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Quantum Calogero models and the degenerate double affine Hecke algebra

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Abstract

Algebraic structure of Calogero-Sutherland-type models is reviewed. Each of the models is related to representations of a degenerate version of the double affine Hecke algebra. From this viewpoint, wavefunctions of the models can be treated in unified manner.

1 Introduction

There are intimate relations between quantum mechanics and special functions. Wavefunctions for some models can explicitly be written in terms of suitable special functions. Recent studies on integrable quantum many-particle systems reveal that wavefunctions of some special cases can be written in terms of multivariable analogue of classical orthogonal polynomials. Furthermore, in those cases, properties of the wavefunctions can be treated in unified manner by using their algebraic structure.

An example of such models is the Sutherland model, which describes interacting particle on a unit circle [Su1, Su2, OP]:

$$\mathcal{H}_S = -N \sum_{j=1}^{N} \frac{\partial^2}{\partial \theta_j^2} + \frac{1}{2} \sum_{j<k} \frac{\beta(\beta-1)}{\sin^2[(\theta_j - \theta_k)/2]},$$

(1.1)

where $\beta$ is coupling constant and we assume $\beta$ is a non-negative positive integer. In this case, wavefunctions can be written by the so-called Jack polynomials.

Another example is the quantum Calogero model confined in harmonic potential [Ca1, Ca2, Su1, OP]:

$$\mathcal{H}_A = \frac{1}{2} \sum_{j=1}^{N} \left( -\frac{\partial^2}{\partial x_j^2} + x_j^2 \right) + \sum_{j<k} \frac{\beta(\beta-1)}{(x_j - x_k)^2},$$

(1.2)

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The subscripts "A" signify that this Hamiltonian is invariant under the action of the Weyl group of $A_{N-1}$-type. There also exists the model invariant under the action of the $B_N$-type Weyl group [OP, Y]:

$$H_B = \frac{1}{2} \sum_{j=1}^{N} \left\{-\frac{\partial^2}{\partial z_j^2} + z_j^2 + \frac{\gamma(\gamma-1)}{z_j^2}\right\} + \sum_{j<k} \left\{\frac{\beta(\beta-1)}{(z_j - z_k)^2} + \frac{\beta(\beta-1)}{(z_j + z_k)^2}\right\}. \quad (1.3)$$

(For the latter convenience, we use the letter "z" as the coordinates of the $B_N$-type model.) We remark that the model associated with the $C_N$-type Weyl group is equivalent to the $B_N$ case, and $D_N$-type model is obtained by setting $\gamma = 0$.

In these cases, polynomial part of wavefunctions can be regarded as multivariable generalization of the Hermite ($A_{N-1}$ case) and Laguerre ($B_N$ case) polynomials and has been studied by several authors [BF1, BF2, BF3, vD, Ka1, Ka2, So, UW]. In fact, these three models share the same algebraic structure, the degenerate double affine Hecke algebra. From this viewpoint, each of the models corresponds to individual representation of the degenerate double affine Hecke algebra. Furthermore, by using the intertwining operators between representations of the degenerate double affine Hecke algebra, several known results on the Jack polynomials can be mapped directly to those of the multivariable Hermite and Laguerre polynomials. As applications, we will construct raising operators and shift operators for such polynomials.

2 Dunkl-type operators and multivariable orthogonal polynomials

2.1 Jack polynomials and the Sutherland model

In this subsection, we define our notation and review the theory of symmetric and nonsymmetric Jack polynomials [M1, O, KS]. There are several ways to characterize the Jack polynomials; Here we define them as eigenfunctions of some operators. We note that we restrict ourselves to the case associated with the $A_{N-1}$-type Weyl group since we only use such case.

In the paper [Du1], Dunkl has introduced differential-exchange operators, now called "Dunkl operators", which are associated with root systems. For the $A_{N-1}$-type root system, the operators are defined as

$$D_j^A = \frac{\partial}{\partial x_j} + \beta \sum_{k \neq j} \frac{1 - s_{jk}}{x_j - x_k} \quad (j = 1, \ldots, N),$$

where $s_{ij}$ are elements of the symmetric group $S_N$. An element $s_{ij}$ acts on functions of $x_1, \ldots, x_N$ as an operator which permutes arguments $x_i$ and $x_j$. We remark that the operators $D_j$ preserve the space of polynomials of variables $x_1, \ldots, x_N$ which we denote $\mathbb{C}[x]$. These operators satisfy the following properties:

$$[D_i^A, D_j^A] = 0, \quad s_{ij} D_i^A = D_i^A s_{ij}, \quad s_{ij} D_k^A = D_k^A s_{ij} \quad (k \neq i, j),$$

$$[D_i^A, x_j] = \delta_{ij} \left(1 + \beta \sum_{k \neq i} s_{ik}\right) - (1 - \delta_{ij}) \beta s_{ij}. $$
Heckman introduced "global" Dunkl operators [He1], which are written as $x_jD^A_j$ in our notation. Heckman's operators do not commute each other. Cherednik introduced another version of Dunkl operators that mutually commute [Ch1] (see also [BGHP, KS]):

$$\hat{D}^A_j = x_jD^A_j + \beta \sum_{k<j} s_{jk}.$$ 

The algebra generated by the elements $x_j^{\pm 1}$, $\hat{D}^A_j$ and $s_{ij}$ is isomorphic to the degenerate double affine Hecke algebra $\mathcal{H}'$ associated with the $A_{N-1}$-type root system [Ch1, Ch2]. We remark that the elements $x_j^{\pm 1}$, $D^A_j$ and $s_{ij}$ also generate $\mathcal{H}'$ since $D^A_j$ and $\hat{D}^A_j$ are related through (2.1).

We denote by $\mathcal{H}'_0$ subalgebra of $\mathcal{H}'$ generated by $\hat{D}^A_j$ and $s_{ij}$, which is isomorphic to the degenerate affine Hecke algebra. We further denote by $\tilde{\mathcal{H}}'$ subalgebra of $\mathcal{H}'$ generated by $x_j$, $\hat{D}^A_j$ and $s_{ij}$. In terms of generators, the defining relations are

$$[\hat{D}^A_i, \hat{D}^A_j] = [x_i, x_j] = 0,$$

$$x_is_{ij} = s_{ij}x_j, \quad x_is_{jk} = s_{jk}x_i \quad (i \neq j, k),$$

$$s_j^2 = 1, \quad s_is_{j+1}s_j = s_{j+1}s_js_{j+1}, \quad [s_i, s_j] = 0 \quad (|i-j| \neq 1),$$

$$\hat{D}^A_{j+1}s_j - s_j\hat{D}^A_j = \beta, \quad s_j\hat{D}^A_{j+1} - \hat{D}^A_j s_j = \beta, \quad [s_i, \hat{D}^A_j] = 0 \quad (j \neq i, i + 1),$$

$$[\hat{D}^A_i, x_j] = \left\{ \begin{array}{ll} x_i + \beta \left( \sum_{k<i} x_k s_{ik} + \sum_{k>i} x_i s_{ik} \right) & (i = j), \\ -\beta x_is_{ij} & (i > j), \\ -\beta x_is_{ij} & (i < j), \end{array} \right.$$ 

where $s_j = s_{j,j+1}$ ($j = 1, \ldots, n - 1$) are the simple transpositions.

Since the operators $\hat{D}^A_j$ commute each other, they can be diagonalized simultaneously by suitable choice of bases of $\mathbb{C}[x]$ [BGHP, O, KS]. Such basis is called non-symmetric Jack polynomials. To define the non-symmetric Jack polynomials, we first introduce the ordering $\prec$; For two pairs $(\lambda, w)$, $(\mu, w')$ where $\lambda$, $\mu$ are partitions and $w, w' \in \mathfrak{S}_N$, we define the ordering $\prec$ as follows:

$$(\lambda, w) \prec (\mu, w) \iff \left\{ \begin{array}{ll} (i) \quad \mu \prec_0 \lambda, \\ (ii) \quad \text{if } \mu = \lambda \text{ then } w' \prec_\text{Bruhat} w, \end{array} \right.$$ 

where $\prec_0$ is the dominance ordering for partitions [M1], and $\prec_\text{Bruhat}$ is the Bruhat ordering for the elements of $\mathfrak{S}_N$ (see, for example, [Hu]).

**Definition 2.1** ([BGHP, O, KS]) An non-symmetric Jack polynomial $E^\lambda_w(x)$, labeled with the partition $\lambda = (\lambda_1, \ldots, \lambda_N)$ and the element $w \in \mathfrak{S}_N$, is characterized by the following properties:

(i) $E^\lambda_w(x) = x^\lambda_w + \sum_{(\mu, w') \prec (\lambda, w)} u^\lambda_{w'w} x^\mu_w$,

(ii) $E^\lambda_w(x)$ is joint eigenfunction for the operators $\hat{D}^A_j$,

where we have used the notation $x^\lambda_w = x^\lambda_{w(1)} \cdots x^\lambda_{w(N)}$. 
Note that our definition of the non-symmetric Jack polynomials is slightly different from the one in the references cited above.

Since the action of $\hat{D}_j^A$ on monomials $x_w^\lambda$ are given by

$$\hat{D}_j^A x_w^\lambda = (w(\lambda + \beta \delta))_j x_w^\lambda + \sum_{(\mu,w')<(\lambda,w)} v_{w,w'}^\lambda x_w^\mu, \quad (2.1)$$

with $\delta = (N-1, N-2, \cdots, 0)$, the action of $\hat{D}_j^A$ on the non-symmetric Jack polynomials can be evaluated as follows [BGHP, O, KS]:

$$\hat{D}_j^A E_w^\lambda(x) = (w(\lambda + \beta \delta))_j E_w^\lambda(x). \quad (2.2)$$

From a physical viewpoint, the operators $\hat{D}_j^A$ are related to the Sutherland model (1.1). To see the relation, we introduce "gauge-transformed" Hamiltonian $\overline{\mathcal{H}}_S$:

$$\overline{\mathcal{H}}_S = \text{Res} \left( \sum_{j=1}^{N} \left\{ \hat{D}_j^A - \frac{\beta}{2} (N-1) \right\} \right)^2$$

$$= \sum_{j=1}^{N} \left( x_j \frac{\partial}{\partial x_j} \right)^2 + \beta \sum_{j<k} (x_j - x_k) \left( x_j \frac{\partial}{\partial x_j} - x_k \frac{\partial}{\partial x_k} \right) + \frac{\beta^2}{12} N(N^2 - 1),$$

where $\text{Res} X$ means that action of $X$ is restricted to symmetric functions of the variables $x_1, \ldots, x_N$. If we make a kind of gauge transformation and a change of variables $x_j = \exp(i \theta_j)$, $\overline{\mathcal{H}}_S$ reduces to the Sutherland Hamiltonian (1.1):

$$\phi^{(\beta)}_{S} o \overline{\mathcal{H}}_S o (\phi^{(\beta)}_{S})^{-1} = \sum_{j=1}^{N} \left( x_j \frac{\partial}{\partial x_j} \right)^2 - \beta(\beta - 1) \sum_{j<k} 2x_j x_k (x_j - x_k)^2 = \mathcal{H}_S,$$

where $\phi^{(\beta)}_{S}(x) = \prod_{j<k} |x_j - x_k|^\beta \prod_{j=1}^{N} x_j^{2(\beta(N-1)/2)}$ is the ground state wavefunction of the model. The symmetric Jack polynomials appear as polynomial part of wavefunctions for excited states.

**Definition 2.2 ([M1])** The symmetric Jack polynomials $J^{(\beta)}_\lambda(x)$ are characterized by the following properties:

(i) $J_\lambda(x) = m_\lambda(x) + \sum_{\mu \subset \lambda} u_{\lambda\mu} m_\mu(x)$,

(ii) $J_\lambda(x)$ are eigenfunctions of the transformed Hamiltonian $\overline{\mathcal{H}}_S$,

where $m_\lambda$ are the monomial symmetric functions.

The symmetric Jack polynomials are obtained by symmetrizing $E_w^\lambda$, i.e.,

$$J_\lambda(x) = \frac{1}{\# \mathfrak{S}_N^{\lambda}} \prod_{v \in \mathfrak{S}_N} v(E_w^\lambda),$$

where $\mathfrak{S}_N^{\lambda}$ is a subgroup of $\mathfrak{S}_N$ that preserve $\lambda$. This relation follows from the fact that the right hand side satisfies both of the defining properties of the Jack polynomials.
As wavefunctions of the Hamiltonian $\mathcal{H}_S$, the following scalar product is naturally introduced:

$$ (f(x), g(x))^{(\rho)}_{\mathcal{J}} = \oint \cdots \oint f(x)g(x^{-1})\phi_S^{(\rho)}(x)\phi_S^{(\rho)}(x^{-1}) \frac{dx_1}{2\pi i x_1} \cdots \frac{dx_N}{2\pi i x_N}, $$

where the integration contour is the unit circle in the complex plane. This scalar product can alternatively be written as

$$ (f(x), g(x))^{(\rho)}_{\mathcal{J}} = (-1)^{\rho \cdot (N-1)/2} \left[ f\overline{g}(\phi_S^{(\rho)})^2 \right]_0, $$

where $[\cdot]_0$ stands for the constant term and $\overline{g} = g(x^{-1})$. By a direct calculation, we see that the operators $\hat{D}_j^A$ are self-adjoint with respect to the scalar product (2.3).

**Proposition 2.3** ([M1]) *The Jack polynomials $J_\lambda(x)$ are pairwise orthogonal with respect to the scalar product (2.3).*

**Proof.** We first introduce generating function of symmetric commuting operators [BGHP, Kal]:

$$ \hat{\Delta}_J(u) = \prod_{j=1}^N (u + \hat{D}_j^A). $$

If we expand $\hat{\Delta}_J(u)$ as polynomial in $u$, the coefficients form a set of symmetric commuting operators which contain $\mathcal{H}_S$. Using (2.2), we can evaluate the action of $\hat{\Delta}_J(u)$ on the Jack polynomials:

$$ \hat{\Delta}_J(u) J_\lambda^{(\beta)}(x) = \prod_{j=1}^N \{u + \lambda_{N-j+1} + \beta(j-1)\} J_\lambda^{(\beta)}(x). $$

(2.4)

Since all the eigenvalues of $\hat{\Delta}_J(u)$ are distinct and the operator $\hat{\Delta}_J(u)$ is self-adjoint, we conclude that the Jack polynomials $J_\lambda(x)$ are pairwise orthogonal with respect to the scalar product (2.3). \qed

The property below follows from the fact that the Jack polynomials form an orthogonal basis of the space of symmetric polynomials $\mathbb{C}[x]^{S_N}$:

$$ (J_\lambda^{(\beta)}(x), m_\mu(x))^{(\rho)}_{\mathcal{J}} = 0 \quad \text{for all } \mu \leq \rho \lambda. $$

One can use this relation instead of the second property of Definition 2.2.

### 2.2 Multivariable Hermite polynomials and $A_{N-1}$-type Calogero model

We introduce an analogue of the creation and annihilation operators:

$$ a_j^\dagger = \frac{1}{\sqrt{2}}(-D_j^A + x_j), \quad a_j = \frac{1}{\sqrt{2}}(D_j^A + x_j). $$
The commutation relations of these operators are the same as those of \( x_j \) and \( D_j^A \) by construction. We then make a gauge transformation on \( a_j^\dagger \) and \( a_j \):

\[
\tilde{a}_j = \phi_A^{-1} \circ a_j \circ \phi_A = \frac{1}{\sqrt{2}} \left( -\frac{\partial}{\partial x_j} + 2x_j - \beta \sum_{k(\neq j)} \frac{1-s_{jk}}{x_j-x_k} \right),
\]

\[
\tilde{a}_j = \phi_A^{-1} \circ a_j \circ \phi_A = \frac{1}{\sqrt{2}} \left( \frac{\partial}{\partial x_j} + \beta \sum_{k(\neq j)} \frac{1-s_{jk}}{x_j-x_k} \right),
\]

with \( \phi_A = \prod_{k=1}^N \exp(-x_k^2/2) \). Since this transformation leaves the commutation relations unchanged, we can introduce the following isomorphism:

\[
\rho^A(x_j) = a_j^\dagger, \quad \rho^A(D_j) = \tilde{a}_j, \quad \rho^A(s_{ij}) = s_{ij}.
\]

It should be remarked that this mapping has already been appeared implicitly in [UW], however, treated only as an isomorphism of the algebra. To construct eigenstates of \( \tilde{\mathfrak{H}}_A \), we should introduce intertwiner between two representations which will be discussed in the followings.

We can obtain a set of commuting operators by applying \( \rho^A \) to \( \hat{D}_j^A \):

\[
\tilde{h}_j^A = \rho^A(\hat{D}_j^A) = a_j^\dagger a_j + \beta \sum_{k(<j)} S_{jk}.
\]

The mapping \( \rho^A \) gives another representation of \( \tilde{\mathfrak{H}}' \) on \( \mathbb{C}[x] \). We then introduce intertwining operator \( \sigma^A \), which is a linear operator on \( \mathbb{C}[x] \) such that

\[
\sigma^A(f(x_1, \ldots, x_N)) = f(\tilde{a}_1^\dagger, \ldots, \tilde{a}_N^\dagger) \cdot 1 \quad \text{for all} \quad f(x_1, \ldots, x_N) \in \mathbb{C}[x].
\]

The intertwiner \( \sigma^A \) enjoys the following property.

**Theorem 2.4** \( \sigma^A(Qf(x)) = \rho^A(Q)\sigma^A(f(x)) \) for all \( Q \in \tilde{\mathfrak{H}}_0' \), \( f(x) \in \mathbb{C}[x] \).

**Proof.** Since both \( Q \) and \( f(x) \) are elements of \( \tilde{\mathfrak{H}}' \), it suffices to prove \( \sigma^A(P \cdot 1) = \rho^A(P) \cdot 1 \) for all \( P \in \tilde{\mathfrak{H}}' \). We then note that every element \( P \) of \( \tilde{\mathfrak{H}}' \) can be represented in the following form:

\[
P = \sum_n \sum_{w \in \mathfrak{S}_N} p_{n,w}(x)(\hat{D}_1^A)^{n_1} \cdots (\hat{D}_N^A)^{n_N} w, \tag{2.5}
\]

where \( p_{n,w}(x) \) are some polynomials. Considering the action of (2.5) on \( 1 \), we have

\[
P \cdot 1 = \sum_{n(n_1=0)} \sum_{w \in \mathfrak{S}_N} p_{n,w}(x) \beta^{n_2} \cdots ((N-1)\beta)^{n_N},
\]

since \( w \cdot 1 = 1 \) for all \( w \in \mathfrak{S}_N \) and \( \hat{D}_j^A \cdot 1 = \beta(j-1) \) for all \( j \). On the other hand, applying \( \rho^A \) to (2.5), we have

\[
\rho^A(P) = \sum_n \sum_{w \in \mathfrak{S}_N} p_{n,w}(\tilde{a}_1^\dagger)(\tilde{h}_1^A)^{n_1} \cdots (\tilde{h}_N^A)^{n_N} w.
\]

Since \( \tilde{h}_j^A \cdot 1 = \beta(j-1) \) for all \( j \), we conclude that \( \sigma^A(P \cdot 1) = \rho^A(P) \cdot 1 \) for all \( P \in \tilde{\mathfrak{H}}' \). \( \square \)
The representation $\rho^A$ is related to the $A_{N-1}$-type Calogero model. If we define $\overline{\mathcal{H}}_A$ as

$$\overline{\mathcal{H}}_A = \text{Res} \left( \sum_{j=1}^{N} \tilde{h}_j^A \right) - \frac{\beta}{2} N(N - 1)$$

$$= \frac{1}{2} \sum_{j=1}^{n} \left( -\frac{\partial^2}{\partial x_j^2} + 2x_j \frac{\partial}{\partial x_j} \right) - \beta \sum_{j<k} \frac{1}{x_j - x_k} \left( \frac{\partial}{\partial x_j} - \frac{\partial}{\partial x_k} \right),$$

we can obtain the $A_{N-1}$-type Calogero Hamiltonian (1.2) via gauge transformation:

$$\mathcal{H}_A = \phi_A^{(\beta)} \circ \overline{\mathcal{H}}_A \circ (\phi_A^{(\beta)})^{-1} + \frac{N}{2} + \frac{\beta}{2} N(N - 1),$$

with $\phi_A^{(\beta)} = \Pi_{j<k} |x_j - x_k|^{\beta} \Pi_{j=N} \exp(-x_j^2/2)$ ground state wavefunction.

We then introduce scalar product for this case:

$$\langle f, g \rangle_H^{(\beta)} = \int_{-\infty}^{\infty} \cdots \int_{-\infty}^{\infty} f(x) g(x) (\phi_A^{(\beta)})^2 dx_1 \cdots dx_N$$

(2.6)

By a direct calculation, we see that the operator $\tilde{a}_j^\dagger$ is adjoint of $\tilde{a}_j$ with respect to the scalar product (2.6) for all $j = 1, \ldots, N$. Note that $x_j (= (\rho^A)^{-1}(\tilde{a}_j))$ is not adjoint of $D_j (= (\rho^A)^{-1}(a_j))$ for the Jack case.

Multivariable Hermite polynomials are defined by using this scalar product [BF1, vD]. In fact, the definition in [BF1] and that in [vD] are slightly different. Here we shall follow [vD]:

**Definition 2.5 ([vD])** Multivariable Hermite polynomials $H_\lambda^{(\beta)}(x)$ are characterized by the following properties:

(i) $H_\lambda^{(\beta)}(x) = m_\lambda(x) + \sum_{\mu <_D \lambda} u_{\langle \mu \rangle}^{\lambda} m_\mu(x)$,

(ii) $\langle H_\lambda^{(\beta)}(x), m_\mu(x) \rangle_H^{(\beta)} = 0$ for all $\mu <_D \lambda$.

Using the intertwiner $\sigma^A$, we can construct an operator representation of $H_\lambda^{(\beta)}(x)$.

**Proposition 2.6 ([Kal, UW])** Multivariable Hermite polynomials $H_\lambda^{(\beta)}(x)$ are related to the Jack polynomials as follows:

$$H_\lambda^{(\beta)}(x) = 2^{-|\lambda|/2} \sigma^A(J_\lambda^{(\beta)}(x)) = 2^{-|\lambda|/2} J_\lambda^{(\beta)}(\tilde{a}_1^\dagger, \ldots, \tilde{a}_N^\dagger) \cdot 1.$$

**Proof.** We can easily see that $2^{|\lambda|/2} J_\lambda(\tilde{a}_1^\dagger, \ldots, \tilde{a}_N^\dagger) \cdot 1$ satisfy the condition (i) of Definition 2.5. Hence it suffices to show (ii). Applying $\sigma^A$ to (2.4), we have

$$\hat{\Delta}_H(u) J_\lambda^{(\beta)}(\tilde{a}_1^\dagger, \ldots, \tilde{a}_N^\dagger) \cdot 1 = \prod_{j=1}^{N} \{ u + \lambda_{N-j+1} + \beta(j - 1) \} J_\lambda^{(\beta)}(\tilde{a}_1^\dagger, \ldots, \tilde{a}_N^\dagger) \cdot 1,$$

where we denote $\hat{\Delta}_H(u) = \rho^A(\hat{\Delta}_J(u)) = \prod_{j=1}^{N} (u + \tilde{h}_j^A)$. Since all the eigenvalues of $\hat{\Delta}_H(u)$ are distinct and the operator $\hat{\Delta}_H(u)$ are self-adjoint with respect to the scalar product (2.6), we conclude that the polynomials $J_\lambda^{(\beta)}(\tilde{a}_1^\dagger) \cdot 1$ are orthogonal with respect to the scalar product.
(2.6). On the other hand, one may know that the polynomials $J_{\lambda}^{(\beta)}(\sim a\dagger)\cdot 1$ form an orthogonal basis of $\mathbb{C}[x]\mathcal{O}_{N}$ by considering the leading term. It follows that $\langle J_{\lambda}^{(\beta)}(\sim a\dagger)\cdot 1, m_{\mu}\rangle_{t}^{(\beta)} = 0$ for all $\mu <_{\rho} \lambda$, which proves the theorem.

It should be noted that Ujino and Wadati [UW] have shown that $J_{\lambda}^{(\beta)}(\sim a\dagger)\cdot 1$ diagonalize the first two of the family of commuting operators that contains $\hat{H}_{A}_{t}$. The proof given here is essentially the same as that given in [Kal].

The scalar product $\langle \cdot, \cdot \rangle_{H}^{(\beta)}$ induces another scalar product on $\mathbb{C}[x]$: $\langle f(x), g(x) \rangle_{A} = \langle f(\sim a\dagger)\cdot 1, g(\sim a\dagger)\cdot 1 \rangle_{H}^{(\beta)}$. This gives another example of scalar product which makes the Jack polynomials orthogonal.

On the other hand, Dunkl [Du2] introduced the scalar product $\langle f(\hat{D}^{A})g(x) \rangle_{0}$. These scalar products coincide up to a constant factor:

$$\langle f(x), g(x) \rangle_{A} = \langle 1, f(\sim a\dagger)\cdot 1, g(\sim a\dagger)\cdot 1 \rangle_{H}^{(\beta)}$$

We shall evaluate the value $\langle 1, 1 \rangle_{H}^{(\beta)}$ in section 4.2. (See Proposition 4.8 below.)

### 2.3 Multivariable Laguerre polynomials and $B_{N}$-type Calogero model

Dunkl operators associated with the $B_{N}$-type root system are defined as follows [Du1, Y]:

$$D_{j}^{B} = \frac{\partial}{\partial z_{j}} + \beta \sum_{k(\neq j)} \left( \frac{1-s_{jk}}{z_{j}-z_{k}} + \frac{1-t_{jk}s_{jk}}{z_{j}+z_{k}} \right) + \frac{1-t_{j}}{z_{j}}, \quad (2.7)$$

where $s_{jk}$ and $t_{j}$ are elements of the $B_{N}$-type Weyl group. An element $s_{ij}$ acts as same as in the $A_{N-1}$-case and $t_{j}$ acts as sign-change, i.e. replaces the coordinate $z_{j}$ by $-z_{j}$. The commutation relations of the $B_{N}$-type Dunkl operators are

$$[D_{i}^{B}, D_{j}^{B}] = 0, \quad s_{ij}D_{i}^{B} = D_{i}^{B} s_{ij}, \quad s_{ij}D_{k}^{B} = D_{k}^{B} s_{ij} \quad (k \neq i, j),$$

$$t_{j}D_{i}^{B} = -D_{i}^{B} t_{j}, \quad t_{j}D_{k}^{B} = D_{k}^{B} t_{j} \quad (k \neq j),$$

$$[D_{i}^{B}, z_{j}] = \delta_{ij} \left\{ 1 + \beta \sum_{k(\neq i)} (s_{ik} + t_{i}t_{k}s_{ik}) + 2\gamma t_{j} \right\} - (1-\delta_{ij}) \beta (s_{ij} - t_{i}t_{k}s_{ik}).$$

We then define Cherednik-type commuting operators associated with (2.7):

$$\hat{D}_{j}^{B} = z_{j}D_{j}^{B} + \beta \sum_{k(<j)} (s_{jk} + t_{j}t_{k}s_{ik}).$$

Note that the operators $\hat{D}_{j}^{B}$ do not coincide with the Cherednik operators associated with the $B_{N}$-type Weyl group.

**Lemma 2.7** All of the operators $\hat{D}_{j}^{B}$, $s_{ij}$, $t_{j}$ and $z_{j}^{2}$ preserve $\mathbb{C}[z_{1}^{2}, \ldots, z_{N}^{2}]$. 
\textbf{Proof.} Only $\hat{D}^B_j$ need to prove. We introduce the notation $\text{Res}^{(t)}(X)$ which means the action of the operator $X$ is restricted to the functions with the symmetry $t_j f(z) = f(z)$. Under this restriction, $\text{Res}^{(t)}(\hat{D}^B_j)$ is equivalent to $2\hat{D}^A_j$ if we make a change of the variables $x_j = z_j^2/2$. Since $\hat{D}^A_j$ preserve $\mathbb{C}[x]$, the operators $\hat{D}^B_j$ preserve $\mathbb{C}[z^2]$.

From these facts, we can define representation $\iota$ of $\tilde{\mathfrak{g}}'$ on $\mathbb{C}[z^2]$:

$$\iota(x_j) = \frac{1}{2} z_j^2, \quad \iota(\hat{D}^A_j) = \frac{1}{2} \hat{D}^B_j, \quad \iota(s_{ij}) = s_{ij}.$$ 

We now introduce creation and annihilation operators for the $B_N$ case:

$$b_j^\dagger = \frac{1}{\sqrt{2}} (-D_j + z_j), \quad b_j = \frac{1}{\sqrt{2}} (D_j + z_j).$$

The commutation relations of these operators are the same as those of $z_j$ and $D_j^B$ by construction. We then make a gauge transformation on $b_j^\dagger$ and $b_j$:

$$\tilde{b}_j^\dagger = \tilde{\phi}_B^{-1} \circ b_j^\dagger \circ \tilde{\phi}_B = \frac{1}{\sqrt{2}} \left\{ -\frac{\partial}{\partial z_j} + 2z_j - \beta \sum_{k \neq j} \left( \frac{1 - s_{jk}}{z_j - z_k} + \frac{1 - t_j t_k s_{jk}}{z_j + z_k} \right) + \gamma \frac{1 - t_j}{z_j} \right\},$$

$$\tilde{b}_j = \tilde{\phi}_B^{-1} \circ b_j \circ \tilde{\phi}_B = \frac{1}{\sqrt{2}} \left\{ \frac{\partial}{\partial z_j} + \beta \sum_{k \neq j} \left( \frac{1 - s_{jk}}{z_j - z_k} + \frac{1 - t_j t_k s_{jk}}{z_j + z_k} \right) + \gamma \frac{1 - t_j}{z_j} \right\},$$

with $\tilde{\phi}_B = \prod_{k=1}^N \exp(-z_k^2/2)$. Since this transformation leaves the commutation relations unchanged, we can define the following algebra isomorphism:

$$\kappa(x_j) = \tilde{b}_j^\dagger, \quad \kappa(D_j^B) = \tilde{b}_j, \quad \kappa(s_{ij}) = s_{ij}, \quad \kappa(t_j) = t_j.$$ 

We then define the operators $\tilde{h}_j^B$ as follows:

$$\tilde{h}_j^B = \kappa(\hat{D}_j^B) = \tilde{b}_j^\dagger \tilde{b}_j + \beta \sum_{k < j} (s_{jk} + t_j t_k s_{ik}).$$

\textbf{Lemma 2.8} The operators $\tilde{h}_j^B$ and $(\tilde{b}_j^\dagger)^2$ preserve $\mathbb{C}[z_1^2, \ldots, z_N^2]$.

\textbf{Proof.} Since the operators $D_j^B$ preserve $\mathbb{C}[z_1, \ldots, z_N]$, it is clear that both $\tilde{h}_j^B$ and $(\tilde{b}_j^\dagger)^2$ also preserve $\mathbb{C}[z_1, \ldots, z_N]$. Then it suffices to prove $[t_i, \tilde{h}_j^B] = [t_i, (\tilde{b}_j^\dagger)^2] = 0$ for all $i, j$, which can be proved by a direct calculation. \hfill \Box

Using both $\iota$ and $\kappa$, we introduce another representation of $\tilde{\mathfrak{g}}'$ on $\mathbb{C}[z^2]$:

$$\rho^B(x_j) = \kappa(\iota(x_j)) = \frac{1}{2} (\tilde{b}_j^\dagger)^2, \quad \rho^B(\hat{D}^A_j) = \kappa(\iota(\hat{D}^A_j)) = \frac{1}{2} \tilde{h}_j^B, \quad \rho^B(s_{ij}) = s_{ij}.$$ 

We introduce a linear map of $\mathbb{C}[x]$ to $\mathbb{C}[z^2]$ by using $\rho^B$:

$$\sigma^B(f(x_1, \ldots, x_N)) = f((\tilde{b}_j^\dagger)^2/2, \ldots, (\tilde{b}_N^\dagger)^2/2) \cdot 1 \quad \text{for all } f(x_1, \ldots, x_N) \in \mathbb{C}[x].$$

As in the $A_{N-1}$-case, the intertwiner $\sigma^B$ enjoys the following property.
Theorem 2.9 \( \sigma^B(Qf(x)) = \rho^B(Q)\sigma^B(f(x)) \cdot 1 \) for all \( Q \in \mathcal{B}_0, f(x_1, \ldots, x_N) \in \mathbb{C}[x] \).

Proof is given in the same fashion as Theorem 2.4, so we omit details.

The operators \( \bar{b}_j^\dagger \) and \( \bar{b}_j \) are related to the \( B_N \)-type Calogero Hamiltonian (1.3); If we define \( \overline{\mathcal{H}}_B \) as

\[
\overline{\mathcal{H}}_B = \text{Res} \left( \sum_{j=1}^N \tilde{h}_j^B \right) - \beta N(N-1) = \frac{1}{2} \sum_{j=1}^n \left( -\frac{\partial^2}{\partial z_j^2} + 2z_j \frac{\partial}{\partial z_j} - \frac{2\gamma}{z_j} \frac{\partial}{\partial z_j} \right) - 2\beta \sum_{j<k} \frac{1}{z_j^2 - z_k^2} \left( z_j \frac{\partial}{\partial z_j} - z_k \frac{\partial}{\partial z_k} \right),
\]

we can obtain the Hamiltonian (1.3) via gauge transformation:

\[
\mathcal{H}_B = \phi_B^{(\beta)} \circ \overline{\mathcal{H}}_B \circ (\phi_B^{(\beta)})^{-1} + \left( \frac{1}{2} + \gamma \right)N + \beta N(N-1),
\]

with \( \phi_B^{(\beta)} = \prod_{j<k} |z_j^2 - z_k^2|^{\beta} \prod_{j=1}^N |z_j|^\gamma \exp(-z_j^2/2) \) ground state wavefunction of (1.3).

Scalar product associated with this model is

\[
\langle f(z), g(z) \rangle_L^{(\beta)} = \int_{-\infty}^{\infty} \cdots \int_{-\infty}^{\infty} f(z)g(z)(\phi_B^{(\beta)})^2 \mathrm{d}z_1 \cdots \mathrm{d}z_N.
\]

By a direct calculation, we can show that the operator \( \bar{b}_j^\dagger \) is adjoint of \( \bar{b}_j \) with respect to the scalar product (2.8), and hence the operator \( \tilde{h}_j^B \) is self-adjoint for all \( j = 1, \ldots, N \).

Now we define multivariable Laguerre polynomials [vD].

**Definition 2.10 ([vD])** Multivariable Laguerre polynomials \( L_\lambda^{(\beta)}(z) \) are characterized by the following properties:

(i) \( L_\lambda^{(\beta)}(z) = m_\lambda(z^2) + \sum_{\mu<\lambda} u_{\lambda\mu}m_\mu(z^2), \)

(ii) \( \langle L_\lambda^{(\beta)}(z), m_\mu(z^2) \rangle_L^{(\beta)} = 0 \) for all \( \mu <_D \lambda. \)

We can construct an operator representation of \( L_\lambda^{(\beta)}(z) \) by using the intertwiner \( \sigma^B \).

**Proposition 2.11 ([Ka2])** Multivariable Laguerre polynomials \( L_\lambda^{(\beta)}(z) \) are related to the Jack polynomials as follows:

\[
L_\lambda^{(\beta)}(z) = \sigma^B(J_\lambda^{(\beta)}(x)) = J_\lambda^{(\beta)}((\bar{b}_j^\dagger)^2/2) \cdot 1.
\]

One can prove this statement in the same way as Proposition 2.6, so we omit details.

**3 Construction of raising operators**

As is shown in the previous section, the multivariable Hermite and Laguerre polynomials are expressed in terms of the Jack polynomials whose arguments are Dunkl-type operators. Some properties of the multivariable Hermite and Laguerre polynomials are obtained directly from those of the Jack polynomials simply by applying \( \rho^A \) or \( \rho^B \). As an example, we will construct raising operators for the polynomials.
Lapointe and Vinet constructed raising operators for the Jack polynomials [LV]. Using their raising operators, they obtained Rodrigues-type formula for the Jack polynomials. Raising operators for the multivariable Hermite polynomials have been constructed by Ujino and Wadati [UW]. The raising operators above are constructed by the use of Dunkl operators of Heckman-type (non-commutative). In those cases, relation to the degenerate double affine Hecke algebra is still unclear.

On the other hand, Kirillov and Noumi gave another expression of raising operators by using Cherednik operators [KN]. In our notation, their raising operators are expressed as the following form:

$$B_m^J = \sum_{k_1 < \cdots < k_m} x_{k_1} x_{k_2} \cdots x_{k_m} \left( \hat{D}_{k_1}^A + \beta(2 - k_1) \right) \times \left( \hat{D}_{k_2}^A + \beta(3 - k_2) \right) \cdots \left( \hat{D}_{k_m}^A + \beta(m - k_m + 1) \right).$$

We recall an important property of these operators.

**Theorem 3.1 ([KN])** Action of the operators $B_m^J \in \tilde{\mathfrak{H}}'$ on the Jack polynomials are given by

$$B_m^J J^{(\beta)(\lambda)}(x) = \prod_{j=1}^{m} (\lambda_j + \beta(m - j + 1)) J^{(\beta)(\lambda \{1^{m})}x),$$

where $\lambda + (1^{m}) = (\lambda_1 + 1, \ldots, \lambda_N + 1)$.

Applying $\sigma^A$ or $\sigma^B$ to $B_m^J$, we obtain raising operators for the Hermite-case or the Laguerre-case respectively:

$$B_m^H = \sum_{k_1 < \cdots < k_m} \tilde{a}_{k_1} \tilde{a}_{k_2} \cdots \tilde{a}_{k_m} (\tilde{h}_{k_1}^A + \beta(2 - k_1)) \times (\tilde{h}_{k_2}^A + \beta(3 - k_2)) \cdots (\tilde{h}_{k_m}^A + \beta(m - k_m + 1)),$$

$$B_m^L = \sum_{k_1 < \cdots < k_m} \tilde{b}_{k_1} \tilde{b}_{k_2} \cdots \tilde{b}_{k_m} (\tilde{h}_{k_1}^B + \beta(2 - k_1)) \times (\tilde{h}_{k_2}^B + \beta(3 - k_2)) \cdots (\tilde{h}_{k_m}^B + \beta(m - k_m + 1)).$$

Form the theorem 3.1 and the propositions 2.6, 2.11, it immediately follows that:

**Proposition 3.2**

(i) $B_m^H H^{(\beta)}(\lambda)(x) = 2^{-m/2} \prod_{j=1}^{m} (\lambda_j + \beta(m - j + 1)) H^{(\beta)}(\lambda \{1^{m})}x),$

(ii) $B_m^L L^{(\beta)}(\lambda)(z) = \prod_{j=1}^{m} (\lambda_j + \beta(m - j + 1)) L^{(\beta)}(\lambda \{1^{m})}z).$

Applying the raising operators repeatedly, one can obtain Rodrigues-type formulas for the multivariable Hermite and Laguerre polynomials:

$$H^{(\beta)}(\lambda)(x) = 2^{\lambda_1/2} \prod_{(i,j) \in \lambda} (\lambda_i - j + \beta(\lambda'_j - i + 1))^{-1} (B^H_N)^{\lambda_N-\lambda_{N-1}} \cdots (B^H_1)^{\lambda_1-1 \lambda_2} 1,$$

$$L^{(\beta)}(\lambda)(z) = \prod_{(i,j) \in \lambda} (\lambda_i - j + \beta(\lambda'_j - i + 1))^{-1} (B^L_N)^{\lambda_N-\lambda_{N-1}} \cdots (B^L_1)^{\lambda_1-1 \lambda_2} 1,$$

where $\lambda' = (\lambda'_1, \lambda'_2, \ldots)$ is the conjugate partition to $\lambda$. 
4 Construction of shift operators

In this section, we construct shift operators for the multivariable Hermite and Laguerre polynomials. Each of such shift operators are related to one of the scalar products (2.3), (2.6), (2.8) (see the theorems 4.2, 4.6 and 4.9 below). However properties related to scalar product cannot be obtained simply by applying $\rho^A$ or $\rho^B$. It require a little more effort to construct shift operators.

It should be noted that the notion of hypergeometric shift operators were introduced originally by Opdam, and Heckman gave an expression by using the Dunkl operators [He2, He1]. Since our construction of shift operators are based on the Cherednik operators, relation to the degenerate double affine Hecke algebra becomes clear.

4.1 Shift operators for the Jack polynomials

In this subsection, we review the method of constructing shift operators for the Jack polynomials by the use of the Cherednik operators. Our method is based on the lecture note of Kirillov Jr. [Ki]; All results given in this section can be obtained by taking limiting procedure on those of [Ki]. However, all proofs given here are algebraic and we avoid using limiting procedure so that we can apply the results to the Hermite and Laguerre cases.

Consider elements of $\mathfrak{H}_0'$ that have the following forms:

$$\mathcal{X} = \prod_{i<j}(x_i - x_j), \quad \mathcal{Y} = \prod_{i<j}(\beta - \hat{D}_i^A + \hat{D}_j^A), \quad \hat{\mathcal{Y}} = \prod_{i<j}(-\beta - \hat{D}_i^A + \hat{D}_j^A). \quad (4.1)$$

We note that these operators preserve $\mathbb{C}[x]$. From the defining relations of $\mathfrak{H}_0'$, we know that

$$\begin{align*}
(s_j + 1)(-\beta - \hat{D}_j^A + \hat{D}_{j+1}^A) &= (-\beta - \hat{D}_{j+1}^A + \hat{D}_j^A)(s_j - 1), \\
(s_j - 1)(\beta - \hat{D}_j^A + \hat{D}_{j+1}^A) &= (\beta - \hat{D}_{j+1}^A + \hat{D}_j^A)(s_j + 1), \\
s_j(c - \hat{D}_j^A + \hat{D}_k^A)(c - \hat{D}_{j+1}^A + \hat{D}_k^A) &= (c - \hat{D}_j^A + \hat{D}_k^A)(c - \hat{D}_{j+1}^A + \hat{D}_k^A)s_j,
\end{align*} \quad (4.2)$$

with $c$ arbitrary constant and $k \neq j, j + 1$. Then, if we define $\mathbb{C}[x]^{\mathfrak{S}_N}$ and $\mathbb{C}[x]^{-\mathfrak{S}_N}$ as

$$\begin{align*}
\mathbb{C}[x]^{\mathfrak{S}_N} &= \{ f(x) \in \mathbb{C}[x] \mid (s_j - 1)f(x) = 0 \}, \\
\mathbb{C}[x]^{-\mathfrak{S}_N} &= \{ f(x) \in \mathbb{C}[x] \mid (s_j + 1)f(x) = 0 \},
\end{align*}$$

we see that

$$\begin{align*}
\mathcal{X} \in \mathbb{C}[x]^{-\mathfrak{S}_N}, \quad \mathcal{Y} \left( \mathbb{C}[x]^{\mathfrak{S}_N} \right) = \mathbb{C}[x]^{-\mathfrak{S}_N}, \quad \hat{\mathcal{Y}} \left( \mathbb{C}[x]^{-\mathfrak{S}_N} \right) = \mathbb{C}[x]^\mathfrak{S}_N. \quad (4.3)
\end{align*}$$

We now introduce shift operators for the Jack polynomials as $G_j = \mathcal{X}^{-1}\mathcal{Y}_j, \hat{G}_j = \hat{\mathcal{Y}}_j\mathcal{X}$. The operators $G_j$ and $\hat{G}_j$ enjoy the following properties:

**Lemma 4.1** ([Ki]) (i) $G_j J_{\lambda}^{(\beta)} \mathfrak{R}_+ \delta, \hat{G}_j J_{\lambda}^{(\beta+1)} \in \mathbb{C}[x]^{\mathfrak{S}_N}$.

(ii) $G_j J_{\lambda}^{(\beta)} = c_{\lambda}^{(\beta+1)} m_{\lambda} + \text{ "lower terms" with respect to } <_D,$

$\text{with } c_{\lambda}^{(\beta+1)} = \prod_{i<j} \{ \lambda_{N-j+1} - \lambda_{N-i+1} + j - i + \beta(j - i - 1) \}.$
(iii) $\hat{G}_j f^{(\beta+1)} = c_\lambda^{(\beta+1)} m_{\lambda+\delta} + \text{"lower terms" with respect to } <_D$, \\
with $c_\lambda^{(\beta+1)} = \prod_{i<j} \{ \lambda_{N-j+1} - \lambda_{N-i+1} + j - i + \beta(j-i+1) \}$.

Proof. (i) Follows from (4.3).

(ii) For the longest element $w_0$ of $\mathfrak{S}_N$, i.e. $w_0(j) = N - j + 1$, equation (2.1) reduces to

$$\hat{D}_j^A x_{w_0}^\lambda = (\lambda_{N-j+1} + \beta(j-1))x_{w_0}^\lambda + \sum_{(\mu,w') < (\lambda,w_0)} u_{w_0 w}^\mu x_{w'}^\mu.$$ 

Using this relation, we can calculate the action of $\mathcal{Y}_j$:

$$\mathcal{Y}_j f^{(\beta)}_{\lambda+\delta} = \prod_{i<j} (\beta - \hat{D}_i^A + \hat{D}_j^A)(x_{w_0}^{\lambda+\delta} + \text{"lower terms" with respect to } <)$$

$$= c_\lambda^{(\beta+1)} x_{w_0}^{\lambda+\delta} + \text{"lower terms" with respect to } <. \quad (4.4)$$

On the other hand, (4.3) implies that $\mathcal{Y}_j f^{(\beta)}_{\lambda+\delta}$ is divisible by $\mathcal{X}$. Together with (i), this concludes the proof.

(iii) Can also be proved in similar way. \(\square\)

The following theorem implies that $\hat{G}_j$ is, in a sense, adjoint of $G_j$:

**Theorem 4.2 ([Ki])** For $f, g \in \mathbb{C}[x]^{\mathfrak{S}_N}$, $\langle G_j f, g \rangle_j^{(\beta+1)} = \langle f, \hat{G}_j g \rangle_j^{(\beta)}$.

To prove this theorem, we introduce symmetrizer $\mathcal{P}_+$ and anti-symmetrizer $\mathcal{P}_-$ as

$$\mathcal{P}_+ = \frac{1}{\# \mathfrak{S}_N} \sum_{w \in \mathfrak{S}_N} w, \quad \mathcal{P}_- = \frac{1}{\# \mathfrak{S}_N} \sum_{w \in \mathfrak{S}_N} (-1)^{l(w)} w,$$

where $l(w)$ is the length of the element $w$. We further prepare a lemma.

**Lemma 4.3 ([Ki])** $\mathcal{P}_-(\mathcal{Y}_j - \hat{\mathcal{Y}}_j) = \sum_j \hat{g}_j(\hat{D}_1^A, \ldots, \hat{D}_N^A)(s_j - 1)$ for some $\hat{g}_j(x_1, \ldots, x_N) \in \mathbb{C}[x]$.

It should be remarked that this lemma is a degenerate version of a lemma given in [Ki]. We will give a proof in Appendix A for reader’s convenience.

Now we go back to the proof of Theorem 4.2.

**Proof of Theorem 4.2.** It is clear that $\mathcal{P}_+$ does not affect constant term of polynomials. Then, for $f, g \in \mathbb{C}[x]^{\mathfrak{S}_N}$, we know that

$$\langle G_j f, g \rangle_j^{(\beta+1)} = (-1)^{\beta(N-1)/2} [\langle \mathcal{X}^{-1} \mathcal{Y}_j f \rangle \bar{\bar{g}}(\phi_0^{(\beta+1)})^2]_0$$

$$= (-1)^{\beta(N-1)/2} [\mathcal{P}_+ \langle (\mathcal{Y}_j f \mathcal{X})(\phi_0^{(\beta)})^2 \rangle \bar{\bar{g}}]_0$$

$$= (-1)^{\beta N(N-1)/2} [\mathcal{P}_- \langle (\mathcal{Y}_j f \mathcal{X})(\phi_0^{(\beta)})^2 \rangle \bar{\bar{g}}]_0.$$ 

From Lemma 4.3, we see that $\mathcal{P}_-(\mathcal{Y}_j - \hat{\mathcal{Y}}_j) f = 0$ for all $f \in \mathbb{C}[x]^{\mathfrak{S}_N}$. Hence we can replace $\mathcal{Y}_j$ by $\hat{\mathcal{Y}}_j$:

$$\langle G_j f, g \rangle_j^{(\beta+1)} = (-1)^{\beta N(N-1)/2} [\mathcal{P}_- \langle (\hat{\mathcal{Y}}_j f \mathcal{X})(\phi_0^{(\beta)})^2 \rangle \bar{\bar{g}}]_0 = \langle \hat{\mathcal{Y}}_j f, \mathcal{X} g \rangle_j^{(\beta)} = \langle f, \hat{G}_j g \rangle_j^{(\beta)}.$$ 

In the last equality, we have used the self-adjointness of $\hat{D}_j^A$. \(\square\)

Using Theorem 4.2, we can evaluate the action of $G_j$ and $\hat{G}_j$ on the Jack polynomials.
Proposition 4.4 ([Kii]) \( G_J^{(\rho)} = c^{(\beta+1)}_\lambda J^{(\beta+1)}_\lambda, \hat{G}_J^{(\rho)} = c^{(\beta+1)}_\lambda J^{(\rho)}_\lambda \), where the constants \( c^{(\beta)}_\lambda \) and \( \hat{c}^{(\rho)}_\lambda \) are defined in Lemma 4.1.

Proof. Assume \( \mu <_D \lambda \). Then we have

\[
\langle G_J^{(\rho)} \rangle = \langle J^{(\rho)} \rangle = \langle J^{(\beta)} \rangle = 0,
\]

where we have used Lemma 4.1 and Theorem 4.2. From this fact, along with Lemma 4.1 (ii), we see that \( G_J^{(\beta+1)} \) coincides with \( c^{(\beta+1)}_\lambda J^{(\beta+1)}_\lambda \). The latter can be proved in similar way.

With these preliminaries, it is possible to derive the following result.

Proposition 4.5 ([M1])

\[
\langle J^{(\rho)} \rangle = N! \prod_{k=1}^{\beta} \frac{\lambda_i - \lambda_j - k + \beta(j - i + 1)}{\lambda_i - \lambda_j + k + \beta(j - i - 1)}.
\]

Proof. From Proposition 4.4, it follows that

\[
\langle J^{(\beta+1)} \rangle = \frac{1}{c^{(\beta+1)}_\lambda} \langle G_J^{(\beta+1)} \rangle = \frac{1}{c^{(\beta+1)}_\lambda} \langle J^{(\beta+1)} \rangle.
\]

Applying this relation repeatedly, we have

\[
\langle J^{(\beta)} \rangle = \prod_{k=0}^{\beta-1} \frac{c^{(\beta-k)}_\lambda}{c^{(\beta-k)}_{\lambda+\beta}} \langle J^{(\rho)} \rangle,
\]

which gives the desired result.

The norm formula (4.5) can be rewritten into the following form:

\[
\langle J^{(\rho)} \rangle = \frac{(N\beta)!}{(\beta)!^N} \prod_{(i,j) \in \lambda} \frac{j - 1 + \beta(N - i + 1)}{j + \beta(N - i)} \cdot \frac{\lambda_i - j + 1 + \beta(\lambda_j - i)}{\lambda_i - j + \beta(\lambda_j - i + 1)}.
\]

A proof of the equivalence between (4.5) and (4.6) is given in Appendix B.

4.2 Shift operators for the multivariable Hermite polynomials

In this section, we will construct shift operators for the multivariable Hermite polynomials. It should be noted that Heckman has constructed shift operators for the Hamiltonian \( H_A \) without harmonic potential [He2]. However, for the application to norm formulas, it is needed to compute actions of the shift shift operators on polynomials explicitly. Our method gives an unified and straightforward way to compute such actions.

To construct shift operators, we first introduce \( \mathcal{Y}_H \) and \( \hat{\mathcal{Y}}_H \) as follows:

\[
\mathcal{Y}_H = \rho^A(\mathcal{Y}_J) = \prod_{i<j}(\beta - \lambda_i^A + \lambda_j^A), \quad \hat{\mathcal{Y}}_H = \rho^A(\hat{\mathcal{Y}}_J) = \prod_{i<j}(\beta - \lambda_i^A + \lambda_j^A).
\]
Using these operators, we define shift operators for the multivariable Hermite polynomials as \( G_H = \mathcal{X}^{-1} \mathcal{Y}_H, \hat{G}_H = \hat{\mathcal{Y}}_H \mathcal{X} \). We stress that we have used same \( \mathcal{X} \) as (4.1), and therefore \( G_H \neq \rho^A(G_1), \hat{G}_H \neq \rho^A(\hat{G}_1) \). This reflects the characteristics of the scalar products (2.3), (2.6).

If we apply \( \rho^A \) to (4.2), we have

\[
\begin{aligned}
(s_j + 1)(-\beta - \tilde{h}_{j+1}^A + \tilde{h}_j^A) &= (-\beta - \tilde{h}_{j+1}^A + \tilde{h}_j^A)(s_j - 1), \\
(s_j - 1)(\beta - \tilde{h}_{j+1}^A + \tilde{h}_j^A) &= (\beta - \tilde{h}_{j+1}^A + \tilde{h}_j^A)(s_j + 1), \\
s_j(c - \tilde{h}_j^A + \tilde{h}_{j+1}^A)(c - \tilde{h}_j^A + \tilde{h}_{j+1}^A) &= (c - \tilde{h}_j^A + \tilde{h}_{j+1}^A)(c - \tilde{h}_j^A + \tilde{h}_{j+1}^A)s_j,
\end{aligned}
\]

with \( c \) arbitrary constant and \( k \neq j, j + 1 \). These relations imply that

\[
\begin{aligned}
\mathcal{Y}_H \left( \mathbb{C}[x]^\mathfrak{S}_N \right) &= \mathbb{C}[x]^{-\mathfrak{S}_N}, \\
\hat{\mathcal{Y}}_H \left( \mathbb{C}[x]^\mathfrak{S}_N \right) &= \mathbb{C}[x]^\mathfrak{S}_N.
\end{aligned}
\]  

(4.7)

Furthermore, if we apply \( \sigma^A \) to (4.4), we see that

\[
\begin{aligned}
\mathcal{Y}_H^H_{\lambda+\delta} &= c^{(\beta+1)}_\lambda x^{+\delta}_{\omega_0} + \text{ "lower terms" with respect to } \prec.
\end{aligned}
\]  

(4.8)

For the proof of the shift relations for the Jack polynomials (Proposition 4.4), Theorem 4.2 played a crucial role. Here we state analogous result for Hermite case:

**Theorem 4.6** For \( f, g \in \mathbb{C}[x]^\mathfrak{S}_N \), \( \langle G_H f, g \rangle_H^{(\beta+1)} = \langle f, \hat{G}_H g \rangle_H^{(\beta)} \).

**Proof.** The proof is similar to that of Theorem 4.2. For \( f, g \in \mathbb{C}[x]^\mathfrak{S}_N \), we know that

\[
\langle G_H f, g \rangle_H^{(\beta+1)} = \int_{-\infty}^{\infty} \cdots \int_{-\infty}^{\infty} \mathcal{P}_-(\mathcal{Y}_H f) g(\phi_A(\beta))^2 dx_1 \cdots dx_N.
\]

On the other hand, applying \( \rho^A \) to Lemma 4.3, we obtain

\[
\mathcal{P}_-(\mathcal{Y}_H - \hat{\mathcal{Y}}_H) = \sum_j \hat{g}_j(\tilde{a}_1, \ldots, \tilde{a}_N)(s_j - 1) \quad \text{for some } \hat{g}_j(x_1, \ldots, x_N) \in \mathbb{C}[x].
\]

From this relation, we find that \( \mathcal{P}_-(\mathcal{Y}_H - \hat{\mathcal{Y}}_H) f = 0 \) for all \( f \in \mathbb{C}[x]^\mathfrak{S}_N \). Hence we can replace \( \mathcal{Y}_H \) by \( \hat{\mathcal{Y}}_H \):

\[
\langle G_H f, g \rangle_H^{(\beta+1)} = \int_{-\infty}^{\infty} \cdots \int_{-\infty}^{\infty} \mathcal{P}_-(\hat{\mathcal{Y}}_H f) g(\phi_A(\beta))^2 dx_1 \cdots dx_N
\]

\[
= \langle \hat{\mathcal{Y}}_H f, \mathcal{X} g \rangle_H^{(\beta)} = \langle f, \hat{G}_H g \rangle_H^{(\beta)}
\]

In the last equality, we have used the self-adjointness of the operator \( \tilde{h}_j^A \).

\[\square\]

Now we are in position to state that:

**Proposition 4.7** \( G_H H_{\lambda+\delta}^{(\beta)} = c^{(\beta+1)}_\lambda H_{\lambda}^{(\beta+1)}, \hat{G}_H H_{\lambda}^{(\beta+1)} = \bar{c}^{(\beta+1)}_\lambda H_{\lambda+\delta}^{(\beta)}, \) where the constants \( c^{(\beta)}_\lambda \) and \( \bar{c}^{(\beta)}_\lambda \) are defined in Lemma 4.1.

**Proof.** From (4.7) and (4.8), we know that \( (c^{(\beta+1)}_\lambda)^{-1} G_H H_{\lambda+\delta}^{(\beta)} \) satisfies the first condition of Definition 2.5. So it suffice to prove the orthogonality which can be shown in the same way as Proposition 4.4. The second equation can be proved in similar way.

\[\square\]

Using Proposition 4.7 and Theorem 4.6, we can prove the norm formula for \( H_{\lambda}^{(\beta)} \).
Proposition 4.8 ([BF1, vD])

\[
\langle H^{(\beta)}_\lambda, H^{(\beta)}_\lambda \rangle_H \frac{\pi^{N/2}N!}{2^{[\lambda]_+ + \beta N(N-1)/2}} \times \prod_{i,j=1}^N (\lambda_j + \beta(N-j))! \prod_{k=1}^\beta \prod_{i,j} \frac{\lambda_i - \lambda_j - k + \beta(j-i+1)}{\lambda_i - \lambda_j + k + \beta(j-i-1)},
\]

where \(|\lambda| = \sum \lambda_j.

Proof. Using Proposition 4.7 and Theorem 4.6, we see that

\[
\langle H^{(\beta+1)}_\lambda, H^{(\beta+1)}_\lambda \rangle_H \frac{c^{(\beta+1)}_\lambda}{c^{(\beta)}_\lambda} \langle H^{(\beta)}_\lambda, H^{(\beta)}_\lambda \rangle_H ^{\lambda+1}.\rho+1).
\]

On the other hand, since \(H^{(\beta=0)}_\lambda(x)\) is direct product of the (one-variable) Hermite polynomials, one can evaluate the norm easily:

\[
\langle H^{(\beta=0)}_\lambda, H^{(\beta=0)}_\lambda \rangle_H = \frac{\pi^{N/2}}{2^{[\lambda]_+}} \prod_{j=1}^N \lambda_j!.
\]

Using these relations, one arrives at the formula above. \(\square\)

The norm formula (4.9) can be rewritten into the following form [BF1]:

\[
\langle H^{(\beta)}_\lambda, H^{(\beta)}_\lambda \rangle_H = \frac{\pi^{N/2}}{2^{[\lambda]_+ + \beta N(N-1)/2}} \prod_{j=1}^N (j\beta)! \prod_{i,j} \frac{\{j-1 + \beta(N-j+1)\} \{\lambda_i - j + 1 + \beta(\lambda'_i - i)\}}{\lambda_i - j + \beta(\lambda'_i - i+1)}.
\]

It should be remarked that other proofs of these formulas have been given via limiting procedure [BF1, vD].

4.3 Shift operators for the multivariable Laguerre polynomials

We first define \(X_L, Y_L\) and \(\hat{Y}_L\) as follows:

\[
X_L = \prod_{i<j} (z_i^2 - z_j^2), \quad Y_L = \prod_{i<j} (\beta - \tilde{h}_i^2, 2 + \tilde{h}_j^2/2), \quad \hat{Y}_L = \prod_{i<j} (-\beta - \tilde{h}_i^2, 2 + \tilde{h}_j^2/2).
\]

After same discussion as in the previous subsection, we see that

\[
Y_L \left(C[z^2]^{\mathcal{E}_N}\right) = C[z^2]^{-\mathcal{E}_N}, \quad \hat{Y}_L \left(C[z^2]^{-\mathcal{E}_N}\right) = C[z^2]^{\mathcal{E}_N},
\]

and

\[
Y^{L(\beta)}_H = c^{(\beta+1)}_{\lambda+\delta} z_{\mathcal{E}_N} \quad \text{"lower terms" with respect to } \prec.
\]

Now we define the shift operators for Laguerre case as \(G_L = X_L^{-1} Y_L, \quad \hat{G}_L = \hat{Y}_L X_L\).

These operators enjoy the following properties:

Theorem 4.9 For \(f, g \in C[z^2]^{\mathcal{E}_N}, \) \(\langle G_L f, g \rangle_L^{(\beta+1)} = \langle f, \hat{G}_L g \rangle_L^{(\beta)}\).
Proof. The proof for this case is also similar to that of Theorem 4.2. For \( f, g \in \mathbb{C}[z^2]^{\mathfrak{S}_N} \), we know that
\[
\langle G_L f, g \rangle_L^{(\beta+1)} = \int_{-\infty}^{\infty} \cdots \int_{-\infty}^{\infty} \mathcal{P}_-(\mathcal{Y}_L f) g \mathcal{X}_L (\phi^{(\beta)}_B) d\mathcal{Z}_1 \cdots d\mathcal{Z}_N.
\]
On the other hand, applying \( \rho^B \) to Lemma 4.3, we find that
\[
\mathcal{P}_-(\mathcal{Y}_L - \hat{\mathcal{Y}}_L) f = \sum_j \hat{g}_j(x_1, \ldots, x_N) \in \mathbb{C}[x].
\]
From this relation, we see that \( \mathcal{P}_-(\mathcal{Y}_L - \hat{\mathcal{Y}}_L) f = 0 \) for all \( f \in \mathbb{C}[z^2]^{\mathfrak{S}_N} \). Hence we can replace \( \mathcal{Y}_L \) by \( \hat{\mathcal{Y}}_L \):
\[
\langle G_L f, g \rangle_L^{(\beta+1)} = \int_{-\infty}^{\infty} \cdots \int_{-\infty}^{\infty} \mathcal{P}_-(\hat{\mathcal{Y}}_L f) g \mathcal{X}_L (\phi^{(\beta)}_B) d\mathcal{Z}_1 \cdots d\mathcal{Z}_N.
\]
In the last equality, we have used the self-adjointness of the operator \( \mathcal{H}_j^B \).

We then state the following results:

**Proposition 4.10** \( G_L L^{(\beta)}_{\lambda+} = c^{(\beta+1)}_\lambda L^{(\beta+1)}_\lambda \), \( \hat{G}_L L^{(\beta+1)}_\lambda = \hat{c}^{(\beta+1)}_\lambda L^{(\beta)}_\lambda \), where the constants \( c^{(\beta)}_\lambda \) and \( \hat{c}^{(\beta)}_\lambda \) are defined in Lemma 4.1.

**Proof.** From (4.11) and (4.12), we know that \( (c^{(\beta+1)}_\lambda)^{-1} G_L L^{(\beta)}_{\lambda+} \) satisfies the first condition of Definition 2.10 up to a constant factor. So it suffice to prove the orthogonality which can be shown in the same way as Proposition 4.4. The second equation can be proved in similar way.

Using Proposition 4.10 and Theorem 4.9, we can prove the norm formula for \( L^{(\beta)}_\lambda \).

**Proposition 4.11 ([BF1, vD])**
\[
\langle L^{(\beta)}_\lambda, L^{(\beta)}_\lambda \rangle_H^{(\beta)} = N! \prod_{j=1}^{N} (\lambda_j + \beta(N - j))!
\times \prod_{j=1}^{N} \Gamma(\lambda_j + \beta(N - j) + \gamma + 1/2) \prod_{k=1 \leq j}^{\beta} \prod_{l=1 \leq i}^{\beta} \frac{\lambda_i - \lambda_j - k + \beta(j - i + 1)}{\lambda_i - \lambda_j + k + \beta(j - i - 1)}
\]
where \( \Gamma(\cdot) \) denotes the gamma function.

**Proof.** The proof of this proposition is similar to the Hermite case. We only note the following formula for the case \( \beta = 0 \):
\[
\langle L^{(\beta=0)}_\lambda, L^{(\beta=0)}_\lambda \rangle_H^{(\beta=0)} = \#\mathfrak{S}_N \prod_{j=1}^{N} \{\lambda_j! \cdot \Gamma(\lambda_j + \gamma + 1/2)\},
\]
which follows from the norm formula of the one-variable Laguerre polynomials. \( \square \)
The norm formula (4.13) can be rewritten into the following form [BF1]:

\[ \langle L_{\lambda}^{(\rho)}, L_{\lambda}^{(\rho)} \rangle_{H} = \frac{\prod_{j=1}^{N} (j \beta)!}{(\beta!)^{N}} \prod_{j=1}^{N} \Gamma(\lambda_j + \beta(N - j) + \gamma + 1/2) \times \prod_{(i,j) \in \lambda} \frac{(j - 1 + \beta(N - i + 1))(\lambda_i - j + 1 + \beta(\lambda_j' - i))}{\lambda_i - j + \beta(\lambda_j' - i + 1)}. \]

It should be remarked that other proofs of these formulas have been given via limiting procedure [BF1, vD].

5 Concluding remarks

In this paper, we have reviewed the common algebraic structure of Calogero-Sutherland-type models, i.e. the degenerate double affine Hecke algebra. From this viewpoint, we can construct the intertwining operators that map the Jack polynomials to the multivariable Hermite and Laguerre polynomials.

We restrict ourselves to symmetric polynomials though the operators \( \sigma^A \) and \( \sigma^B \) are applicable to non-symmetric case, i.e. we can obtain the non-symmetric counterparts of the multivariable Hermite and Laguerre polynomials:

\[ E_{w}^{(H)}(x) = 2^{-|\lambda|/2} E_{w}^{\lambda}(\hat{g}^T) \cdot 1, \quad E_{w}^{(L)}(x) = E_{w}^{\lambda}((\hat{g}^T)^2/2) \cdot 1. \]

Baker and Forrester named these polynomials non-symmetric Hermite and Laguerre polynomials respectively, and studied their properties [BF2, BF3]. We note that some of their results may be obtained directly from the corresponding properties of the Jack polynomials by applying the intertwiners.

Our constructs are based on the degenerate double affine Hecke algebra, so it is expected that the results given here extend to non-degenerate case. Baker and Forrester studied isomorphism between affine Hecke algebras that maps the Macdonald polynomials to the multivariable Al-Salam and Carlitz polynomials [BF4]. On the other hand, van Diejen proposed difference counterpart of the Hamiltonians \( \mathcal{H}_A \) and \( \mathcal{H}_B \) [vD]. It would be nice to clarify algebraic structure of his models. We hope to report on them in the near future.

Appendices

Appendix A: Proof of Lemma 4.3

In Appendix A, we will give a proof of Lemma 4.3. We remark again that the proof given in this section is a limiting case of [Ki].

We begin with seeing some properties of the anti-symmetrizer.

Lemma A.1 ([Ki])

(i) The anti-symmetrizer \( \mathcal{P}_- \) is divisible by \( 1 + (-1)^{l(w_0)}w_0 \) both on the left and on the right.

(ii) For all \( j = 1, \ldots, N - 1 \), the anti-symmetrizer \( \mathcal{P}_- \) is divisible by \( 1 - s_j \) both on the left and on the right.
Proof. (i) \( \mathfrak{S}_N \) can be divided into pairs \((w, w_0 w)\). Then, rewriting into the summation over such pairs, we have

\[
P_-= \sum_{(w, w_0 w)} \{-(-1)^{(w)} w + (-1)^{(w_0 w)} w_0 w\} = \sum_{(w, w_0 w)} (-1)^{(w)} w \{1 + (-1)^{(w_0)} w_0\}.
\]

Divisibility on the left is proved similarly.

(ii) Can also be proved by similar discussion. \(\square\)

From Lemma A.1 (ii), we know that \( \text{Ker} \ P_- \supset \sum_j \text{Ker} (1-\mathfrak{s}_j) \).

To describe kernel of the anti-symmetrizer, we first investigate kernels of \(1-\mathfrak{s}_j\) and their union.

Lemma A.2 ([Ki])

(i) Let \( V \) be a representation of \( \mathfrak{S}_N \), and denote \( V_j = \text{Ker} (1-\mathfrak{s}_j) \), \( V' = \sum_j V_j \). Then \( V' \) is \( \mathfrak{S}_N \)-invariant.

(ii) Assume \( V \) is a finite-dimensional irreducible representation of \( \mathfrak{S}_N \). Then we have

\[
V' = \begin{cases} 
0 & \text{if } V \text{ is the sign representation}, \\
V & \text{(otherwise)}. 
\end{cases}
\]

Proof. (i) From the definition of \( V_j \), it follows that \( \mathfrak{s}_j(s_i v) = s_i v \) for all \( v \in s_i V_j \). If we introduce \( v_\pm = (v \pm s_i v)/2 \), we see that \( s_j(v_+ - v_-) = v_+ - v_- \) which means \( v_+ - v_- \in V_j \). Since \( v_+ \in V_i \) by definition, we obtain \( v = v_+ + v_- \in V_i + V_j \). This leads to \( s_i V_j \subseteq V_i + V_j \), which concludes the proof.

(ii) From (i), it follows that \( V' \) is a subrepresentation. Due to the irreducibility, \( V' \) can be either \( 0 \) or \( V \). If \( V' = 0 \), then we have \( V_j = 0 \) for all \( j \). This means that \( 1-\mathfrak{s}_j \) is invertible, i.e. for all \( v \in V \), there exists \( u \) such that \( v = (1-\mathfrak{s}_j)u \). Then we obtain \( s_j v = -v \) for all \( v \in V \), i.e. \( V \) is the sign representation. \(\square\)

From Lemma A.2(ii), it immediately follows:

\[
\text{Ker} \ P_- = \sum_j \text{Ker} \ (1-\mathfrak{s}_j)
\]

for any finite-dimensional representation of \( \mathfrak{S}_N \). Note this identity also holds for the representation of \( \mathfrak{S}_N \) in the space of polynomials \( \mathbb{C}[x] \), since this representation is a direct sum of finite-dimensional representations.

We now introduce operators \( \hat{s}_j \) as

\[
\hat{s}_j = s_j + \beta \frac{s_j - 1}{x_j - x_{j+1}}.
\]

Using these operators, we can define another representation of the degenerate affine Hecke algebra \( \mathfrak{H}_0' \) on \( \mathbb{C}[x] \):

\[
\rho'(\hat{D}_j^A) = x_j, \quad \rho'(\hat{s}_j) = \hat{s}_j.
\]

Using the isomorphism \( \rho' \), we introduce deformed anti-symmetrizer \( P_-^{(\beta)} \) as

\[
P_-^{(\beta)} = \rho'(P_-) = \frac{1}{\#\mathfrak{S}_N} \sum_{w \in \mathfrak{S}_N} (-1)^{(\hat{w})} \hat{w},
\]

where \( \hat{w} = \rho'(w) \).
Lemma A.3 ([Ki]) Ker $\mathcal{P}_-(\beta) = \text{Ker } P$ for the action of $\mathcal{P}_-(\beta)$ in $\mathbb{C}[x]$:

Proof. By similar discussion to Proposition A.1, we know that $\mathcal{P}_-(\beta)$ is divisible by $s_j - 1$ both on the left and on the right for every $j = 1, \ldots, N - 1$. Hence we have

$$\text{Ker } \mathcal{P}_-(\beta) \supset \sum_j \text{Ker } (1 - s_j) = \text{Ker } P,$$

for all $f \in \mathbb{C}[x]^\mathfrak{S}_N$. On the other hand, if we denote $\mathbb{C}[x]_n$ as space of polynomials of order $n$, it is clear that $\mathcal{P}_-(\beta)$ preserves $\mathbb{C}[x]_n$. Since $\dim(\text{Ker } \mathcal{P}_-(\beta))$ cannot decrease under specialization, it follows that $\dim(\text{Ker } \mathcal{P}_-(\beta)) = \dim(\text{Ker } P)$ and hence we have

$$\dim(\text{Ker } \mathcal{P}_-(\beta)) = \dim(\text{Ker } P).$$

Thus it follows from (A.1) and (A.2) that $\text{Ker } \mathcal{P}_-(\beta) = \sum_j \text{Ker } (1 - s_j).$ □

We then define $\mathcal{Y}'$ and $\hat{\mathcal{Y}}'$ as

$$\mathcal{Y}' = \rho'(\mathcal{Y}_1) = \prod_{i<j}(\beta - x_i + x_j), \quad \hat{\mathcal{Y}}' = \rho'(\hat{\mathcal{Y}}_1) = \prod_{i<j}(-\beta - x_i + x_j).$$

Lemma A.4 ([Ki])

$$\mathcal{P}_-(\mathcal{Y}' - \hat{\mathcal{Y}}')f = 0 \quad \text{for all } f \in \mathbb{C}[x]^{\mathfrak{S}_N}.$$

Proof. We can show that

$$(1 + (-1)^{N(N-1)/2}w_0)(\mathcal{Y}' - \hat{\mathcal{Y}}') = (\mathcal{Y}' - \hat{\mathcal{Y}}')(1 - w_0).$$

by the direct calculation. Considering the action on $\mathbb{C}[x]^{\mathfrak{S}_N}$, we have

$$(1 + (-1)^{N(N-1)/2}w_0)(\mathcal{Y}' - \hat{\mathcal{Y}}')f = 0,$$

for all $f \in \mathbb{C}[x]^{\mathfrak{S}_N}$. Using this formula and Proposition A.1 (i), we obtain the desirous result. □

From Lemma A.3 and Lemma A.4, we know that

$$\mathcal{P}_-(\beta)(\mathcal{Y}' - \hat{\mathcal{Y}}')f = 0 \quad \text{for all } f \in \mathbb{C}[x]^{\mathfrak{S}_N}.$$

On the other hand, the following statement can easily be proved:

Lemma A.5 Let $\hat{A}$ be a operator of the form, $\hat{A} = \sum_{w \in \mathfrak{S}_N} g_w \hat{w}$ with $g_w \in \mathbb{C}[x]$. If $\hat{A}f = 0$ for all $f \in \mathbb{C}[x]^{\mathfrak{S}_N}$, then $\hat{A}$ can be represented in the following form:

$$\hat{A} = \sum_i \hat{g}_j(s_j - 1) \quad \text{for some } \hat{g}_j \in \mathbb{C}[x]^{\mathfrak{S}_N}.$$
Proof. The operator $\hat{A}$ can be rewritten as
$$\hat{A} = \sum_{j_1, \ldots, j_k} \hat{g}'_{j_1, \ldots, j_k} (\hat{s}_{j_1} - 1) \cdots (\hat{s}_{j_k} - 1) + \hat{g}_0'$$
for some $\hat{g}'_{j_1, \ldots, j_k} \in \mathbb{C}[x]$.
Then the assumption of the proposition means $\hat{g}_0' = 0$, which gives the desirous result. \qed

Applying Lemma A.5 to Lemma A.4, we conclude that
$$P_{-}^{(\rho')} (Y - \tilde{Y}') = \sum_j \hat{g}_j(x_1, \ldots, x_N)(\hat{s}_j - 1)$$
for some $\hat{g}_j(x_1, \ldots, x_N) \in \mathbb{C}[x]$. (A.3)

Applying $(\rho')^{-1}$ completes the proof of Lemma 4.3.

Appendix B: Equivalence of the two expressions for the norm formula

In Appendix B, we will give a proof of equivalence between two expressions of the norm formulas. We first begin with considering the Jack case.

Let $\lambda$ be a partitions satisfying the following conditions (see Figure 1 below):
$$\lambda_{p-1} > \lambda_p = \cdots = \lambda_{p+r_1} = \cdots = \lambda_{p+r_1 + r_2 - 2}$$
$$> \cdots > \lambda_{p+r_1 + \cdots + r_m - 1} = \cdots = \lambda_{p+r_1 + \cdots + r_m - 1} > \lambda_{p+r_1 + \cdots + r_m} = \cdots = 0,$n
$$\lambda_1' = \cdots = \lambda_{s_1}' = \cdots = \lambda_{s_1+s_2}'$$
$$> \cdots > \lambda_{s_1+s_2}' = \cdots = \lambda_{s_1+s_2+s_3}' = \cdots = \lambda_{s_1+s_2+s_3+s_m}' = \cdots = 0.$n

We further define $\mu$ as $\mu = (\lambda_1, \cdots, \lambda_p+1, \cdots, \lambda_N)$.

Calculating the ratio $\langle J^{(\beta)}_{\mu}, J^{(\beta)}_{\mu} \rangle_{J} / \langle J^{(\beta)}_{\lambda}, J^{(\beta)}_{\lambda} \rangle_{J}$ by using (4.5) or (4.6), one can show that both cases reduce to
$$\frac{\langle J^{(\beta)}_{\mu}, J^{(\beta)}_{\mu} \rangle_{J}}{\langle J^{(\beta)}_{\lambda}, J^{(\beta)}_{\lambda} \rangle_{J}} = \prod_{i=1}^{p-1} \frac{\lambda_i - \lambda_p + \beta(p-i)}{\lambda_i - \lambda_p - 1 + \beta(p-i)} \cdot \frac{\lambda_i - \lambda_p - 1 + \beta(p-i)}{\lambda_i - \lambda_p + \beta(p-i - 1)} \cdot \frac{\lambda_i - \lambda_p - 1 + \beta(p-i - 1)}{\lambda_i - \lambda_p - 1 + \beta(p-i - 1)}.$$
\[
\begin{align*}
& \times \frac{s_m + 1 + \beta(r_1 - 1)}{1 + \beta(r_1 - 1)} \cdots \frac{s_m + \cdots + s_1 + 1 + \beta(r_1 + \cdots + r_m - 1)}{s_m + \cdots + s_2 + 1 + \beta(r_1 + \cdots + r_m - 1)} \\
& \times \frac{\beta r_1}{s_m + \beta r_1} \cdot \frac{s_m + \beta(r_1 + r_2)}{s_m + s_m-1 + \beta(r_1 + r_2)} \cdots \frac{s_m + \cdots + s_1 + \beta(r_1 + \cdots + r_m)}{s_m + \cdots + s_1 + 1 + \beta(N - p + 1) + 1} \\
& \times \frac{\beta}{s_m + \cdots + s_1 + 1 + \beta(N - p)}
\end{align*}
\]

On the other hand, if we consider the simplest case \( \lambda = \phi \), both (4.5) and (4.6) reduce to \((1, 1)^{[p]} = (\beta N)!/(\beta!)^N \). Hence, by induction, we conclude that (4.5) and (4.6) are equivalent for all \( \lambda \).

The Hermite and Laguerre cases can be proved in the similar fashion.

References


