In this note we consider a natural mapping between the following two spaces:

\[ E = \{ (\text{certain}) \text{ linear ordinary differential equations on (marked)} \]
\[ \text{Riemann surfaces s.t. the local monodromy representations}\]
\[ \text{around the singular points are as specified} \}, \]

\[ R = \{ \text{representation classes of } \pi_{1}(\text{punctured surface}) \text{ s.t. the local representations around the punctures are as specified} \}. \]

Let us denote by \( F \) the mapping that assigns the elements of \( E \) their monodromy representations and consider for instance a one-parameter family of differential equations lying in the fiber \( F^{-1}(r) \) above a point \( r \in R \). The characteristic feature of that family is of course that the corresponding monodromy of each element of the family is always the same \( r \) and therefore we call it an isomonodromic family. Our goal in this note will be to give an infinitesimal description of isomonodromic families in terms of a completely integrable system of partial differential equations on some local coordinate parameters of \( E \). More specifically, that amounts to describing the tangential directions to the fibers \( F^{-1}(r) \) and can be carried out in the following geometric manner. The key observation of our method is that there exists a natural symplectic structure \( \omega \) on the space \( R \) of representations [2]. (A symplectic structure \( \omega \) on \( R \) is, by definition, a closed nondegenerate 2-form on \( R \).) By pulling back the 2-form \( \omega \) onto \( E \) by the mapping \( F \), we obtain a possibly degenerate closed 2-form on \( E \); and that 2-form can then be used to describe the tangential directions to the fibers of \( F \) as follows: For a tangent vector \( \xi \) to \( E \) at a point \( p \in E \), we have

\[ \xi \text{ is tangent to a fiber of } F \]
\[ \iff d_{p}F(\xi) = 0 \]
\[ \iff \omega(d_{p}F(\xi), \cdot) \equiv 0 \quad \text{since } \omega \text{ is nondegenerate} \]
\[ \iff F^{*}\omega(\xi, \cdot) \equiv 0 \quad \text{where we assume } d_{p}F \text{ is surjective}. \]
The surjectivity condition will always be satisfied in our discussion below.) It thus follows that the problem of describing the tangential directions to the fibers of $F$ can be reduced to determining precisely the vectors $\xi$ such that $F^*\omega(\xi, \cdot) \equiv 0$ (we say that $F^*\omega$ is degenerate in the direction $\xi$ if $F^*\omega(\xi, \cdot) \equiv 0$) and consequently we shall first write out the pulled-back 2-form $F^*\omega$ explicitly in terms of some local coordinates and then determine the directions making $F^*\omega$ degenerate. (The distribution $\{ \xi \in TE; F^*\omega(\xi, \cdot) \equiv 0 \}$ is obviously integrable (since the fibers of $F$ are precisely the maximal integral manifolds) and therefore defines a foliation on $E$, which will be called the null-foliation of $F^*\omega$.)

Let $X$ be a compact Riemann surface of genus one, and $H = \{ \tau \in \mathbb{C}; \text{Im}\tau > 0 \}$ the upper half-plane. Selecting a suitable $\tau \in H$, one can represent $X$ as the quotient $\mathbb{C}/\mathbb{Z} \cdot 1 + \mathbb{Z} \cdot \tau$; and then equations on $X$ can be represented as equations on $\mathbb{C}$ with doubly periodic coefficients. Consider the Fuchsian equation

$$
\frac{d^2y}{dz^2} = q(z) y,
$$

$$
q(z) = k + \sum_{i=0}^{m} \left[ H_i \zeta(z-t_i, \tau) + \frac{1}{4} (\theta_i^2 - 1) \wp(z-t_i, \tau) \right] + \sum_{\alpha=0}^{m} \left[ -\mu_\alpha \zeta(z-\lambda_\alpha, \tau) + \frac{3}{4} \wp(z-\lambda_\alpha, \tau) \right],
$$

$$
\sum_{i=0}^{m} H_i - \sum_{\alpha=0}^{m} \mu_\alpha = 0,
$$

where $\zeta(z, \tau)$ and $\wp(z, \tau)$ denote Weierstrass’ $\zeta$-function and $\wp$-function with fundamental periods 1, $\tau$ and $t_0$ will always be normalized so that $t_0 = 0$. It has its (regular) singularities at $[t_i]$ ($i = 0, \ldots, m$) and $[\lambda_\alpha]$ ($\alpha = 0, \ldots, m$) with characteristic exponents $\frac{1}{2}(1 \pm \theta_i)$ and $\frac{1}{2}(1 \pm 2)$ respectively ($[z]$ denotes the congruence class of a point $z \in \mathbb{C}$) and determines its monodromy representation

$$
\rho: \pi_1(X \setminus \{ [t_0], \ldots, [t_m], [\lambda_0], \ldots, [\lambda_m] \}) \to \text{SL}(2, \mathbb{C})
$$

up to conjugacy. If we assume here that (i) the parameters $\theta_i$’s are not integers, and that (ii) the singularities $[\lambda_\alpha]$’s are not logarithmic (i.e., apparent), then the local monodromies around the $[t_i]$’s and $[\lambda_\alpha]$’s become respectively (conjugate to)

$$
\left( -\exp(\pi\sqrt{-1} \theta_i) \ 0 \right) \text{ and } \left( 0 \ -\exp(-\pi\sqrt{-1} \theta_i) \right).
$$

Keeping this in mind and viewing the $\theta_i$’s as fixed (non-integral) constants, let us define our space $E$ to be the set of equations having the form (1) and satisfying
assumption (ii). Since it follows from (2) and assumption (ii) (under a generic condition) that the parameters $k, H_i (i = 0, \ldots, m)$ are described as certain functions of the other parameters of equation (1), we find that the space $E$ thus defined can be locally parametrized by the parameters $(\vec{t}, \tau, \vec{\lambda}, \vec{\mu})$. (We have introduced here the vector notation $\vec{t} = (t_1, \ldots, t_m) (t_0 = 0)$, $\vec{\lambda} = (\lambda_0, \ldots, \lambda_m)$, $\vec{\mu} = (\mu_0, \ldots, \mu_m)$.) Having finished the (local) description of the space $E$ of equations, we are now ready to write out the specific form of the pulled-back 2-form $F^*\omega$ (in terms of the coordinate parameters $(\vec{t}, \tau, \vec{\lambda}, \vec{\mu})$). (In view of the construction of $E$, the space $R$ of representations will correspondingly be defined as

$$R = \{ \rho: \pi_1(X \setminus \{2m + 2 \text{ points}\}) \to \text{SL}(2, \mathbb{C}) ; \text{the local representations around the punctures are as in (3) up to conjugacy}\}/\sim,$$

where $/\sim$ means taking the quotient space by the conjugate action of the group $\text{SL}(2, \mathbb{C})$ on $R$.)

**Theorem 1 [3].** In terms of the local parameters $(\vec{t}, \tau, \vec{\lambda}, \vec{\mu})$ of $E$, the pulled-back 2-form $F^*\omega$ takes the form

$$F^*\omega = -2 \left( \sum_{\alpha=0}^{m} d\mu_\alpha \wedge d\lambda_\alpha - \sum_{i=1}^{m} dH_i \wedge dt_i - dK \wedge d\tau \right),$$

where

$$K = \frac{1}{2\pi\sqrt{-1}} \left[ k + \eta_1(\tau) \left( \sum_{\alpha=0}^{m} \lambda_\alpha \mu_\alpha - \sum_{i=1}^{m} t_i H_i \right) \right]$$

and the term $\eta_1(\tau)$ is defined by $\eta_1(\tau) = \zeta(z + 1, \tau) - \zeta(z, \tau)$.

In particular, it follows from this result that if we consider the space $E_0$ of differential equations (again having the form (1) and satisfying assumption (ii)) on the fixed elliptic curve $X$ (we therefore regard the parameter $\tau$ as a fixed constant), then the resulting 2-form $F^*\omega$ becomes

$$-2 \left( \sum_{\alpha=0}^{m} d\mu_\alpha \wedge d\lambda_\alpha - \sum_{i=1}^{m} dH_i \wedge dt_i \right);$$

and that formula can indeed be viewed as a special instance of Iwasaki's result [2].

As explained earlier, we turn next to describing the null-foliation of the 2-form $F^*\omega$, which is given by integrating the distribution $\{ \xi \in TE ; F^*\omega(\xi, \cdot) \equiv 0 \}$. For this we first note that the first term $-2 \left( \sum_{\alpha=0}^{m} d\mu_\alpha \wedge d\lambda_\alpha \right)$ of the right-hand side
of (5) is nondegenerate on the (locally defined) $(\bar{\lambda}, \bar{\mu})$-space. From this simple but useful observation it follows that the leaves of the null-foliation of $F^*\omega$ are transversal to the $(\bar{\lambda}, \bar{\mu})$-directions, or equivalently that any tangent vector $\xi$ satisfying $F^*\omega(\xi, \cdot) \equiv 0$ must be a linear combination of the vectors having the form

$$\begin{align*}
\mathcal{H}_i &= \frac{\partial}{\partial t_i} + \sum_{\alpha=0}^{m} \left( A^i_{\alpha} \frac{\partial}{\partial \lambda_{\alpha}} + B^i_{\alpha} \frac{\partial}{\partial \mu_{\alpha}} \right) \quad (i = 1, \ldots, m) \\
\mathcal{H}_\tau &= \frac{\partial}{\partial \tau} + \sum_{\alpha=0}^{m} \left( C_{\alpha} \frac{\partial}{\partial \lambda_{\alpha}} + D_{\alpha} \frac{\partial}{\partial \mu_{\alpha}} \right),
\end{align*}$$

where $A^i_{\alpha}, B^i_{\alpha}, C_{\alpha}, D_{\alpha}$ are some complex numbers. Moreover a simple calculation shows that the vectors $\mathcal{H}_i$'s and $\mathcal{H}_\tau$ above make $F^*\omega$ degenerate precisely when

$$\begin{align*}
A^i_{\alpha} &= \frac{\partial H_i}{\partial \mu_{\alpha}} \\
B^i_{\alpha} &= - \frac{\partial H_i}{\partial \lambda_{\alpha}} \\
\sum_{\alpha=0}^{m} \left( \frac{\partial H_j}{\partial \mu_{\alpha}} \frac{\partial H_i}{\partial \lambda_{\alpha}} - \frac{\partial H_i}{\partial \lambda_{\alpha}} \frac{\partial H_j}{\partial \mu_{\alpha}} \right) &= \frac{\partial H_j}{\partial t_i} - \frac{\partial H_i}{\partial t_j} \quad (j = 1, \ldots, m)
\end{align*}$$

and

$$\begin{align*}
C_{\alpha} &= \frac{\partial K}{\partial \mu_{\alpha}} \\
D_{\alpha} &= - \frac{\partial K}{\partial \lambda_{\alpha}} \\
\sum_{\alpha=0}^{m} \left( \frac{\partial H_j}{\partial \mu_{\alpha}} \frac{\partial K}{\partial \lambda_{\alpha}} - \frac{\partial H_j}{\partial \lambda_{\alpha}} \frac{\partial K}{\partial \mu_{\alpha}} \right) &= \frac{\partial H_j}{\partial \tau} - \frac{\partial K}{\partial t_j} \quad (j = 1, \ldots, m)
\end{align*}$$

respectively (see [4, pp.10-11]). To describe the null-foliation of $F^*\omega$, it thus remains to prove (or disprove) the third formulas of (6) and (7). Although they can be shown directly by substituting into them the specific forms of the functions (4), the calculation needed is quite complicated and lengthy. Instead, following Iwasaki [1], we have proved them via the residue calculus of certain meromorphic differentials on Riemann surfaces. (The third formula of (6) has already been shown by Okamoto [5] and Iwasaki [1].) In summary then, one concludes that the distribution in question has as a local basis the vector fields

$$\begin{align*}
\mathcal{H}_i &= \frac{\partial}{\partial t_i} + \sum_{\alpha=0}^{m} \left( \frac{\partial H_i}{\partial \mu_{\alpha}} \frac{\partial}{\partial \lambda_{\alpha}} - \frac{\partial H_i}{\partial \lambda_{\alpha}} \frac{\partial}{\partial \mu_{\alpha}} \right) \quad (i = 1, \ldots, m) \\
\mathcal{H}_\tau &= \frac{\partial}{\partial \tau} + \sum_{\alpha=0}^{m} \left( \frac{\partial K}{\partial \mu_{\alpha}} \frac{\partial}{\partial \lambda_{\alpha}} - \frac{\partial K}{\partial \lambda_{\alpha}} \frac{\partial}{\partial \mu_{\alpha}} \right).
\end{align*}$$

However, since these vector fields just describe the time-evolution directions of the parameters $\lambda_{\alpha}$'s and $\mu_{\alpha}$'s with respect to the time-parameters $t_i$'s and $\tau$ with Hamiltonians $H_i$'s and $K$ (see [4, pp.10-11]), we finally obtain the following result.
Theorem 2. The null-foliation of the pulled-back 2-form $F^*\omega$ is locally described by the completely integrable Hamiltonian system

$$\begin{align*}
\frac{d\lambda_\alpha}{\partial \tau} &= \sum_{i=1}^{m} \frac{\partial H_i}{\partial \mu_\alpha} dt_i + \frac{\partial K}{\partial \mu_\alpha} d\tau \\
\frac{d\mu_\alpha}{\partial \tau} &= -\sum_{i=1}^{m} \frac{\partial H_i}{\partial \lambda_\alpha} dt_i - \frac{\partial K}{\partial \lambda_\alpha} d\tau.
\end{align*}$$

(\alpha = 0, \ldots, m)

Just as before, if we regard the parameter $\tau$ as a fixed constant, the resulting Hamiltonian system becomes

$$\begin{align*}
\frac{d\lambda_\alpha}{\partial \tau} &= \sum_{i=1}^{m} \frac{\partial H_i}{\partial \mu_\alpha} dt_i \\
\frac{d\mu_\alpha}{\partial \tau} &= -\sum_{i=1}^{m} \frac{\partial H_i}{\partial \lambda_\alpha} dt_i;
\end{align*}$$

(\alpha = 0, \ldots, m)

and that system has been obtained by Okamoto [5] and Iwasaki [1].

REFERENCES


