V_L^{T_{\chi_1}}(\tau)$ and $V_L^{T_{\chi_2}}(\tau)$. Since $M^{\tau} \otimes M_t^0$ is contained in the subalgebra $(V_L)^{\tau}$ of fixed points of τ in V_L , we can deal with $(V_L^{T_{\chi_j}}(\tau), Y^{\tau}(\cdot, z))$ as an $M^{\tau} \otimes M_t^0$ -module.

The decomposition of $V_L^{T_{\chi_j}}(\tau)$ as a τ -twisted $M \otimes M_t^0$ -module was studied in [10]. Set

$$\begin{split} &M_T^0(\tau) = \{u \in V_L^{T_{\chi_0}}(\tau) \mid (\omega^3)_1 u = 0\}, \\ &W_T^0(\tau) = \{u \in V_L^{T_{\chi_0}}(\tau) \mid (\omega^3)_1 u = \frac{2}{5}u\}, \\ &M_T^j(\tau) = \{u \in V_L^{T_{\chi_j}}(\tau) \mid (\omega^3)_1 u = \frac{2}{3}u\}, \quad j = 1, 2, \\ &W_T^j(\tau) = \{u \in V_L^{T_{\chi_j}}(\tau) \mid (\omega^3)_1 u = \frac{1}{15}u\}, \quad j = 1, 2. \end{split}$$

Then, by [10, Proposition 6.8], $M_T^j(\tau)$ and $W_T^j(\tau)$, j=0,1,2, are irreducible τ -twisted M-modules. Furthermore, for j=0,1,2,

$$V_L^{T_{\chi_j}}(\tau) \cong M_T^j(\tau) \otimes M_t^j \oplus W_T^j(\tau) \otimes W_t^j$$

as τ -twisted $M \otimes M_t^0$ -modules.

There are at most two inequivalent irreducible τ -twisted M-modules by Lemma 4.1 and [3, Theorem 10.2]. Then, looking at the smallest weight of $M_T^j(\tau)$ and $W_T^j(\tau)$, we have that $M_T^0(\tau) \cong M_T^1(\tau) \cong M_T^2(\tau)$ and $W_T^0(\tau) \cong W_T^1(\tau) \cong W_T^2(\tau)$ and that $M_T^0(\tau) \not\cong W_T^0(\tau)$ as τ -twisted M-modules. We denote $M_T^0(\tau)$ by $M_T(\tau)$ and $M_T^0(\tau)$ by $M_T(\tau)$. We conclude that there are exactly two inequivalent irreducible τ -twisted M-modules, which are represented by $M_T(\tau)$ and $M_T(\tau)$. As τ -twisted $M \otimes \operatorname{Vir}(\omega^3)$ -modules, we have

$$V_L^{T_{\mathbf{x_0}}}(\tau) \cong M_T(\tau) \otimes \left(L(\frac{4}{5}, 0) + L(\frac{4}{5}, 3)\right) \oplus W_T(\tau) \otimes \left(L(\frac{4}{5}, \frac{2}{5}) + L(\frac{4}{5}, \frac{7}{5})\right), \tag{3.4}$$

$$V_L^{T_{\chi_1}}(\tau) \cong V_L^{T_{\chi_2}}(\tau) \cong M_T(\tau) \otimes L(\frac{4}{5}, \frac{2}{3}) \oplus W_T(\tau) \otimes L(\frac{4}{5}, \frac{1}{15}). \tag{3.5}$$

For $\epsilon = 0, 1, 2$, let

$$M_T(\tau)(\epsilon) = \{u \in M_T(\tau) \mid \tau u = \xi^{\epsilon} u\},$$

 $W_T(\tau)(\epsilon) = \{u \in W_T(\tau) \mid \tau u = \xi^{\epsilon} u\}.$

Those 6 modules for M^{τ} are inequivalent irreducible modules by [14, Theorem 2]. Their top levels are of dimension one. Those top levels and the eigenvalues for the action of $L^{\tau}(0) = \omega_1$ and $J^{\tau}(0) = J_2$ are collected in Table 2.

By a similar argument we obtain 6 inequivalent irreducible M^{τ} -modules inside τ^2 -twisted M-modules. The results are collected in Table 3.

We have obtained 20 irreducible M^{τ} -modules. Now we turn to the vertex operator algebra W generated by ω and J in M^{τ} . Note that

$$L(n)\mathbf{1} = 0 \text{ for } n \ge -1,$$
 $L(-2)\mathbf{1} = \omega,$ (3.6)

$$J(n)\mathbf{1} = 0 \text{ for } n > -2,$$
 $J(-3)\mathbf{1} = J,$ (3.7)

Then by using commutation relations (2.2), (2.3), and (2.4), we see that

Table 2: irreducible M^{τ} -modules in $M_{T}(\tau)$ and $W_{T}(\tau)$ irred. module | top level | $L^{\tau}(0)$ | $J^{\tau}(0)$ | $M_{T}(\tau)(0)$ | $C1 \otimes v$ | $\frac{1}{9}$ | $\frac{14}{81}\sqrt{-3}$ | $M_{T}(\tau)(1)$ | $Ch_{2}(-\frac{1}{3})^{2} \otimes v$ | $\frac{1}{9} + \frac{2}{3}$ | $-\frac{238}{81}\sqrt{-3}$ | $M_{T}(\tau)(2)$ | $C(\frac{4}{3}h_{1}(-\frac{2}{3})^{2} \otimes v + h_{2}(-\frac{1}{3})^{4} \otimes v)$ | $\frac{1}{9} + \frac{4}{3}$ | $\frac{374}{81}\sqrt{-3}$ | $W_{T}(\tau)(0)$ | $Ch_{2}(-\frac{1}{3}) \otimes v$ | $\frac{2}{45}$ | $-\frac{4}{81}\sqrt{-3}$ | $W_{T}(\tau)(1)$ | $Ch_{1}(-\frac{2}{3}) \otimes v$ | $\frac{2}{45} + \frac{1}{3}$ | $-\frac{22}{81}\sqrt{-3}$ | $-\frac{22}{81}\sqrt{-3}$ | $W_{T}(\tau)(2)$ | $Ch_{2}(-\frac{1}{3})^{3} \otimes v$ | $\frac{2}{45} + \frac{2}{3}$ | $\frac{176}{81}\sqrt{-3}$ | $-\frac{27}{81}\sqrt{-3}$ | $-\frac{27}{81}\sqrt$

Table 3: irreducible M^{τ} -modules in $M_{T}(\tau^{2})$ and $W_{T}(\tau^{2})$ irred. module | top level | $L^{\tau^{2}}(0)$ | $J^{\tau^{2}}(0)$ | $M_{T}(\tau^{2})(0)$ | $\mathbb{C}1 \otimes v$ | $\frac{1}{9}$ | $-\frac{14}{81}\sqrt{-3}$ | $M_{T}(\tau^{2})(1)$ | $\mathbb{C}h'_{2}(-\frac{1}{3})^{2} \otimes v$ | $\frac{1}{9} + \frac{2}{3}$ | $\frac{238}{81}\sqrt{-3}$ | $M_{T}(\tau^{2})(2)$ | $\mathbb{C}(\frac{4}{3}h'_{1}(-\frac{2}{3})^{2} \otimes v + h'_{2}(-\frac{1}{3})^{4} \otimes v)$ | $\frac{1}{9} + \frac{4}{3}$ | $-\frac{374}{81}\sqrt{-3}$ | $W_{T}(\tau^{2})(0)$ | $\mathbb{C}h'_{2}(-\frac{1}{3}) \otimes v$ | $\frac{2}{45}$ | $\frac{4}{81}\sqrt{-3}$ | $\frac{4}{81}\sqrt{-3}$ | $W_{T}(\tau^{2})(1)$ | $\mathbb{C}h'_{1}(-\frac{2}{3}) \otimes v$ | $\frac{2}{45} + \frac{1}{3}$ | $\frac{22}{81}\sqrt{-3}$ | $\frac{2}{45} + \frac{1}{3}$ | $\frac{22}{81}\sqrt{-3}$ | $\frac{2}{45} + \frac{1}{3}$ | $\frac{276}{81}\sqrt{-3}$ | $\frac{176}{81}\sqrt{-3}$ | \frac

Lemma 3.2 W is spanned by the vectors of the form

$$L(-m_1)\cdots L(-m_p)J(-n_1)\cdots J(-n_q)\mathbf{1}$$
 (3.8)

with $m_1 \ge \cdots \ge m_p \ge 2$, $n_1 \ge \cdots \ge n_q \ge 3$, $p = 0, 1, 2, \ldots$, and $q = 0, 1, 2, \ldots$

A vector $v \in \mathcal{W}$ of weight h is called a singular vector for \mathcal{W} if it satisfies

- (1) L(0)v = hv,
- (2) L(n)v = 0 and J(n)v = 0 for $n \ge 1$.

Note that v is not necessarily an eigenvector for J(0). By commutation relations (2.2) and (2.3), it is easy to show that the condition (2) holds if v satisfies

$$(2)' L(1)v = L(2)v = J(1)v = 0.$$

We consider W as a space spanned by the vectors of the form (3.8). The weight of such a vector is $m_1 + \cdots + m_p + n_1 + \cdots + n_q$. Using a computer algebra system Risa/Asir we have the following lemma.

Lemma 3.3 Let v be a linear combination of the vectors of the form (3.8) of weight h. Under the conditions (3.6) and (3.7) and the commutation relations (2.2), (2.3), and (2.4), we have L(1)v = L(2)v = J(1)v = 0 only if v = 0 in the case $h \le 11$. In the case h = 12, there exists a unique, up to scalar multiple, linear combination \mathbf{v}^{12} which satisfies $L(1)\mathbf{v}^{12} = L(2)\mathbf{v}^{12} = J(1)\mathbf{v}^{12} = 0$. The explicit form of \mathbf{v}^{12} is given in Appendix. We also have $J(0)\mathbf{v}^{12} = 0$.

By a general theory of lattice vertex operator algebras, V_L possesses an invariant positive definite hermitian form. From this it follows that

Proposition 3.4 The singular vector $\mathbf{v}^{12} = 0$.

Using the explicit form of \mathbf{v}^{12} , $J(-1)\mathbf{v}^{12}$, $J(-2)\mathbf{v}^{12}$, and $J(-1)\mathbf{v}^{12}$, we can determine the Zhu algebra $Z(\mathcal{W})$ of \mathcal{W} . The standard reference to Zhu's theory is [15, Section 2]. The Zhu algebra $Z(\mathcal{W})$ is a quotient vector space $A(\mathcal{W}) = \mathcal{W}/O(\mathcal{W})$ equipped with a commutative associative algebra structure with respect to an operation *. We denote the image of a vector $v \in \mathcal{W}$ in $A(\mathcal{W}) = \mathcal{W}/O(\mathcal{W})$ by [v]. We can define an algebra homomorphism from $\mathbb{C}[x,y]$ onto $A(\mathcal{W})$ by $x \mapsto [\omega]$ and $y \mapsto [J]$. The primary decomposition of its kernel \mathcal{I} is such that $\mathcal{I} = \bigcap_{i=1}^{20} \mathcal{P}_i$, where \mathcal{P}_i , $1 \le i \le 20$ are

$$\begin{array}{lll} \langle x,\,y\rangle, & \langle 5x-8,\,y\rangle, \\ \langle 2x-1,\,y\rangle, & \langle 10x-1,\,y\rangle, \\ \langle x-2,\,y-12\sqrt{-3}\rangle, & \langle x-2,\,y+12\sqrt{-3}\rangle, \\ \langle 5x-3,\,y-2\sqrt{-3}\rangle, & \langle 5x-3,\,y+2\sqrt{-3}\rangle, \\ \langle 9x-1,\,81y-14\sqrt{-3}\rangle, & \langle 9x-1,\,81y+14\sqrt{-3}\rangle, \\ \langle 9x-7,\,81y-238\sqrt{-3}\rangle, & \langle 9x-7,\,81y+238\sqrt{-3}\rangle, \\ \langle 9x-13,\,81y-374\sqrt{-3}\rangle, & \langle 9x-13,\,81y+374\sqrt{-3}\rangle, \\ \langle 45x-2,\,81y-4\sqrt{-3}\rangle, & \langle 45x-2,\,81y+4\sqrt{-3}\rangle, \\ \langle 45x-17,\,81y-22\sqrt{-3}\rangle, & \langle 45x-17,\,81y+22\sqrt{-3}\rangle, \\ \langle 45x-32,\,81y-176\sqrt{-3}\rangle, & \langle 45x-32,\,81y+176\sqrt{-3}\rangle, \end{array}$$

These primary ideals correspond to the 20 irreducible M^{τ} -modules listed in Tables 1, 2, and 3. The correspondence is given by substituting x and y with the eigenvalues for L(0) and J(0) on the top levels of 20 irreducible modules. The eigenvalues are the zeros of those primary ideals.

By Zhu's theory we obtain the classification of irreducible W-modules, namely,

Theorem 3.5 (1) $M^{\tau} = \mathcal{W}$.

- (2) $A(M^{\tau}) \cong \bigoplus_{i=1}^{20} \mathbb{C}[x,y]/\mathcal{P}_i$ is a 20-dimensional commutative associative algebra.
- (3) There are exactly 20 inequivalent irreducible M^{τ} -modules. Their representatives are listed in Tables 1, 2, and 3 in Section 2, namely, $M(\epsilon)$, $W(\epsilon)$, M_k^c , W_k^c , $M_T(\tau^i)(\epsilon)$, and $W_T(\tau^i)(\epsilon)$ for $\epsilon = 0, 1, 2$ and i = 1, 2.

Appendix

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\mathbf{v}^{12} = -(5877264800/3501)L(-12)\mathbf{1} + (3404072000/3501)L(-10)L(-2)\mathbf{1}
       -(2653990000/3501)L(-9)L(-3)\mathbf{1} - (266376800/3501)L(-8)L(-4)\mathbf{1}
      \hspace*{35pt} + (282988000/1167)L(-8)L(-2)^2 {\bf 1} - (23744800/1167)L(-7)L(-5) {\bf 1}
       -(30824000/1167)L(-7)L(-3)L(-2)\mathbf{1} + (1242377600/1167)L(-6)^2\mathbf{1}
      -(61947200/3501)L(-6)L(-4)L(-2)\mathbf{1} - (1313806000/1167)L(-6)L(-3)^2\mathbf{1}
      -(45496000/1167)L(-6)L(-2)^3\mathbf{1} - (3046768400/3501)L(-5)^2L(-2)\mathbf{1}
      +(299424800/1167)L(-5)L(-4)L(-3)\mathbf{1}+(2347094000/3501)L(-5)L(-3)L(-2)^2\mathbf{1}
      -(17280400/1167)L(-4)^3\mathbf{1} - (2036373200/3501)L(-4)^2L(-2)^2\mathbf{1}
      +(82996000/3501)L(-4)L(-3)^{2}L(-2)\mathbf{1}+(1074512000/3501)L(-4)L(-2)^{4}\mathbf{1}
      +(511628125/3501)L(-3)^4\mathbf{1}-(418850000/3501)L(-3)^2L(-2)^3\mathbf{1}
      \hspace*{35pt} - (59680000/3501)L(-2)^{6} \mathbf{1} - (505200/389)L(-6)J(-3)^{2} \mathbf{1}
      +(3380480/1167)L(-4)L(-2)J(-3)^2\mathbf{1}+1150L(-3)^2J(-3)^2\mathbf{1}
      \hspace*{35pt} - (184400/1167)L(-2)^3J(-3)^2{\bf 1} + (3788680/1167)L(-5)J(-4)J(-3){\bf 1}
      -\left(8788400/3501\right)L(-3)L(-2)J(-4)J(-3)\mathbf{1}-(12761440/3501)L(-4)J(-5)J(-3)\mathbf{1}
      -(5727500/10503)L(-4)J(-4)^2\mathbf{1} + (352400/389)L(-2)^2J(-5)J(-3)\mathbf{1}
      + (5727500/10503)L(-2)^2J(-4)^2\mathbf{1} + (1593900/389)L(-3)J(-6)J(-3)\mathbf{1}
      +\ (12935800/10503)L(-3)J(-5)J(-4)\mathbf{1} + (4108000/3501)L(-2)J(-7)J(-3)\mathbf{1}
      -(2811800/1167)L(-2)J(-6)J(-4)1 - (3131600/10503)L(-2)J(-5)^21
      -(14904160/3501)J(-9)J(-3)\mathbf{1} + (32677600/10503)J(-8)J(-4)\mathbf{1}
      +(9423200/10503)J(-7)J(-5)1+(2432375/1167)J(-6)^21
      +J(-3)^41.
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References

- [1] P. Bouwknegt, J. McCarthy and K. Pilch, *The W*₃ Algebra, Lecture Notes in Physics, **m42**, Springer, Berlin 1996.
- [2] C. Dong and J. Lepowsky, The algebraic structure of relative twisted vertex operators, J. Pure and Applied Algebra 110(1996), 259–295.
- [3] C. Dong, H. Li and G. Mason, Modular-invariance of trace functions in orbifold theory and generalized moonshine, *Comm. Math. Phys.* **214** (2000), 1–56.
- [4] C. Dong, H. Li, G. Mason and S. P. Norton, Associative subalgebras of the Griess algebra and related topics, in: *Proc. of the Conference on the Monster and Lie algebras at The Ohio State University, May 1996*, ed. by J. Ferrar and K. Harada, Walter de Gruyter, Berlin-New York, 1998, 27–42.
- [5] C. Dong, C. Lam, K. Tanabe, H. Yamada and K. Yokoyama, \mathbb{Z}_3 symmetry and W_3 algebra in lattice vertex operator algebras, preprint.
- [6] C. Dong and G. Yamskulna, Vertex operator algebras, generalized doubles and dual pairs, *Math. Z.* **241** (2002), 397–423.
- [7] V. A. Fateev and A. B. Zamolodchikov, Conformal quantum field theory models in two dimensions having Z₃ symmetry, *Nuclear Physics* **B280** (1987), 644–660.
- [8] I. B. Frenkel, J. Lepowsky and A. Meurman, Vertex Operator Algebras and the Monster, Pure and Applied Math., Vol. 134, Academic Press, 1988.
- [9] K. Kitazume, C. Lam and H. Yamada, Decomposition of the moonshine vertex operator algebra as Virasoro modules, J. Algebra, 226 (2000), 893–919.
- [10] K. Kitazume, C. Lam and H. Yamada, 3-state Potts model, moonshine vertex operator algebra and 3A elements of the monster group, to appear in *International Mathematics Research Notices*.
- [11] M. Kitazume, M. Miyamoto and H. Yamada, Ternary codes and vertex operator algebras, J. Algebra, 223 (2000), 379–395.
- [12] C. Lam and H. Yamada, $\mathbb{Z}_2 \times \mathbb{Z}_2$ codes and vertex operator algebras, J. Algebra **224** (2000), 268–291.
- [13] J. Lepowsky, Calculus of twisted vertex operators, *Proc. Natl. Acad. Sci. USA* 82 (1985), 8295–8299.
- [14] M. Miyamoto and K. Tanabe, Uniform product of $A_{g,n}(V)$ for an orbifold model V and G-twisted Zhu algebra, math.QA/0112054.
- [15] Y. Zhu, Modular invariance of characters of vertex operator algebras, J. Amer. Math. Soc. 9 (1996), 237-302.