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Asymptotic decay toward the planar rarefaction waves of solutions for viscous conservation laws in several space dimensions

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Abstract

This paper concerns the asymptotic decay rate toward the planar rarefaction waves of the solutions for the scalar viscous conservation laws in two or more space dimensions. This is proved by a result on the decay rate of solutions for one dimensional scalar viscous conservation laws and by using an $L^2$-energy method with a weight of time.

1 Introduction and main result

In this paper, we present the asymptotic decay rate, toward the planar rarefaction waves, of the solutions for scalar viscous conservation laws in two or more space dimensions. Since the proof of the result for the case in more than two dimension will be identical to that for the case in two dimension, we only discuss the equation of the following form:

\[
\begin{align*}
    u_t + (f(u))_x + (g(u))_y &= u_{xx} + u_{yy}, \\
    u(0, x, y) &= u_0(x, y),
\end{align*}
\]

(1) (2)

where $u = u(t, x, y)$ is a scalar function of time $t \geq 0$ and position $(x, y) \in R^2$. We assume that nonlinear flux functions $f$ and $g$ are smooth and also assume that $f$ is convex i.e., for a fixed constant $\alpha > 0$,

\[ f''(u) \geq \alpha. \]

(3)

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The initial condition satisfies
\[ u_0(x, y) \to u_\pm \quad \text{as} \quad x \to \pm \infty, \] (4)
where \( u_\pm \) are constants satisfying \( u_- < u_+ \).

A planar rarefaction wave is a weak solution of the following problem
\[ r_t + (f(r))_x = 0, \] (5)
\[ r(0, x) = r_0(x), \] (6)
where \( r_0(x) \) is given by
\[ r_0(x) = \begin{cases} 
  u_-, & \text{for } x < 0, \\
  u_+, & \text{for } x > 0.
\end{cases} \] (7)

Then \( r(t, x) \) is given explicitly,
\[ r(t, x) = \begin{cases} 
  u_-, & \text{for } x < a(u_-)t, \\
  a^{-1}(x/t), & \text{for } a(u_-)t < x < a(u_+)t, \\
  u_+, & \text{for } a(u_+)t < x,
\end{cases} \] (8)
where \( a = a(u) \) is defined by
\[ a(u) = f'(u). \] (9)

Note
\[ a'(u) \geq \alpha > 0. \] (10)

The stability of rarefaction waves was originally considered by Il'in and Oleinik [3], and has recently been studied by many authors [8, 5, 6, 7, 9, 10].

Harabetian [1] first studied the asymptotic decay rate toward rarefaction waves of the solutions of the scalar viscous conservation laws in one space dimension. Hattori and Nishihara [2] showed a more precise result on the decay rate toward rarefaction waves of solutions to the one dimensional Burgers equation instead of the scalar viscous conservation law in one space dimension. Xin [11] first proved the asymptotic stability of planar rarefaction waves for the several dimensional scalar viscous conservation laws, but the paper [11] did not refer to the decay rate of solutions to (1)-(2) toward rarefaction waves.

In this paper we give the asymptotic convergence rate toward the planar rarefaction wave \( r(t, x) \) of the solution \( u(t, x, y) \) for (1)-(2) in \( L^\infty(R_y; L^2(R_x)) \). To state our result, following Matsumura and Nishihara [6], we introduce the smooth rarefaction wave, which is a smooth solution of the following problem:
\[ w_t + (f(w))_x = 0, \] (11)
\[ w(0, x) = w_0(x), \] (12)
where $w_0(x)$ is defined by
\[ w_0(x) = \frac{u_+ + u_-}{2} + \frac{u_+ - u_-}{2} \kappa \int_0^x (1 + \xi^2)^{-1} d\xi, \tag{13} \]
where
\[ \kappa = \left( \int_0^\infty (1 + \xi^2)^{-1} d\xi \right)^{-1}. \]

We here introduce notations used throughout this paper. For $1 \leq p \leq \infty$, $L^p(R^N)$ is the usual Lebesgue space on $R^N$. For positive integers $m$, $W^{m,p}(R^N)$ is the space of all functions whose weak derivatives up to $m$-th order belong to $L^p(R^N)$. $H^m(R^N)$ denotes $W^{m,2}(R^N)$.

Now we are in position to state our main theorem.

**Theorem 1** Suppose that $u_0 - w_0 \in (H^2 \cap L^1)(R^2)$. Then, there exist positive constants $\delta_0$ and $\delta_0'$ such that if
\[ \|u_0 - w_0\|_{H^2(R^2)} \leq \delta_0 \quad \text{and} \quad |u_+ - u_-| \leq \delta_0', \tag{14} \]
then the problem $(1)-(2)$ has a smooth unique global solution $u(t, x, y)$ satisfying
\[ \sup_{y \in R} \|u(t) - r(t)\|_{L^2(R_x)} \leq C t^{-1/4} \log(2 + t) \quad \text{for } t > 0, \tag{15} \]
where $C$ is a positive constant depending on $u_0$ and $|u_+ - u_-|$.

It is possible to say that our method to prove Theorem 1 is valid for more space dimensional case, for the proof is identical.

The rest of the paper is organized as follows. In Section 2, we show that the original planar rarefaction wave $r(t, x)$ in (8) are approximated by the smooth rarefaction wave $w(t, x)$ in (11) in $L^2(R)$ at the asymptotic rate $O(t^{-1/4})$ as $t \to \infty$. In Section 3, it is shown by using an $L^2$-energy method with polynomial and logarithmic weight of time that the asymptotic behavior of the smooth rarefaction wave $w(t, x)$ in $L^2(R)$ is described by the solution of the viscous scalar conservation law in one space dimension, that is, they converge to each other in $L^2(R)$ at the asymptotic rate $O(t^{-1/4} \log t)$ as $t \to \infty$. Finally, by making use of the result in Sections 2 and 3, we give the proof of Theorem 1. $L^2$-energy method with weight of time also plays an crucial role here.

## 2 Convergence of $w(t, x) - r(t, x)$

Throughout this paper, $C$ denote generic positive constants.

The smooth rarefaction wave has the following properties.
Lemma 1 ([5, 6, 7]) (i) $u_- < w(t, x) < u_+$, $w_x(t, x) > 0$ for $(t, x) \in [0, \infty) \times R$.
(ii) For all $p$ with $1 \leq p \leq \infty$ there is a constant $C_p$ such that
\[
\|w_x(t)\|_{L^p(R)} \leq C_p \min(d, d^1/p(1+1/p)),
\]
\[
\|w_{xx}(t)\|_{L^p(R)} \leq C_p \min(d, d^{-(p-1)/2p(1+(p-1)/2p)}),
\]
for $t > 0$, where
\[
d = u_+ - u_-.
\]
Furthermore we need the following lemma throughout this paper.

Lemma 2 (i) For all $p$ with $1 \leq p \leq \infty$, there is a constant $C_{p,d}$ such that
\[
\|w_{xxx}(t)\|_{L^p(R)} \leq C_{p,d}(1+t)^{-(1+(2p-1)/2p)},
\]
for $t \geq 0$.
(ii) For all $p$ with $1 < p \leq \infty$, there is a constant $C_{p,d}$ such that
\[
\|w(t) - r(t)\|_{L^p(R)} \leq C_{p,d}t^{-(p-1)/2p},
\]
for $t > 0$.

For the proof of Lemma 2, see [4].

3 Approximation of $w(t, x)$ by a solution of one dimensional scalar viscous conservation law

In this section we study the convergence rate between the smooth rarefaction wave $w(t, x)$ and a solution $U(t, x)$ of the scalar viscous conservation law in one space dimension:
\[
\begin{cases}
U_t + (f(U))_x = U_{xx}, \\
U(0, x) = U_0(x).
\end{cases}
\]

Our aim in this section is to obtain a detailed asymptotic behavior of $U(t, x)$ in large time to be able to get the convergence rate toward $r(t, x)$ of solutions to the two dimensional scalar viscous conservation laws.

To do this, we decompose the solution as
\[
U(t, x) = w(t, x) + v(t, x).
\]
Then, the problem (21) is reduced to
\begin{align*}
\frac{\partial v}{\partial t} + (a(w)v)_x + (F(w, v)v^2)_x &= v_{xx} + w_{xx}, \\
v(0, x) &= v_0(x) \equiv U_0(x) - w_0(x),
\end{align*}
where
\[ F(w, v) = \frac{f(w + v) - f(w) - f'(w)v}{v^2}. \]
Note that $F$ is a smooth and bounded function of $(w, v)$. Thus, the problem we consider from now on becomes (22)-(23). We begin by showing the local existence result.

**Lemma 3 (local existence)** Suppose that $v_0 \in H^2(R) \cap L^1(R)$. Then there is a positive constant $T_0$ depending on $||v_0||_{H^2(R) \cap L^1(R)}$ and $d = u_+ - u_-$ such that the problem (22)-(23) has a unique solution $v(t, x)$ satisfying
\begin{align*}
v &\in C^0([0, T_0); H^2(R)) \cap C^1([0, T_0); L^2(R)) \\
&\quad \cap L^2([0, T_0); H^3(R)) \cap C^0([0, T_0); L^1(R)),
\end{align*}
where
\[ \Psi(v)(t) = G(t) * v_0 - \int_0^t \partial_x G(t - s) * [a(w)v + F(w, v)v^2 - w_x](s) ds, \]
where $G(t, x)$ is the Gauss kernel in one space dimension and $*$ denotes the convolution with respect to the space variable. It then follows from the expression of $\Psi$ and the lack of the space-integrability of $\partial_x v_0(x)$ that $v_x(t, \cdot)$ has the order $O(t^{-1/2})$ in $L^1(R)$, which shows the claim mentioned above.

Next, we state a priori estimate of $v$.

**Lemma 4 (a priori estimate)** Suppose that $v(t, x)$ is a solution of (22)-(23) belonging to the class as in (25) with $0 \leq t \leq T$.

(i) There holds
\[ ||v(t)||_{L^1(R)} \leq ||v_0||_{L^1(R)} + C_d \log(1 + t), \]
where $C_d$ is a constant depending on the size of $d = u_+ - u_-$. 

(ii) There exists a constant $\delta_1 > 0$ such that if
\[ N(T) \equiv \sup_{0 \leq t \leq T} ||v(t)||_{H^2(R)} \leq \delta_1, \]
and $d \leq \delta_1$, then $v(t, x)$ satisfies
\[
||\partial^k_x v(t)||^2_{L^2(R)} + \int_0^t \int_R w_x ||\partial^k_x v(s, x)||^2 ds dx
+ \int_0^t ||\partial^k_x v_x(s)||^2_{L^2(R)} ds \leq C(||v_0||^2_{H^k(R)} + \omega(d)), \quad \text{for } 0 \leq t \leq T, \quad k = 0, 1, 2, \tag{28}
\]
where $\omega(d)$ is a constant satisfying
\[
\omega(d) \to 0 \quad \text{as } d \to 0.
\]

**Proof of Lemma 4.** The $H^2(R)$-bound of $v$ (28) can be proved by a method as in [11] as well as by applying Lemma 1, so we omit the proof of it.

Here we only show (26). Let $j_\delta(\lambda)$ be the usual smoothing kernel in $R^1$, i.e.,
\[
j_\delta(\lambda) = \delta^{-1}j(\lambda/\delta), \tag{29}
\]
where $j$ is a smooth function which has a compact support and satisfies $\int_R j(\lambda) d\lambda = 1$. Let $\phi_\delta$ be the convolution of the sign function and $j_\delta$, i.e.,
\[
\phi_\delta(\lambda) = (j_\delta \ast \text{sign})(\lambda), \tag{30}
\]
and put
\[
\Phi_\delta(\lambda) = \int_0^\lambda \phi_\delta(\xi) d\xi. \tag{31}
\]
Note
\[
\phi_\delta'(\lambda) = 2j_\delta(\lambda). \tag{32}
\]

Multiplying (22) by $\phi_\delta(v)$ and integrating it with respect to $t$ and $x$, we have
\[
\int_R \Phi_\delta(v) dx + \int_0^t \int_R \phi_\delta(v) \{a(w)v + F(w, v)v^2\}_x ds dx ds
= \int_R \Phi_\delta(v_0) dx + \int_0^t \int_R \phi_\delta(v) v_x dx ds + \int_0^t \int_R \phi_\delta(v) w_x dx ds. \tag{33}
\]

**Claim:** There holds
\[
H(t) = \int_0^t \int_R \phi_\delta(v) \{a(w)v + F(w, v)v^2\}_x dx ds \to 0 \quad \text{as } \delta \to 0 \text{ for each } t. \tag{34}
\]

To see this, integrating by parts and making use of (32), we get
\[
H(t) = - \int_0^t \int_R 2j_\delta(v_x) \{a(w)v + F(w, v)v^2\} x dx ds
= -2 \int_0^t \int_R \left( \int_0^v j_\delta(\xi) d\xi \right) x \{a(w) + F(w, v)v\} dx ds
= 2 \int_0^t \int_0^\infty j_\delta(\xi) d\xi \{a'(w)w_x + F_w w_x v + F_v v_x v + F_v x\} dx ds. \tag{35}
\]
It follows from the definition of $j_{\delta}$
\[
\left| \int_{0}^{v} j_{\delta}(\xi) \xi d\xi \right| \leq \delta \int_{0}^{\infty} j(\xi) \xi d\xi \leq C\delta.
\]
Then, the integrand in (35) does not exceed
\[
C\delta(w_{x} + w_{x}|v| + |v_{x}| + |v_x|),
\]
which tends to 0 as $\delta \to 0$ for almost every $(t, x)$, and integrable on $(0, t) \times \mathbb{R}$ for each $t$ in view of Remark under Lemma 3. Thus the Lebesgue dominated convergence theorem implies (34). The claim has been verified.

Noting
\[
\int_{0}^{t} \int_{\mathbb{R}} \phi_{\delta}(v)v_{xx} dx ds = - \int_{0}^{t} \int_{\mathbb{R}} \phi_{\delta}'(v)v_{x}^{2} dx ds \leq 0,
\]
and making use of Lemma 1, we let $\delta \to 0$ in (33) and obtain the desired estimate (26). The proof of Lemma 4 is complete.

Combining Lemmas 3 and 4, we obtain the global existence result.

**Theorem 2 (global existence)** Suppose that $v_{0} \in H^{2}(\mathbb{R}) \cap L^{1}(\mathbb{R})$. Then, if
\[
\max(||v_{0}||_{H^{2}(\mathbb{R})}, d) < \delta,
\]
then the problem (22)-(23) has a unique global solution $v(t, x)$ satisfying
\[
v \in C^{0}([0, \infty); H^{2}(\mathbb{R}) \cap C^{1}([0, \infty); L^{2}(\mathbb{R})) \cap C^{0}([0, \infty); L^{1}(\mathbb{R}))
\]
\[
v_{x} \in L^{2}([0, \infty); H^{2}(\mathbb{R})),
\]
\[
t^{1/2}v_{x} \in C^{0}([0, \infty); L^{1}(\mathbb{R}))
\]
and the estimates (26) and (28) hold for any $T > 0$.

Our main result in this section is the following decay estimate of $v(t, x)$, which states the convergence rate of $U(t, \cdot)$ toward the smooth rarefaction wave $w(t, \cdot)$ in $H^{2}(\mathbb{R})$.

**Theorem 3 (decay estimate)** Let $v(t, x)$ be the solution of (22)-(23) obtained in Theorem 2. Then, for any $\epsilon > 0$ there exists a constant $C > 0$ such that the following decay estimates hold for $v(t, x)$:

\[
(1 + t)^{k+1/2+\epsilon}||\partial_{x}^{k}v(t)||^{2}_{L^{2}(\mathbb{R})} + \int_{0}^{t} (1 + s)^{k+1/2+\epsilon} \int_{\mathbb{R}} w_{x}(s, x) |\partial_{x}^{k}v(s, x)|^{2} dx ds
\]
\[
+ \int_{0}^{t} (1 + s)^{k+1/2+\epsilon} ||\partial_{x}^{k}v_{x}(s)||_{L^{2}(\mathbb{R})}^{2} ds \leq C\gamma_{k}^{2}(1 + t)^{\epsilon} \rho_{k}(t)^{2},
\]
where $\gamma_{k} = ||v_{0}||_{(H^{k} \cap L^{1})(\mathbb{R})} + C_{d}$ and $\rho_{k}(t) = \log^{2k+1-1}(2 + t)$ for $k = 0, 1, 2$. 
Corollary 1 (Convergence of $U - r$) There exists a constant $C > 0$ depending on $U_0$ and $d$ such that
\[ \|U(t) - r(t)\|_{L^2(R)} \leq Ct^{-1/4}\log(2+t) \quad \text{for } t > 0. \] (37)

Proof of Corollary 1. Combining Theorem 3 with Lemma 2, we can obtain (37). The proof of Corollary 1 is complete.

Proof of Theorem 3. We first show (36) with $k = 0$. When $N(t)$ is small ($N(t)$ is defined by (27)), by applying the $L^2$-energy method as in [11], we find
\[
\frac{1}{2} \frac{d}{dt} \|v(t)\|_{L^2(R)}^2 + \frac{1}{2} \alpha \int_R w_x v^2 dx + \|v_x(t)\|_{L^2(R)}^2 \leq C \left| \int_R vw_{xx} dx \right|. \] (38)

By Lemma 1, the right hand side of (38) is estimated as follows:
\[
\left| \int_R vw_{xx} dx \right| \leq C \|v(t)\|_{L^\infty(R)} \|w_{xx}(t)\|_{L^1(R)} \leq C_d \|v(t)\|_{L^\infty(R)} (1+t)^{-1}. \] (39)

Multiplying (38) by $(1+t)^{1/2+\epsilon}$ and taking into account of (39), we have
\[
\frac{1}{2} \frac{d}{dt} (1+t)^{1/2+\epsilon} \|v(t)\|_{L^2(R)}^2 + \frac{1}{2} \alpha (1+t)^{1/2+\epsilon} \int_R w_x v^2 dx \\
+ (1+t)^{1/2+\epsilon} \|v_x(t)\|_{L^2(R)}^2 \\
\leq C (1+t)^{-1/2+\epsilon} \|v(t)\|_{L^2(R)}^2 + C_d (1+t)^{-1/2+\epsilon} \|v(t)\|_{L^\infty(R)} \\
\leq C \|v(t)\|_{L^1(R)} + C_d (1+t)^{-1/2+\epsilon} \|v(t)\|_{L^\infty(R)}. \] (40)

Making use of (26) and Sobolev inequality
\[
\|f\|_{L^\infty(R)} \leq C \|f\|_{L^1(R)}^{1/3} \|f_x\|_{L^2(R)}^{2/3}, \] (41)
we compute in (40):
\[
\frac{1}{2} \frac{d}{dt} (1+t)^{1/2+\epsilon} \|v(t)\|_{L^2(R)}^2 + \frac{1}{2} \alpha (1+t)^{1/2+\epsilon} \int_R w_x v^2 dx \\
+ (1+t)^{1/2+\epsilon} \|v_x(t)\|_{L^2(R)}^2 \\
\leq C (\|v(t)\|_{L^1(R)} + C_d (1+t)^{-1/2+\epsilon} \|v(t)\|_{L^1(R)}^{1/3} \|v_x(t)\|_{L^2(R)}^{2/3} \\
\leq C (\|v_0\|_{L^1(R)} + C_d (1+t)^{-1/2+\epsilon} \log^{4/3} (2+t) \|v_x(t)\|_{L^2(R)}^{2/3} \\
\leq C (\|v_0\|_{L^1(R)} + C_d (1+t)^{-1+\epsilon} \log^2 (2+t) \\
+ \frac{1}{2} (1+t)^{1/2+\epsilon} \|v_x(t)\|_{L^2(R)}^2),
\]
that is,

$$
\frac{d}{dt}(1 + t)^{1/2+\epsilon} \|v(t)\|_{L^2(R)}^2 + \alpha(1 + t)^{1/2+\epsilon} \int_R w_x v^2 \, dx \\
+ (1 + t)^{1/2+\epsilon} \|v_x(t)\|_{L^2(R)}^2 \leq C(\|v_0\|_{L^1(R)} + C_d)^2 (1 + t)^{-1+\epsilon} \log^2 (2 + t).
$$

Integrating (42) with respect to time from 0 to $t$, we get (36) with $k = 0$. In particular, we obtain

$$
\|v(t)\|_{L^2(R)} \leq C\gamma_0 (1 + t)^{-1/4} \log (2 + t).
$$

Next we derive (36) with $k = 1$. Making the $L^2$-energy equality on $v_x$ and multiplying it by $(1 + t)^{3/2+\epsilon}$, we have

$$
\frac{1}{2} \frac{d}{dt} (1 + t)^{3/2+\epsilon} \|v_x(t)\|_{L^2(R)}^2 + (1 + t)^{3/2+\epsilon} \int_R v_x (a(w)v)_{xx} \, dx \\
+ (1 + t)^{3/2+\epsilon} \int_R v_x (F(w, v) v^2)_{xx} \, dx + (1 + t)^{3/2+\epsilon} \|v_{xx}(t)\|_{L^2(R)}^2 \\
= \frac{1}{2} (3/2 + \epsilon) (1 + t)^{3/2+\epsilon} \|v_x(t)\|_{L^2(R)}^2 + (1 + t)^{3/2+\epsilon} \int_R v_x w_{xx} \, dx.
$$

We first study the second term in the left hand side of (44). Integration by parts gives

$$(1 + t)^{3/2+\epsilon} \int_R v_x (a(w)v)_{xx} \, dx \\
= -(1 + t)^{3/2+\epsilon} \int_R v_{xx} a'(w) w_x v \, dx - (1 + t)^{3/2+\epsilon} \int_R a(w) v_x v_{xx} \, dx.
$$

The first term of the right hand side in (45) is estimated by making use of Lemma 1 as follows:

$$
\left| (1 + t)^{3/2+\epsilon} \int_R v_{xx} a'(w) w_x v \, dx \right| \\
\leq \frac{1}{8} (1 + t)^{3/2+\epsilon} \|v_{xx}(t)\|_{L^2(R)}^2 + C(1 + t)^{3/2+\epsilon} \|w_x(t)\|_{L^\infty(R)} \int_R w_x v^2 \, dx \\
\leq \frac{1}{8} (1 + t)^{3/2+\epsilon} \|v_{xx}(t)\|_{L^2(R)}^2 + C_d (1 + t)^{1/2+\epsilon} \int_R w_x v^2 \, dx.
$$

On the other hand, integrating by parts, we find that the second term of the right hand side in (45) is estimated as

$$
-(1 + t)^{3/2+\epsilon} \int_R a(w) v_x v_{xx} \, dx \\
= \frac{1}{2} (1 + t)^{3/2+\epsilon} \int_R d'(w) w_x v_x^2 \, dx \geq \frac{1}{2} \alpha (1 + t)^{3/2+\epsilon} \int_R w_x v_x^2 \, dx.
$$
Secondly we estimate the third term in the left hand side of (44). Integration by parts gives

\[ \left| (1 + t)^{3/2 + \epsilon} \int_{R} v_{x}(F(w, v)v^{2})_{xx}dx \right| \]

\[ = \left| -(1 + t)^{3/2 + \epsilon} \int_{R} v_{xx}(F(w, v)v^{2})_{x}dx \right| \]

\[ = \left| -(1 + t)^{3/2 + \epsilon} \int_{R} v_{xx}(F_{w}w_{x}v^{2} + F_{v}v_{x}v^{2} + F \cdot 2vv_{x})dx \right| \]

\[ \leq \frac{1}{8} (1 + t)^{3/2 + \epsilon} \int_{R} w_{x}^{2}v^{4}dx + C(1 + t)^{3/2 + \epsilon} \int_{R} v^{4}v_{x}^{2}dx + C(1 + t)^{3/2 + \epsilon} \int_{R} v^{2}v_{x}^{2}dx \]

\[ \equiv \frac{1}{8} (1 + t)^{3/2 + \epsilon} \int_{R} (v_{x}(t))^{2}L^{2}(R) + J_{1} + J_{2} + J_{3}. \]  

(48) We estimate the last three terms. For \( J_{1} \), from the Lemma 1, (41), and (28), we compute:

\[ J_{1} \leq C(1 + t)^{3/2 + \epsilon} \| v_{x}(t) \|_{L^{\infty}(R)} \| v(t) \|_{L^{\infty}(R)}^{2} \int_{R} w_{x}v^{2}dx \]

\[ \leq C(1 + t)^{1/2 + \epsilon} \int_{R} w_{x}v^{2}dx. \]  

(49) For \( J_{2} \), as in \( J_{1} \), we estimate:

\[ J_{2} \leq C(1 + t)^{3/2 + \epsilon} \sup_{0 \leq s \leq t} \| v(s) \|_{L^{\infty}(R)}^{2} \| v(t) \|_{L^{2}(R)}^{2} \| v_{x}(t) \|_{L^{\infty}(R)}^{2} \]

\[ \leq C(1 + t)^{3/2 + \epsilon} \| v(t) \|_{L^{2}(R)}^{2} \cdot \| v_{x}(t) \|_{L^{2}(R)} \| v_{xx}(t) \|_{L^{2}(R)}. \]  

(50) Furthermore, we apply (43) in the right hand side of (50). Then, \( J_{2} \)

\[ \leq C(1 + t)^{3/2 + \epsilon} \log^{2}(2 + t) \cdot \| v_{x}(t) \|_{L^{2}(R)} \| v_{xx}(t) \|_{L^{2}(R)} \]

\[ = C(1 + t)^{1+\epsilon} (3/2+\epsilon/2) \log^{2}(2 + t) \| v_{x}(t) \|_{L^{2}(R)} \cdot (1 + t)^{(3/2+\epsilon)/2} \| v_{xx}(t) \|_{L^{2}(R)} \]

\[ \leq C(1 + t)^{1/2 + \epsilon} \log^{4}(2 + t) \| v_{x}(t) \|_{L^{2}(R)}^{4} + \frac{1}{8} (1 + t)^{3/2 + \epsilon} \| v_{xx}(t) \|_{L^{2}(R)}^{2}. \]  

(51) \( J_{3} \) is majorized by the same bound as that of \( J_{2} \), that is,

\[ J_{3} \leq C(1 + t)^{1/2 + \epsilon} \log^{4}(2 + t) \| v_{x}(t) \|_{L^{2}(R)}^{4} + \frac{1}{8} (1 + t)^{3/2 + \epsilon} \| v_{xx}(t) \|_{L^{2}(R)}^{2}. \]  

(52) Collecting (49), (51) and (52), we arrive at the estimate

\[ \left| (1 + t)^{3/2 + \epsilon} \int_{R} v_{x}(F(w, v)v^{2})_{xx}dx \right| \]

\[ \leq C(1 + t)^{1/2 + \epsilon} \int_{R} w_{x}v^{2}dx + C(1 + t)^{1/2 + \epsilon} \log^{4}(2 + t) \| v_{x}(t) \|_{L^{2}(R)}^{2} \]

\[ + \frac{1}{4} (1 + t)^{3/2 + \epsilon} \| v_{xx}(t) \|_{L^{2}(R)}^{2}. \]  

(53)
Finally, we estimate the second term of the right hand side in (44). Integration by parts and Lemma 1 gives

\[
\left| (1 + t)^{3/2+\varepsilon} \int_R v_x w_{xxx} \, dx \right| \\
= \left| -(1 + t)^{3/2+\varepsilon} \int_R v_{xx} w_{xx} \, dx \right| \\
\leq \frac{1}{8} (1 + t)^{3/2+\varepsilon} \| v_{xx}(t) \|^2_{L^2(R)} + C(1 + t)^{3/2+\varepsilon} \| w_{xx}(t) \|^2_{L^2(R)} \\
\leq \frac{1}{8} (1 + t)^{3/2+\varepsilon} \| v_{xx}(t) \|^2_{L^2(R)} + C_d (1 + t)^{-1+\varepsilon}.
\]

Collecting all the estimates (46)-(54), we arrive at

\[
\frac{1}{2} \frac{d}{dt} (1 + t)^{3/2+\varepsilon} \| v_x(t) \|^2_{L^2(R)} + \frac{\alpha}{2} (1 + t)^{3/2+\varepsilon} \int_R w_x v_x^2 \, dx \\
+ \frac{3}{8} (1 + t)^{3/2+\varepsilon} \| v_{xx}(t) \|^2_{L^2(R)} \leq \frac{1}{2} (3/2 + \varepsilon) (1 + t)^{1/2+\varepsilon} \| v_x(t) \|^2_{L^2(R)} \\
+ C(1 + t)^{1/2+\varepsilon} \int_R w_x v^2 \, dx + C(1 + t)^{1/2+\varepsilon} \log^4 (2 + t) \| v_x(t) \|^2_{L^2(R)} + C_d (1 + t)^{-1+\varepsilon}.
\]

Integrating (55) with respect to \( t \) and making use of (36), we get

\[
(1 + t)^{3/2+\varepsilon} \| v_x(t) \|^2_{L^2(R)} + \int_0^t (1 + s)^{3/2+\varepsilon} \int_R w_x v_x^2 \, dx \, ds + \int_0^t (1 + s)^{3/2+\varepsilon} \| v_{xx}(t) \|^2_{L^2(R)} \leq \| v_{0x} \|^2_{L^2(R)} + \gamma_0^2 (1 + t)^4 \log^2 (2 + t) + \gamma_0^2 (1 + t)^4 \log^6 (2 + t) + C_d (1 + t)^{\varepsilon}.
\]

Clearing up the above inequality, we arrive at (36) with \( k = 1 \). In particular, we have

\[
\| v_x(t) \|_{L^2(R)} \leq \gamma_1 (1 + t)^{-3/4} \log^3 (2 + t).
\]

Finally we show (36) with \( k = 2 \). Similar procedure to derive (44) also gives

\[
\frac{1}{2} \frac{d}{dt} (1 + t)^{5/2+\varepsilon} \| v_{xx}(t) \|^2_{L^2(R)} + (1 + t)^{5/2+\varepsilon} \int_R v_{xx}(a(w)v)_{xxx} \, dx \\
+ (1 + t)^{5/2+\varepsilon} \int_R v_{xx}(F(w, v)v^2)_{xxx} \, dx + (1 + t)^{5/2+\varepsilon} \| v_{xxx}(t) \|^2_{L^2(R)} \leq \frac{1}{2} (5/2 + \varepsilon) (1 + t)^{3/2+\varepsilon} \| v_{xx}(t) \|^2_{L^2(R)} + (1 + t)^{5/2+\varepsilon} \int_R v_{xx} w_{xxx} \, dx.
\]

First we study the second term of the left hand side in (56). Integration by parts gives

\[
(1 + t)^{5/2+\varepsilon} \int_R v_{xx}(a(w)v)_{xxx} \, dx
\]
\[ \equiv (1 + t)^{5/2 + \epsilon} \int_R v_{xxx}a''(w)w^2 v dx 
\] 
\[ + (1 + t)^{5/2 + \epsilon} \int_R v_{xxx}a'(w)w_{xx} v dx 
\] 
\[ - 2(1 + t)^{5/2 + \epsilon} \int_R v_{xxx}a'(w)w_x^2 v dx 
\] 
\[ - (1 + t)^{5/2 + \epsilon} \int_R v_{xxx}a(w)v_{xx} v dx 
\] 
\[ \equiv K_1 + K_2 + K_3 + K_4. \quad (57) \]

\( K_1 \) is estimated by using Lemma 1 as follows:

\[ |K_1| \leq C(1 + t)^{5/2 + \epsilon} \left( \|v_{xxx}(t)\|_{L^2(R)} \right) \left( \|v_x^2(t)\|_{L^2(R)} \right) \]

\[ \leq \frac{1}{16}(1 + t)^{5/2 + \epsilon} \left( \|v_{xxx}(t)\|_{L^2(R)} \right)^2 + C(1 + t)^{5/2 + \epsilon} \left( \|w_x(t)\|_{L^\infty(R)} \right) \int_R w_x^2 v^2 dx 
\]

\[ \leq \frac{1}{16}(1 + t)^{5/2 + \epsilon} \left( \|v_{xxx}(t)\|_{L^2(R)} \right)^2 + C(1 + t)^{-1/2 + \epsilon} \int_R w_x^2 v^2 dx. \quad (58) \]

Computing \( K_2 \) as above,

\[ |K_2| \leq \frac{1}{16}(1 + t)^{5/2 + \epsilon} \left( \|v_{xxx}(t)\|_{L^2(R)} \right)^2 + C(1 + t)^{5/2 + \epsilon} \int_R w_x^2 v^2 dx. \]

Applying (43) as well as (41) and Lemma 1 to the last integral,

\[ \int_R w_x^2 v^2 dx \leq \left( \|w_x(t)\|_{L^2(R)} \right)^2 \left( \|v(t)\|_{L^2(R)} \right) \left( \|v_x(t)\|_{L^2(R)} \right) 
\]

\[ \leq \gamma_1^2 (1 + t)^{-5/2} \cdot (1 + t)^{-1/2} \log(2 + t) \cdot (1 + t)^{-3/4} \log^3(2 + t) 
\]

\[ \leq \gamma_1^2 (1 + t)^{-7/2} \log^4(2 + t). \]

Hence,

\[ |K_2| \leq \frac{1}{16}(1 + t)^{5/2 + \epsilon} \left( \|v_{xxx}(t)\|_{L^2(R)} \right)^2 + \gamma_1^2 (1 + t)^{-7/2} \log^4(2 + t). \quad (59) \]

For \( K_3 \), in a similar way,

\[ |K_3| \leq \frac{1}{16}(1 + t)^{5/2 + \epsilon} \left( \|v_{xxx}(t)\|_{L^2(R)} \right)^2 + C(1 + t)^{5/2 + \epsilon} \left( \|w_x(t)\|_{L^\infty(R)} \right) \left( \|v_x(t)\|_{L^2(R)} \right)^2 
\]

\[ \leq \frac{1}{16}(1 + t)^{5/2 + \epsilon} \left( \|v_{xxx}(t)\|_{L^2(R)} \right)^2 + C(1 + t)^{1/2 + \epsilon} \left( \|v_x(t)\|_{L^2(R)} \right)^2. \quad (60) \]

For \( K_4 \), we integrate by parts to get

\[ K_4 \geq \frac{1}{2} \alpha(1 + t)^{5/2 + \epsilon} \int_R w_x^2 v^2 dx. \quad (61) \]
Next we estimate the third term of the left hand side in (56). Estimating it as before, we obtain

\[
(1 + t)^{5/2 + \epsilon} \int_{R} v_{xx}(F(w, v)v^2)_{xxx} \, dx \nln(1 + t)^{5/2 + \epsilon} \int_{R} v_{xx}(F(w, v)v^2)_{xxx} \, dx \leq (1 + t)^{5/2 + \epsilon} \|v_{xxx}(t)\|_{L^2(R)} \|f(w, v)v^2\|_{L^2(R)} + C(1 + t)^{3/2 + \epsilon} \|v_{xx}w_{x}(t)\|_{L^2(R)}^2 + C(1 + t)^{3/2 + \epsilon} \|v_{xx}w_{x}(t)\|_{L^2(R)}^2 + C(1 + t)^{3/2 + \epsilon} \|v_{xx}w_{x}(t)\|_{L^2(R)}^2 + C(1 + t)^{3/2 + \epsilon} \|v_{xx}w_{x}(t)\|_{L^2(R)}^2 \\
\equiv \frac{1}{16}(1 + t)^{5/2 + \epsilon} \|v_{xxx}(t)\|_{L^2(R)}^2 + K_5 + K_6 + K_7 + K_8 + K_9 + K_10 + K_{11} + K_{12}.
\]

The estimates of \(K_5-K_{12}\) can be done in a similar way for the estimates of the second term of the left hand side in (56). We present only the results of them:

\[
\begin{align*}
K_5 &\leq C\|v_{x}(t)\|_{L^2(R)}^2, \\
K_6 &\leq C(1 + t)^{3/2 + \epsilon} \|v_{x}(t)\|_{L^2(R)}^2, \\
K_7 &\leq C\|v_{x}(t)\|_{L^2(R)}^2, \\
K_8 &\leq C\|v_{x}(t)\|_{L^2(R)}^2 + C(1 + t)^{3/2 + \epsilon} \|v_{xx}(t)\|_{L^2(R)}^2, \\
K_9 &\leq C(1 + t)^{3/2 + \epsilon} \|v_{xx}(t)\|_{L^2(R)}^2, \\
K_{10} &\leq \gamma_1^2(1 + t)^{-1+\epsilon} + C(1 + t)^{3/2 + \epsilon} \|v_{xx}(t)\|_{L^2(R)}^2, \\
K_{11} &\leq C(1 + t)^{3/2 + \epsilon} \|v_{x}(t)\|_{L^2(R)}^2 \log^{12}(2 + t) + C(1 + t)^{3/2 + \epsilon} \|v_{xx}(t)\|_{L^2(R)}^2 \log^{4}(2 + t), \\
K_{12} &\leq \frac{1}{16}(1 + t)^{5/2 + \epsilon} \|v_{xxx}(t)\|_{L^2(R)}^2 + C(1 + t)^{3/2 + \epsilon} \|v_{xx}(t)\|_{L^2(R)}^2 \log^{4}(2 + t).
\end{align*}
\]

Collecting (62)-(69), we get the estimate

\[
(1 + t)^{5/2 + \epsilon} \int_{R} v_{xx}(F(w, v)v^2)_{xxx} \, dx \leq C(1 + t)^{3/2 + \epsilon} \|v_{x}(t)\|_{L^2(R)} + C(1 + t)^{3/2 + \epsilon} \|v_{xx}(t)\|_{L^2(R)} + \gamma_1^2(1 + t)^{-1+\epsilon} \\
+ C(1 + t)^{3/2 + \epsilon} \|v_{x}(t)\|_{L^2(R)}^2 \log^{12}(2 + t) + C(1 + t)^{3/2 + \epsilon} \|v_{xx}(t)\|_{L^2(R)}^2 \log^{4}(2 + t) \\
+ C(1 + t)^{3/2 + \epsilon} \|v_{xx}(t)\|_{L^2(R)}^2 \log^{4}(2 + t)
\]

Finally we estimate the second term of the right hand side of (56). Making use of Lemma 2,

\[
(1 + t)^{5/2 + \epsilon} \int_{R} v_{xx}w_{xxxx} \, dx
\]
\[
\leq \frac{1}{16}(1 + t)^{5/2 + \epsilon} \|v_{xxx}(t)\|_{L^2(\mathbb{R})}^2 + C(1 + t)^{5/2 + \epsilon} \|w_{xxx}(t)\|_{L^2(\mathbb{R})}^2 \\
\leq \frac{1}{16}(1 + t)^{5/2 + \epsilon} \|v_{xxx}(t)\|_{L^2(\mathbb{R})}^2 + C_d(1 + t)^{-1 + \epsilon} 
\] (71)

Now, collecting the estimates (58)-(61), (70), and (71), we obtain
\[
\frac{1}{2} \frac{d}{dt} (1 + t)^{s/2 + \epsilon} \|u_{xx}(t)\|_{L^2(\mathbb{R})}^2 + C \|u_{xxx}(t)\|_{L^2(\mathbb{R})}^2 + C(1 + t)^{5/2 + \epsilon} \|v_{xxx}(t)\|_{L^2(\mathbb{R})}^2 \\
\leq C(1 + t)^{-1/2 + \epsilon} \int_{\mathbb{R}} w_x u_x^2 dX + C_d(1 + t)^{-1 + \epsilon} + C(1 + t)^{1/2 + \epsilon} \|v_x(t)\|_{L^2(\mathbb{R})}^2 \log^4(2 + t).
\]

Integrating the above inequality with respect to time from 0 to \(t\) with the aid of (36) with \(k = 0, 1\), we arrive at (36) with \(k = 2\). The proof of Theorem 3 is complete.

We close this section to state one more property of \(U(t, x)\) obtained by Xin [11], which plays an essential role in the next section.

**Lemma 5 ([11])** Suppose that \(U_0(x)\) has the following properties:
\[
\frac{d}{dx} U_0(x) > 0 \quad \text{and} \quad \frac{d^2}{dx^2} U_0(x) \leq C \frac{d}{dx} U_0(x),
\] (72)
for any \(x \in \mathbb{R}\). Then, \(U(t, x)\) satisfies
\[
U_x(t, x) > 0 \quad \text{for } t > 0 \text{ and } x \in \mathbb{R},
\] (73)
and
\[
|U_{xx}(t, x)| \leq C U_x(t, x) \quad \text{for } t > 0 \text{ and } x \in \mathbb{R}.
\] (74)

Note that if we choose \(v_0(x) \equiv 0\) in (23), then \(U_0(x)\) in (21) satisfies (72).

### 4 convergence of \(u(t, x, y) - r(t, x)\)

In this section we prove Theorem 1. Throughout this section, \(v_0(x) \equiv 0\) is adopted.

First, as in the work of Xin [11], we decompose the solution \(u(t, x, y)\) as follows:
\[
u(t, x, y) = U(t, x) + V(t, x, y),
\]
where $U(t, x)$ is the solution of (21) obtained in Section 3. Then, the problem (1)-(2) is reduced to
\[
\begin{align*}
V_t + (f'(U)V)_x + (F(U, V)V^2)_x + (g'(U)V)_y + (G(U, V)V^2)_y &= V_{xx} + V_{yy}, \\
V(0, x, y) &= V_0(x, y) \equiv u_0(x) - w_0(x),
\end{align*}
\] (75)
where $F(U, V)$ is in (24) and $G(U, V)$ is defined by
\[
G(U, V) = \frac{g(U + V) - g(U) - g'(U)V}{V^2}.
\] (76)

From now on, we study the problem (75).

Our main purpose in this section is to derive the decay estimate for $V$.

Throughout this section, we use the notation $\partial^k$ as in the meaning
\[
\partial^k = \sum_{i+j=k} \partial_i \partial_j xy.
\] (77)

For the problem (75), Xin [11] showed the following global existence result.

**Theorem 4 (global existence [11])** There exists a constant $\delta_2$ such that if $\|V_0\|_{H^3(R^2)} \leq \delta_2$, then the problem (75) has a unique global solution $V(t, x, y)$ satisfying
\[
\|V(t)\|_{H^3(R^2)} + \int_0^t \iint_{R^2} U_x V^2(s, x, y) dxdyds \\
+ \int_0^t \|\partial^{1} V(t)\|_{H^3(R^2)}^2 ds \leq C\|V_0\|_{H^3(R^2)}^2
\] (78)
for all $t \geq 0$.

Furthermore, when the integrability of $V_0$ is imposed, we have

**Lemma 6 (L^1-estimate)** Suppose further in Theorem 4 that $V_0 \in L^1(R^2)$. Then, the solution $V(t, x, y)$ also satisfies
\[
\|V(t)\|_{L^1(R^2)} \leq \|V_0\|_{L^1(R^2)}.
\] (79)

The proof of Lemma 6 can be done in a similar way as in having derived (26).

Our main result in this section is the following decay estimate of $V$.

**Theorem 5 (decay estimate)** Let $V(t, x, y)$ be the solution of (75) obtained in Theorem 4 and Lemma 6. Then, for any $\epsilon > 0$ there exists a constant $C > 0$ such that the following estimate holds:
\[
(1 + t)^{k+1+\epsilon} \|\partial^k V(t)\|_{L^2(R^2)}^2 + \int_0^t (1 + s)^{k+1+\epsilon} \iint_{R^2} U_x |\partial^k V|^2 dxdyds \\
+ \int_0^t (1 + s)^{k+1+\epsilon} \|\partial^{k+1} V(s)\|_{L^2(R^2)}^2 ds \leq Cd(1 + t)^{k+1+\epsilon} \theta_k(t)^2 \|V_0\|_{(H^k \cap L^1)(R^2)}^2,
\] (80)
where $\theta_0(t) = 1$, $\theta_1(t) = \log^5(2 + t)$ and $\theta_2(t) = \log^{33/2}(2 + t).
Here, admitting Theorem 5 temporarily, we prove Theorem 1.

**Proof of Theorem 1.** First we write

\[ \sup_{y \in R} \| u(t) - r(t) \|_{L^2(R_x)} \leq \sup_{y \in R} \| u(t) - U(t) \|_{L^2(R_x)} + \| U(t) - r(t) \|_{L^2(R)}. \]  

(81)

It follows from Corollary 1 that the second term of the right hand side in (81) does not exceed \( C t^{-1/4} \log(2 + t) \). So, we only estimate the first term of the right hand side of (81). Note that the following Sobolev inequality holds:

\[ \sup_{y \in R} \| f \|_{L^2(R_x)} \leq C \| f \|_{L^2(R)}^{1/2} \| f_y \|_{L^2(R)}^{1/2}, \]  

(82)

for functions \( f = f(x, y) \). Then, it follows from (80) and (82)

\[ \sup_{y \in R} \| V(t) \|_{L^2(R_x)} \leq C \| V(t) \|_{L^2(R)}^{1/2} \| V_y(t) \|_{L^2(R)}^{1/2} \leq C(1 + t)^{-3/4} \log^{5/4}(2 + t), \]

which means

\[ \sup_{y \in R} \| u(t) - U(t) \|_{L^2(R_x)} \leq C(1 + t)^{-3/4} \log^{5/4}(2 + t). \]  

(83)

From the above, we arrive at

\[ \sup_{y \in R} \| u(t) - r(t) \|_{L^2(R_x)} \leq C t^{-1/4} \log(2 + t), \]

which gives (15). The proof of Theorem 1 is complete.

It remains to prove Theorem 5.

**Proof of Theorem 5.** By using the \( L^2 \)-energy method as treated by Xin [11], we get

\[ \frac{d}{dt} \| V(t) \|_{L^2(R)}^2 + \iint_{R^2} U_x V^2 dxdy + C \| \partial^1 V(t) \|_{L^2(R)}^2 \leq 0. \]  

(84)

Multiplying (84) by \( (1 + t)^{1+\varepsilon} \), we have

\[ \frac{d}{dt} (1 + t)^{1+\varepsilon} \| V(t) \|_{L^2(R)}^2 + (1 + t)^{1+\varepsilon} \iint_{R^2} U_x V^2 dxdy + C(1 + t)^{1+\varepsilon} \| \partial^1 V(t) \|_{L^2(R)}^2 \leq C(1 + t)^{\varepsilon} \| V(t) \|_{L^2(R)}^2. \]  

(85)

With the aid of (79) and Sobolev inequality

\[ \| f \|_{L^\infty(R^2)} \leq C \| f \|_{L^2(R)}^{1/2} \| \partial^1 f \|_{L^2(R)}^{1/2} \]  

for \( f = f(x, y) \),

(86)
we continue the computations:

$$
\frac{d}{dt}(1+t)^{1+\varepsilon}\|V(t)\|_{L^{2}(R^{2})}^{2} + (1+t)^{1+\varepsilon} \int_{R^{2}} U_{x} V^{2} dx dy \\
+ C(1+t)^{1+\varepsilon}\|\partial^{1}V(t)\|_{L^{2}(R^{2})}^{2} \\
\leq C(1+t)^{1+\varepsilon}\|V_{0}\|_{L^{1}(R^{2})} \cdot (1+t)^{1+\varepsilon}/2\|\partial^{1}V(t)\|_{L^{2}(R^{2})} \\
\leq Cr^{-1}(1+t)^{1+\varepsilon}\|\partial^{1}V(t)\|_{L^{2}(R^{2})}^{2},
$$

(87)

where $r$ is a positive constant sufficiently small. Integrating (87) with respect to time from 0 to $t$, we obtain

$$(1+t)^{1+\varepsilon}\|V(t)\|_{L^{2}(R^{2})}^{2} + \int_{0}^{t} (1+s)^{1+\varepsilon} \int_{R^{2}} U_{x} V^{2} dx dy ds \\
+ C\int_{0}^{t} (1+s)^{1+\varepsilon}\|\partial^{1}V(s)\|_{L^{2}(R^{2})}^{2} ds \leq \|V_{0}\|_{L^{2}(R^{2})} + C(1+t)^{1+\varepsilon}\|V_{0}\|_{L^{1}(R^{2})}^{2},
$$

which gives (80) with $k = 0$.

Secondly, we derive (80) with $k = 1$. $L^{2}$-energy method gives

$$
\frac{d}{dt}\|\partial^{1}V(t)\|_{L^{2}(R^{2})}^{2} + \int_{R^{2}} U_{x} |\partial^{1}V|^{2} dx dy + C\|\partial^{2}V(t)\|_{L^{2}(R^{2})}^{2} \\
\leq C\|U_{x}(t)\|_{L^{\infty}(R)} \left( \int_{R^{2}} U_{x} V^{2} dx dy + \|\partial^{1}V(t)\|_{L^{2}(R^{2})}^{2} \right) \\
+ C\|V(t)\|_{L^{2}(R^{2})} \cdot \|\partial^{1}V(t)\|_{L^{2}(R^{2})}^{2}.
$$

(88)

Multipluing (88) by $(1+t)^{2+\varepsilon}$, we have

$$
\frac{d}{dt}(1+t)^{2+\varepsilon}\|\partial^{1}V(t)\|_{L^{2}(R^{2})}^{2} + (1+t)^{2+\varepsilon} \int_{R^{2}} U_{x} |\partial^{1}V|^{2} dx dy \\
+ C(1+t)^{2+\varepsilon}\|\partial^{2}V(t)\|_{L^{2}(R^{2})}^{2} \leq C(1+t)^{1+\varepsilon}\|\partial^{1}V(t)\|_{L^{2}(R^{2})}^{2} \\
+ C(1+t)^{2+\varepsilon}\|U_{x}(t)\|_{L^{\infty}(R)} \left( \int_{R^{2}} U_{x} V^{2} dx dy + \|\partial^{1}V(t)\|_{L^{2}(R^{2})}^{2} \right) \\
+ C(1+t)^{2+\varepsilon}\|V(t)\|_{L^{2}(R^{2})} \cdot \|\partial^{1}V(t)\|_{L^{2}(R^{2})}^{2}.
$$

(89)

It should be noted that the following estimate holds:

$$
\|U_{x}(t)\|_{L^{\infty}(R)} \leq C_{d}(1+t)^{-1}\log^{5}(2+t).
$$

(90)

This can be obtained by using Lemma 1 and Theorem 3 in the equality

$$
U_{x} = w_{x} + v_{x}.
$$
Then, integrating (89) with respect to time from 0 to \( t \) and making use of (90), we estimate:

\[
(1 + t)^{2+\varepsilon}||\partial^1 V(t)||^2_{L^2(R^2)} + \int_0^t (1 + s)^{2+\varepsilon} \int_{R^2} U_x |\partial^1 V(s)|^2 dx dy ds
+ \int_0^t (1 + s)^{2+\varepsilon} ||\partial^2 V(s)||^2_{L^2(R^2)} ds
\leq ||\partial^1 V_0||^2_{L^2(R^2)} + C \int_0^t (1 + s)^{1+\varepsilon} ||\partial^1 V(s)||^2_{L^2(R^2)} ds
+ C_d \int_0^t (1 + s)^{1+\varepsilon} \log(2 + s) \left( \int_{R^2} U_x V^2 dx dy + ||\partial^1 V(s)||^2_{L^2(R^2)} \right) ds
+ C \int_0^t (1 + s)^{2+\varepsilon} ||V(s)||^2_{L^2(R^2)} ||\partial^1 V(s)||^2_{L^2(R^2)} ds.
\]

Taking account of (80) with \( k = 0 \) and \( L^2 \)-bound of \( \partial^2 V(t) \), we get

\[
(1 + t)^{2+\varepsilon}||\partial^1 V(t)||^2_{L^2(R^2)} + \int_0^t (1 + s)^{2+\varepsilon} \int_{R^2} U_x |\partial^1 V(s)|^2 dx dy ds
+ \int_0^t (1 + s)^{2+\varepsilon} ||\partial^2 V(s)||^2_{L^2(R^2)} ds
\leq ||\partial^1 V_0||^2_{L^2(R^2)} + C(1 + t)^{\varepsilon} ||V_0||^2_{L^2(R^2)} + C_d (1 + t)^{\varepsilon} \log^5(2 + t) ||V_0||^2_{L^2(R^2)}.
\]

Clearing up the right hand side of the above inequality, we get (80) with \( k = 1 \).

Finally, we derive (80) with \( k = 2 \). \( L^2 \)-energy method gives

\[
\frac{d}{dt} ||\partial^2 V(t)||^2_{L^2(R^2)} + \int_{R^2} U_x |\partial^2 V|^2 dx dy dS + C ||\partial^2 V(t)||^2_{L^2(R^2)}
\leq C ||(f'(U) + g'(U))\partial^1 V||^2_{L^2(R^2)} + C ||(f'(U) + g'(U))\partial^1 V||^2_{L^2(R^2)}
+ C ||\partial^2 (F(U, V)V^2)||^2_{L^2(R^2)} + C ||\partial^2 (G(U, V)V^2)||^2_{L^2(R^2)}. \tag{91}
\]

Multiplying (91) by \( (1 + t)^{3+\varepsilon} \) and integrating it with respect to time from 0 to \( t \), we have

\[
(1 + t)^{3+\varepsilon}||\partial^2 V(t)||^2_{L^2(R^2)} + \int_0^t (1 + s)^{3+\varepsilon} \int_{R^2} U_x |\partial^2 V|^2 dx dy ds
+ C \int_0^t (1 + s)^{3+\varepsilon} ||\partial^3 V(s)||^2_{L^2(R^2)} ds
\leq ||\partial^2 V_0||^2_{L^2(R^2)} + C \int_0^t (1 + s)^{2+\varepsilon} ||\partial^2 V(s)||^2_{L^2(R^2)} ds
+ C \int_0^t (1 + s)^{3+\varepsilon} ((f'(U) + g'(U))\partial^1 V(s))||^2_{L^2(R^2)} ds
+ C \int_0^t (1 + s)^{3+\varepsilon} ((f'(U) + g'(U))\partial^1 V(s))||^2_{L^2(R^2)} ds
+ C \int_0^t (1 + s)^{3+\varepsilon} ||\partial^2 (F(U, V)V^2)||^2_{L^2(R^2)} + ||\partial^2 (G(U, V)V^2)||^2_{L^2(R^2)} ds. \tag{92}
\]
We here only estimate the term
\[ \Lambda(t) \equiv \int_0^t (1 + s)^3 \|U_{xx}(s)V(s)\|_{L^2(R^2)}^2 ds, \]
which arises from the first term of the right hand side in (92). The rest terms in (92) can be treated as before. Note
\[ \|U_{xx}(s)V(s)\|_{L^2(R^2)}^2 \leq \|U_{xx}(s)\|_{L^2(R_1)}^2 \sup_{x \in R} \int_R V(s, x, y)^2 dy. \]
It then follows from Lemma 1, (36) with \( k = 2 \), and (80) with \( k = 0, 1 \)
\[ \Lambda(t) \leq \int_0^t (1 + s)^{3+\varepsilon} \|U_{xx}(s)\|_{L^2(R_1)}^2 \|V(s)\|_{L^2(R^2)} \|V_x(s)\|_{L^2(R^2)} ds \]
\[ \leq C_d \|V_0\|_{(H^1 \cap L^1)(R^2)} \int_0^t (1 + s)^{-1+\varepsilon} \log^{3/2}(2 + s) ds \]
\[ \leq C_d \|V_0\|_{(H^1 \cap L^1)(R^2)} (1 + t)^{\varepsilon} \log^{3/2}(2 + t), \]
which gives the right hand side of (80) with \( k = 2 \). The proof of Theorem 5 is complete.

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References


