A Study of the Relativistic Euler Equation

Tetu Makino (牧野 哲)
Faculty of Engineering,
Yamaguchi University (山口大工)

This is a joint work with Cheng-Hsiung Hsu (National Central University, Chungli, Taiwan) and Song-Sun Lin (National Chiao Tung University, Hsinchu, Taiwan).

1 Introduction

In this article we study the Cauchy problem to the one-dimensional relativistic Euler equation

\[ \frac{\partial}{\partial t} \frac{\rho + Pu^2/c^4}{1 - u^2/c^2} + \frac{\partial}{\partial x} \frac{(\rho + P/c^2)u}{1 - u^2/c^2} = 0, \]
\[ \frac{\partial}{\partial t} \frac{P + \rho u^2}{1 - u^2/c^2} + \frac{\partial}{\partial x} \frac{(\rho + P/c^2)u}{1 - u^2/c^2} = 0, \]

(1.1)

\[ \rho|_{t=0} = \rho_0(x), \quad u|_{t=0} = u_0(x). \]  

(1.2)

Here $c$ is a positive constant, the speed of light, and $P$ is a given function of $\rho$. The equation (1.1) governs the one dimensional motion of a perfect gas in the Minkowski space-time. When $c \to \infty$, (1.1) tends to the usual Euler equation of gas dynamics

\[ \rho_t + (\rho u)_x = 0, \]
\[ (\rho u)_t + (P + \rho u^2)_x = 0. \]

(1.3)

Many mathematical investigations for this non-relativistic Euler equation were done. But the first mathematical investigation for the relativistic Euler equation (1.1) was done recently by Smoller and Temple [6]. They assume $P = \sigma^2 \rho$, where $\sigma$ is a positive constant $< c$. Under this assumption, they showed that if the initial data $\rho_0(x)$ and $u_0(x)$ satisfy

\[ T.V. \log \rho_0 < \infty, \quad T.V. \log \frac{c + u_0}{c - u_0} < \infty, \]
then there exists a global weak solution to the Cauchy problem (1.1)(1.2). The result was obtained by Glimm’s scheme and it is the relativistic version of Nishida’s result [5] for the non-relativistic problem.

However we would like to consider a more realistic equation of states. We keep in mind the equation of state for a neutron stars, which is given by

$$P = Kc^5 f(y), \quad \rho = Kc^3 g(y)$$

$$f(y) = \int_0^y \frac{q^4}{\sqrt{1 + q^2}} dq,$$

$$g(y) = 3 \int_0^y q^2 \sqrt{1 + q^2} dq.$$  

For this equation of state, we have $P \sim \frac{c^2}{3} \rho$ as $\rho \to \infty$ but $P \sim \frac{1}{5} K^{2/3} \rho^{5/3}$ as $\rho \to 0$. So we assume the following properties of the function $P(\rho)$:

\begin{itemize}
  \item[(A):] $P(\rho) > 0$, \quad $0 < dP/d\rho < c^2$, \quad $0 < d^2 P/d\rho^2$
\end{itemize}

for $\rho > 0$, and

$$P = A \rho^\gamma (1 + \left[\rho^{\gamma-1} / c^2\right]_1)$$

as $\rho \to 0$. Here $A$ and $\gamma$ are positive constants and

$$\gamma = 1 + \frac{2}{2N + 1},$$

$N$ being a positive integer, and $[X]_1$ denotes a convergent power series of the form \( \sum_{k \geq 1} a_k X^k \).

The result which we want to generalize to the relativistic problem is those by G.-Q. Chen et al [2]. So we assume that the initial data $\rho_0(x), u_0(x)$ satisfy

$$0 \leq \rho_0(x) \leq M_0, \quad \left| \frac{c}{2} \log \frac{c + u_0(x)}{c - u_0(x)} \right| \leq M_0.$$  

A weak solution of (1.1)(1.2) is defined as follows.

We write

$$E = \frac{\rho + Pu^2/c^4}{1 - u^2/c^2},$$

$$F = \frac{(\rho + P/c^2)u}{1 - u^2/c^2},$$

$$G = \frac{P + pu^2}{1 - u^2/c^2},$$

$$U = (E, F)^T, \quad f(U) = (F, G)^T.$$  

Then (1.1) can be written as

$$U_t + f(U)_x = 0.$$
Let us denote by $U_0(x)$ the initial data. Then a weak solution $U(t, x)$ is a bounded measurable function which satisfies
\[
\int \int (U \Phi_t + f(U) \Phi_x) dx dt + \int U_0(x) \Phi(0, x) dx = 0
\]
for any test function $\Phi \in C_0^\infty([0, +\infty) \times \mathbb{R})$.

2 Riemann problems

The Riemann problem is the problem to the special initial data of the form
\[
U_0(x) = \begin{cases} 
U_L & \text{if } x < 0 \\
U_R & \text{if } x > 0 
\end{cases}
\]
In order to solve this we introduce the Riemann invariants
\[
w = x + y, \quad z = x - y
\]
where
\[
x = \frac{c}{2} \log \frac{c+u}{c-u}, \quad y = \int_0^\rho \frac{\sqrt{P'}}{\rho + P/c^2} d\rho.
\]
Then (1.1) is diagonalized as
\[
w_t + \lambda_2 w_x = 0, \quad z_t + \lambda_1 z_x = 0,
\]
where
\[
\lambda_1 = \frac{u - \sqrt{P'}}{1 - \sqrt{P'u/c^2}}, \quad \lambda_2 = \frac{u + \sqrt{P'}}{1 + \sqrt{P'u/c^2}}.
\]
the possible states $U = U_R$ connected to $U_L$ on the right by rarefaction waves are
\[
R_1: \quad w = w_L, z > z_L
\]
and
\[
R_2: \quad w > w_L, z = z_L.
\]
The Rankine Hugoniot jump condition
\[
\sigma[U] = [f(U)],
\]
where $[U] = U_R - U_L$, $[f(U)] = f(U_R) - f(U_L)$, gives the shock curve
\[
\frac{(u_R - u_L)^2}{(1 - u_R^2/c^2)(1 - u_L^2/c^2)} = \frac{(\rho_R - \rho_L)(P_R - P_L)}{(\rho_L + P_L/c^2)(\rho_R + P_R/c^2)}.
\]
Along this curve we have shocks
\[
S_1: \quad \rho_L < \rho_R, u_R < u_L,
\]
\[
S_2: \quad \rho_R < \rho_L, u_R < u_L.
\]
The Riemann problem can be solved uniquely by using these rarefaction waves and shock waves and vacuum state. The detailed discussion can be found in J. Chen [1].

If we look at a region of the form

$$\Sigma_B = \{(w, z) | -B \leq z \leq w \leq B\},$$

we have the following

**Proposition 1** If the initial data $U_L, U_R$ belong to $\Sigma_B$ for some large $B$, then the solution of the Riemann problem is confined to $\Sigma_B$.

Moreover if we consider the image of $\Sigma_B$ in the $(E, F)$-space, we have

**Proposition 2** The region $\Sigma_B$ is convex in the $(E, F)$-plane.

Proof. Let us consider the above hedge $F = F(E)$ which corresponds to $w = B, -B < z < B$. We have to show $d^2 F/dE^2 < 0$. Along the hedge $w = B$, we have

$$u = c \tanh \left( \frac{1}{c} \left( B - \int_0^\rho \frac{\sqrt{P'}}{\rho + P/c^2} d\rho \right) \right),$$

from which

$$\frac{du}{d\rho} = -(1-u^2/c^2) \frac{\sqrt{P'}}{\rho + P/c^2}.$$

By a direct calculation we have

$$\frac{dF}{dE} = \frac{u - \sqrt{P'}}{1 - \sqrt{P'u/c^2}} = \lambda_1.$$

Differentiating once more we have

$$\frac{d^2 F}{dE^2} = -\frac{1-u^2/c^2}{(1-\sqrt{P'u/c^2})^4} \left( \frac{P''}{2\sqrt{P'}} + \frac{P'}{c^2} \frac{\sqrt{P'}}{\rho + P/c^2} \right) < 0.$$

This was to be seen. QED.

From Proposition 2, we have

**Proposition 3** If $U(s), s \in [a, b]$, is confined to a region $\Sigma_B$, then the average

$$\frac{1}{b-a} \int_a^b U(s) ds$$

belongs to $\Sigma_B$.

Let us look at the shock wave which connects the left state $U_L$ to the right state $U_R$ with the shock speed $\sigma$.

The right state $U_R$ and $\sigma$ are parametrized by $\rho = \rho_R$. Then we have the following fact, which will be used in Section 4.
Proposition 4 Along $S_1(\rho_L < \rho)$, we have $d\sigma/d\rho < 0$, and along $S_2(\rho < \rho_L)$ we have $d\sigma/d\rho > 0$.

Proof. Without loss of generality we can assume $u_L = 0$. Then $u = u_R$ is given by

$$u = -\sqrt{\frac{[\rho][P]}{(\rho_L + P/c^2)(\rho + P_L/c^2)}},$$

where $[\rho] = \rho - \rho_L, [P] = P - P_L$. We have

$$\sigma = \frac{[F]}{[E]} = \frac{(\rho + P/c^2)u}{\rho + Pu^2/c^4 - \rho L(1-u^2/c^2)}.$$

By a direct but tedious computations, we have

$$\frac{d\sigma}{d\rho} = \frac{(\rho + P/c^2)(\rho_L + P_L/c^2)[\rho]X}{2(\rho + Pu^2/c^4 - \rho L(1-u^2/c^2))^2u(\rho L + P/c^2)(\rho + P_L/c^2)^2},$$

$$X = (\rho + P_L/c^2)(\rho + P/c^2)[P]' + (\rho + P_L/c^2)(-2\rho + P + P_L/c^2)\frac{[P]}{c^2} - (\rho L + P/c^2)[P]^2/c^2.$$

Since $P'' > 0$ we know $[P] \leq P'[\rho]$. Thus

$$X \geq (\rho + P_L/c^2)(\rho + P/c^2)[P] + (\rho + P_L/c^2)(-2\rho + P + P_L/c^2)\frac{[P]}{c^2} - (\rho L + P/c^2)[P]^2/c^2.$$

But

$$1 > \frac{[\rho] - [P]/c^2}{[\rho]} = 1 - P'(\rho_L + \theta(\rho - \rho_L))/c^2 > 0.$$

Using this, it's easy to see $X > 0$ both when $[\rho] > 0$ and when $[\rho] < 0$. Since $u < 0$, this completes the proof. QED.

3 Entropies

A pair of functions $\eta$ and $q$ is called an entropy-entropy flux if it satisfies the equation

$$D_U q = D_U \eta. D_U f.$$

Using the Riemann invariants, we can write (3.1) as

$$\frac{\partial q}{\partial w} = \lambda_2 \frac{\partial \eta}{\partial w}, \quad \frac{\partial q}{\partial z} = \lambda_1 \frac{\partial \eta}{\partial z}.$$
By eliminating $q$ from the equation, we get the following second order equation:
\[
\frac{\partial^2 \eta}{\partial w \partial z} + Q \left( J \frac{\partial \eta}{\partial w} - \frac{1}{J} \frac{\partial \eta}{\partial z} \right) = 0,
\]
\[(3.2)\]
where
\[
Q = \frac{1}{4\sqrt{P'}} \left( 1 - \frac{P'}{c^2} - \frac{\rho + P/c^2}{2P} \right),
\]
\[
J = \frac{1 - \sqrt{P'}u/c^2}{1 + \sqrt{P'}u/c^2}.
\]
Since this equation tends to the Euler-Poisson-Darboux equation
\[
\frac{\partial^2 \eta}{\partial w \partial z} + \frac{N}{w-z} \left( \frac{\partial \eta}{\partial w} - \frac{\partial \eta}{\partial z} \right) = 0
\]
\[(3.3)\]
as $c \to \infty$, we shall call (3.2) the relativistic Euler-Poisson-Darboux equation.

Among entropies of (3.3) when $c = \infty$ the kinetic energy
\[
\eta = \frac{1}{2} \rho u^2 + \frac{P}{\gamma - 1}
\]
\[(3.4)\]
plays an important role. Therefore we want to find an entropy of (3.2) which tends to (3.4) as $c \to \infty$. Let us look for an entropy-entropy flux of the form
\[
\eta = H(\rho, u^2), \quad q = Q(\rho, u^2)u.
\]
Inserting this to the equation it is easy to find an entropy-entropy flux
\[
\eta^* = -\frac{\Psi(\rho)}{(1-u^2/c^2)^{1/2}} + c^2 \left( \frac{\rho + Pu^2/c^4}{1-u^2/c^2} \right),
\]
\[
q^* = \left( \frac{\Psi(\rho)}{(1-u^2/c^2)^{1/2}} + c^2 \frac{\rho + P/c^2}{1-u^2/c^2} \right)u,
\]
\[
\Psi = \exp \left( \int_{1}^{\beta} \frac{d\rho}{\rho + P/c^2 + K_0} \right),
\]
\[(3.5),(3.6),(3.7)\]
where $K_0$ is determined so that $\eta^*$ tends to the kinetic energy (3.4) as $c = \infty$. We call the entropy $\eta^*$ defined by (3.5) the relativistic standard entropy. The important fact is

**Proposition 5** The Hessian $D^2_{U} \eta^*$ is positive definite. For any fixed $B$ there is a positive constant $k$ such that
\[
(\xi | D^2_{U} \eta^* (U) \xi ) \geq k|\xi|^2,
\]
for any $U \in \Sigma_B$ and $\xi = (\xi_0, \xi_1)$ with $|\xi|^2 = \xi_0^2 + \xi_1^2$. 
Proof. The proof is due to direct but tedious calculations. We note
\[
\begin{align*}
\frac{\partial \rho}{\partial E} &= \frac{1 + u^2/c^2}{1 - P'u^2/c^4}, \\
\frac{\partial u}{\partial E} &= \frac{(1 + P'/c^4)(1 - u^2/c^2)u}{(\rho + P/c^2)(1 - P'u^2/c^4)}, \\
\frac{\partial \rho}{\partial F} &= -\frac{2u/c^2}{1 - P'u^2/c^4}, \\
\frac{\partial u}{\partial F} &= \frac{(1 - u^2/c^2)(1 + P'u^2/c^4)}{(\rho + P/c^2)(1 - P'u^2/c^4)},
\end{align*}
\]
Using these, we have
\[
\begin{align*}
\frac{\partial \eta^*}{\partial E} &= -\frac{\Psi}{(\rho + P/c^2)(1 - u^2/c^2)^{1/2} + c^2}, \\
\frac{\partial \eta^*}{\partial F} &= -\frac{\Psi u/c^2}{(\rho + P/c^2)(1 - u^2/c^2)^{1/2}}, \\
\frac{\partial^2 \eta^*}{\partial E^2} &= \frac{\Psi}{c^2}, \\
\frac{\partial^2 \eta^*}{\partial E \partial F} &= \frac{\Psi u/c^2}{(\rho + P/c^2)(1 - u^2/c^2)^{1/2}}(2P'/c^2 + 1 + P'u^2/c^4), \\
\frac{\partial^2 \eta^*}{\partial F^2} &= \frac{\Psi}{c^2}.
\end{align*}
\]
Therefore we get
\[
\begin{align*}
(\xi | D_0^2 \eta^* \xi) &= \eta_{EE} \xi_0^2 + 2 \eta_{EF} \xi_0 \xi_1 + \eta_{FF} \xi_1^2 \\
&= \frac{\Psi}{c^2}, \\
Z &= (P' + 2P'u^2/c^2 + u^2)\xi_0^2 - 2(2P'/c^2 + 1 + P'u^2/c^4)u\xi_0 \xi_1 + \\
&+ (1 + 3P'u^2/c^4)\xi_1^2 \\
&\quad \geq \frac{2P'(1 - u^2/c^2)^2(1 - P'u^2/c^4)}{A + C + \sqrt{(A - C)^2 + 4B^2}}(\xi_0^2 + \xi_1^2), \\
A &= P' + 2P'u^2/c^2 + u^2, \\
B &= (2P'/c^2 + 1 + P'u^2/c^4)u, \\
C &= 1 + 3P'u^2/c^4.
\end{align*}
\]
This completes the proof. QED.

4 Construction of approximate solutions

Let us construct approximate solutions using the Godunov scheme. The construction is similar if we use the Lax-Friedrichs scheme.
Suppose that the initial data $U_0(x)$ is confined to an invariant region $\Sigma_B$. Put $\Lambda_0 = \sup\{\lambda_j(U)| j = 1, 2, U \in \Sigma_B\}$. Fixing $\Lambda_1 > \Lambda_0$, we take mesh lengths $\Delta x, \Delta t$ such that $\Delta x = \Lambda_1 \Delta t$. We denote $\Delta = \Delta x$.

Let us construct the approximate solution $U^\Delta(t, x)$. First we put

$$U_0^\Delta(x) = U_0(x)\chi_{[-1/\Delta, 1/\Delta]}.$$ 

We define

$$U^\Delta(+0, x) = \frac{1}{2\Delta x} \int_{2j\Delta x}^{(2j+2)\Delta x} U_0^\Delta(x)dx$$

for $2j\Delta x < x \leq (2j+2)\Delta x$. Solving the Riemann problem on each interval $[2(j-1)\Delta, 2(j+1)\Delta]$, we define $U^\Delta(t, x)$ for $0 \leq t < \Delta t$. Since the Courant-Friedrichs-Levi condition is satisfied, the wave from the center $2j\Delta$ does not intersect. If $U^\Delta(t, x)$ for $0 \leq t < n\Delta t$ has been defined, then we define

$$U^\Delta(n\Delta t, x) = \frac{1}{2\Delta} \int_{2j\Delta}^{(2j+2)\Delta} U^\Delta(n\Delta t - 0, x)dx$$

for $2j\Delta < x \leq (2j+2)\Delta$. Solving the Riemann problem, we define $U^\Delta(t, x)$ for $n\Delta t < t < (n+1)\Delta t$.

By Proposition 1 and 3, it is inductively guaranteed that $U^\Delta$ remains in $\Sigma_B$, say,

**Proposition 6** The approximate solution $U^\Delta(t, x)$ satisfies $U^\Delta(t, x) \in \Sigma_B$, therefore,

$$0 \leq \rho^\Delta(t, x) \leq M, \quad \frac{c}{2} \log \frac{c + u^\Delta(t, x)}{c - u^\Delta(t, x)} \leq M.$$

Moreover we shall prove

**Proposition 7** For any test function $\Phi$ it holds that

$$\int \int \Phi(tU^\Delta + \Phi_x f(U^\Delta))dxdt + \int \Phi(0, x)U_0^\Delta(x)dx = O(\Delta^{1/2}).$$

In order to prove Proposition 7, we prepare

**Proposition 8** For any shock wave from $U_L$ to $U_R$ with the shock speed $\sigma$ and for any convex entropy $\eta$, we have

$$\sigma[\eta] - [q] \geq 0,$$

where $[\eta] = \eta(U_R) - \eta(U_L), [q] = q(U_R) - q(U_L)$. 
Proof. The right state of shocks can be parametrized by $\rho = \rho_R$. Putting $Q(\rho) = \sigma[\eta] - [q]$, we shall see $dQ/d\rho \geq 0$ along $S_1 : [\rho] > 0$ and $dQ/d\rho \leq 0$ along $S_2 : [\rho] < 0$. Using the equation (3.1) and the differentiation of the Rankine-Hugoniot condition, we have

$$
\frac{dQ}{d\rho} = \frac{d\sigma}{d\rho}([\eta] - D_U\eta(U).[U])
= -\frac{d\sigma}{d\rho} \int_0^1 \theta(U - U_L|D^2_U\eta(U_L + \theta(U - U_L).(U - U_L))d\theta.
$$

We supposed $D^2_U\eta \geq 0$. By Proposition 4, we know $d\sigma/d\rho < 0$ on $S_1$ and $d\sigma/d\rho > 0$ on $S_2$. QED.

Proof of Proposition 7.

We fix $T$ to consider $U^\Delta$ on $0 \leq t \leq T$. First we shall show

$$
\sum_{j,n} \int_{2j\Delta}^{(2j+2)\Delta} |U(n\Delta t - 0, x) - U(n\Delta t + 0, (2j+1)\Delta)|^2 dx \leq C.
$$

(4.1)

Let us consider the standard entropy $\eta^*$. Then we have

$$
0 = \int \eta^*(U(T, x))dx - \int \eta^*(U(0, x))dx + L + \Sigma,
$$

$$
L = \sum_{j,n} \int_{2j\Delta}^{(2j+2)\Delta} (\eta^*(U(n\Delta t - 0, x)) - \eta^*(U(n\Delta t + 0, (2j+1)\Delta))]dx,
$$

$$
\Sigma = \int_0^T \sum_{shocks} \int \sigma[\eta^*] - [q^*])dt.
$$

We write $U_0 = U(n\Delta t + 0, (2j+1)\Delta), U_1 = U(n\Delta t - 0, x)$. Since

$$
U_0 = \frac{1}{2\Delta} \int_{2j\Delta}^{(2j+2)\Delta} U_1 dx,
$$

we see

$$
L = \sum_{j,n} \int_{2j\Delta}^{(2j+2)\Delta} \int_0^1 (1 - \theta)(U_1 - U_0|D^2_U\eta^*(U_0 + \theta(U_1 - U_0)).(U_1 - U_0))d\theta dx
\geq 0.
$$

On the other hand we have $\Sigma \geq 0$ from Proposition 8. Thus $L \leq C, \Sigma \leq C$.

But from Proposition 5, we have $D^2_U\eta \geq k$. Therefore

$$
C \geq L \geq \frac{k}{2} \sum_{j,n} \int_{2j\Delta}^{(2j+2)\Delta} |U_1 - U_0|^2 dx.
$$
Thus we get (4.1).

Now let us consider a test function $\Phi$. Put

$$ J = \int \int (\Phi_t U^\Delta + \Phi_x f(U^\Delta))dxdt + \int \Phi(0,x)U_0^\Delta dx. $$

Since $U^\Delta$ is a weak solution on each time strip $n\Delta t < t < (n+1)\Delta t$, we have

$$ J = \sum_{n} \int \Phi(n\Delta t,x)(U(n\Delta t - 0,x) - U(n\Delta t + 0,x))dx $$

$$ = J_1 + J_2, $$

$$ J_1 = \sum_{j,n} \int_{2j\Delta}^{(2j+2)\Delta} \Phi(n\Delta t,j\Delta)(U(n\Delta t - 0,x) - U(n\Delta t + 0,x))dx, $$

$$ J_2 = \sum_{j,n} \int_{2j\Delta}^{(2j+2)\Delta} (\Phi(t,x) - \Phi(n\Delta t,j\Delta))(U(n\Delta t - 0,x) - U(n\Delta t + 0,x))dx. $$

Since

$$ U(n\Delta t + 0,x) = \frac{1}{2\Delta} \int_{2j\Delta}^{(2j+2)\Delta} U(n\Delta t - 0,x)dx $$

for $2j\Delta < x < (2j+2)\Delta$, we see $J_1 = 0$. It follows from (4.1) that

$$ |J_2| \leq C\Delta^{1/2}||\Phi||C^1(\sum_{j,n} \int_{2j\Delta}^{(2j+2)\Delta} |U(n\Delta t - 0,x) - U(n\Delta t + 0,x)|^2dx)^{1/2} $$

$$ \leq C'\Delta^{1/2}. $$

Here we have used $T/\Delta t = O(1/\Delta)$. QED.

Summing up, we have the following theorem.

**Theorem 1** The approximate solution $U^\Delta(t,x)$ satisfies

$$ 0 \leq \rho^\Delta(t,x) \leq M, \quad \left| \frac{c}{2} \log \frac{c + u^\Delta(t,x)}{c - u^\Delta(t,x)} \right| \leq M $$

and

$$ \int \int (\Phi_t U^\Delta + \Phi_x f(U^\Delta))dxdt + \int \Phi(0,x)U_0^\Delta(x) = O(\Delta^{1/2}) $$

for any test function $\Phi$.

We expect that $U^\Delta$ tends to a weak solution everywhere. For the non-relativistic gas dynamics, this was done by DiPerna [3] and G.Q.Chen et al [2]. In their proof the Darboux formula

$$ \eta = \int_w^w ((w - s)(s - z))^N \phi(s)ds $$

which gives solutions of the Euler-Poisson-Darboux equation (3.3) , $\phi$ being arbitrary, plays an important role. Section 6 will be devoted to find such an integral formula for the relativistic Euler-Poisson-Darboux equation (3.2).
5 Remark

We note that
\[
\lambda_2 - \lambda_1 = \frac{\sqrt{P'}(1 - u^2/c^2)}{1 - u^2 P'/c^4} > 0,
\]
\[
\frac{\partial \lambda_1}{\partial z} = \frac{1 - u^2/c^2}{2(1 - \sqrt{P'}u/c^2)}(1 - \frac{\rho + P/c^2}{2P'}) > 0,
\]
\[
\frac{\partial \lambda_2}{\partial w} = \frac{1 - u^2/c^2}{2(1 + \sqrt{P'}u/c^2)}(1 - \frac{\rho + P/c^2}{2P'}) > 0
\]
for \( \rho > 0 \) and \( |u| < c \).

This says that the system is strictly hyperbolic and genuinely nonlinear on \( \rho > 0 \). Therefore the Glimm’s theory can be applied if
\[
||U_0(x) - U^*||_{L^\infty} + T.V.U_0
\]
is sufficiently small, where \( U^* \) is a constant state such that \( \rho^* > 0,|u^*| < c \). But the vacuum may not be covered by this application of the general theorem.

6 Generalized Darboux formula

In this section we seek an integration formula for solutions of the relativistic Euler-Poisson-Darboux equation. Let us introduce the variables
\[
x = \frac{c}{2} \log \frac{c + u}{c - u}, \quad y = \int_0^\rho \frac{\sqrt{P'}}{\rho + P/c^2} d\rho.
\]
Then the relativistic Euler-Poisson-Darboux equation is
\[
(EPD) \quad \eta_{xx} - \eta_{yy} + A(x, y)\eta_y + B(x, y)\eta_x = 0,
\]
where
\[
A(x, y) = \frac{1}{\sqrt{P'}}(1 - \frac{P'}{c^2} - \frac{\rho + P/c^2}{2P'}) \frac{1 + P'u^2/c^4}{1 - P'u^2/c^3},
\]
\[
B(x, y) = -\frac{2u/c^2}{1 - P'u^2/c^4}(1 - \frac{P'}{c^2} - \frac{\rho + P/c^2}{2P'})P''.
\]
The coefficients \( A \) and \( B \) are of the form
\[
A = \frac{2N}{y} + a, \quad a = \frac{y}{c^2}(a_0 + [x^2/c^2, y^2/c^2]_1),
\]
\[
B = -\frac{4N}{N + 1} \frac{x}{c^2}(1 + [x^2/c^2, y^2/c^2]_1),
\]
where $[X, Y]_1$ denotes a convergent power series $\sum_{j+k \geq 1} c_{jk} X^j Y^k$. In order to remove the singularity in $A$, we use the trick of Weinstein [7]. We introduce the sequence of variables $\eta_j, j = 0, 1, \ldots, N$ by

$$\frac{\partial \eta_j}{\partial y} = y\eta_{j+1},$$

or

$$\eta_j(x, y) = \int_0^y Y\eta_{j+1}(x, Y)dY,$$

where $\eta_0 = \eta$. The sequence of formal integro-differential operators $L_j$ is defined by

$$L_j V = \frac{2}{y} (\frac{N-j}{y} + a) V_y + BV_x + j\tilde{a} V + \sum_{k=1}^{j} F_{jk} V_x + \sum_{k=1}^{j} H_{jk} V,$$

where

$$\tilde{a} = \frac{\partial a}{\partial y} + \frac{a}{y} = \frac{1}{c^2} [x^2/c^2, y^2/c^2].$$

The coefficients $F_{jk}$ and $H_{jk}$ are determined inductively by

$$F_{j+1,k} = \begin{cases} F_{j1} + \frac{1}{y} \frac{\partial B}{\partial y} & \text{if } k = 1 \\ F_{jk} + \frac{1}{y} \frac{\partial}{\partial y} F_{j,k-1} & \text{if } k \geq 2 \end{cases}$$

$$H_{j+1,k} = \begin{cases} H_{j1} + \frac{1}{y} \frac{\partial \tilde{a}}{\partial y} & \text{if } k = 1 \\ H_{jk} + \frac{1}{y} \frac{\partial}{\partial y} H_{j,k-1} & \text{if } k \geq 2 \end{cases}$$

It is easy to see that $F_{jk}$ are of the form $\frac{1}{y} [x^2/c^2, y^2/c^2]_0$ and $H_{jk}$ are of the form $\frac{1}{y^2} [x^2/c^2, y^2/c^2]_0$. By the definition we have formally

$$\frac{1}{y} \frac{\partial}{\partial y} (L_j \eta_j) = L_{j+1} \eta_{j+1}.$$

Now we consider the equation $L_N V = 0$ for $V = \eta_N$ with the initial conditions

$$V = 0, \quad V_y = 2^{N+1} N! \phi(x), \quad \text{at } y = 0.$$

The problem is

$$V_{yy} - V_{xx} = a V_y + BV_x + N\tilde{a}V +$$

$$+ \sum_{k=1}^{N} F_k l^k V_x + \sum_{k=1}^{N} H_k l^k V,$$

$$V = 0, \quad V_y = 2^{N+1} N! \phi(x), \quad \text{at } y = 0.$$
Proposition 9 If $\phi \in C^1(R)$, then the problem (Q) admits a unique solution $V$ in $C^2(R \times [0, \infty))$.

Proof. Let us denote by $H(x, y, V)$ the right hand side of the equation $L_N = 0$. Then (Q) is transformed to the integral equation

$$V(x, y) = 2^N N! \int_{x-y}^{x+y} \phi(\xi)d\xi + \frac{1}{2} \int_{0}^{y} \int_{x-y+}^{x+y} H(X, Y, V)dXdY.$$ 

We can solve this integral equation by the iteration

$$V_0(x, y) = 2^N N! \int_{x-y}^{x+y} \phi(\xi)d\xi,$$

$$V^{n+1}(x, y) = 2^N N! \int_{x-y}^{x+y} \phi(\xi)d\xi + \frac{1}{2} \int_{0}^{y} \int_{x-y+}^{x+y} H(X, Y, V^n)dXdY.$$ 

Fixing $L$ arbitrarily, we consider $|x| \leq L$. Then it is easy to get the estimates

$$|V^{n+1}(x, y) - V^n(x, y)| \leq \frac{M^{n+1}y^{n+1}}{(n+1)!}.$$ 

Therefore $V^n$ tends to a limit $V$ uniformly on $|x| \leq L, 0 \leq y \leq L$. The limit is the unique solution of (Q). QED.

Now we put

$$\eta_N = V, \quad \eta_{N-k} = I\eta_{N-k+1}.$$ 

Since $\eta_{N-k}$ and its derivatives of order $\leq 2$ all vanish on $y = 0$ for $k \geq 1$, we see $\eta = \eta_0$ gives a solution of the relativistic Euler-Poisson-Darboux equation (EPD).

Next we give an integral formula for the solution $V$ of (Q).

Proposition 10 There is a $C^{N+2}$-function $G(x, y, \xi)$ of $|x| < \infty, y \geq 0, x-y \leq \xi \leq x+y$ such that the solution $V$ of (Q) satisfies

$$V(x, y) = \int_{x-y}^{x+y} G(x, y, \xi)\phi(\xi)d\xi.$$ 

Moreover

$$G = 2^N N! + O(y/c^2),$$

$$\partial_x^p \partial_\xi^p \partial^2_x G = O(1/c^2) \quad \text{for } 1 \leq p_1 + p_2 + p_3 \leq N + 2.$$ 

Proof. We consider the approximate solution $V^n(x, y)$ which appeared in the iteration of the proof of Proposition 9. By writing $H$ as

$$H = (aV)_y + (BV)_x + bV + \sum (F_kI^kV)_x + \sum \tilde{H}_kI^kV,$$
where

$$b = N\tilde{a} - a_y - B_x = \frac{1}{c^2}[x^2/c^2, y^2/c^2],$$

$$\tilde{H}_k = H_k - (F_k)_x = \frac{1}{c^2}[x^2/c^2, y^2/c^2],$$

it is easy to see inductively that there is a kernel $G^n(x, y, \xi)$ such that

$$V^n(x, y) = \int_{x-y}^{x+y} G^n(x, y, \xi)\phi(\xi)d\xi.$$

In fact $G^0 = 2$ and $G^n$ are determined inductively by the formula

$$G^{n+1} = 2 + \frac{1}{2}(G^n_{II} + G^n_{I} + G^n_{III} + \sum G^n_{IVk} + \sum G^n_{Vk}),$$

where

$$G_I = \int_{(-x+y+\xi)/2}^{y} a(x - y + Y, Y)G(x - y + Y, Y, \xi)dY + \int_{(x+y-\xi)/2}^{y} a(x + y - Y, Y)G(x + y - Y, Y, \xi)dY,$$

$$G_{II} = \int_{(x+y-\xi)/2}^{y} B(x + y - Y, Y)G(x + y - Y, Y, \xi)dY - \int_{(-x+y+\xi)/2}^{y} B(x - y + Y, Y)G(x - y + Y, Y, \xi)dY,$$

$$G_{III} = \int \int_{D(x,y,\xi)} b(x, Y)G(x, Y, \xi)dXdY,$$

where

$$D(x, y, \xi) = \{(X, Y)|X - Y \leq \xi \leq X + Y, x - y + Y \leq X \leq x + y - Y, 0 \leq Y \leq y\},$$

$$G_{IVk} = \int_{(-x+y+\xi)/2}^{y} F_k(x + y - Y, Y)J^k G(x + y - Y, Y, \xi)dY + \int_{(-x+y+\xi)/2}^{y} F_k(x - y + Y, Y)J^k G(x - y + Y, Y, \xi)dY,$$

where

$$JG(x, y, \xi) = \int_{|x-\xi|}^{y} YG(x, Y, \xi)dY,$$

and

$$G_{Vk} = \int \int_{D(x,y,\xi)} \tilde{H}_k(X, Y)J^k G(X, Y, \xi)dXdY.$$

It is easy to see inductively that

$$|G^{n+1}(x, y, \xi) - G^n(x, y, \xi)| \leq \frac{M^{n+1}y^{n+1}}{(n+1)!}.$$
therefore $G^n$ converges to a limit $G$ uniformly and (6.1) holds. Moreover we can differentiate $G^{n+1}$ supposing that $G^n$ is differentiable. In fact we have

\[
G_{I,x} = \frac{1}{2}aG((x - y + \xi)/2, (-x + y + \xi)/2, \xi)
- \frac{1}{2}aG((x + y + \xi)/2, (x + y - \xi)/2, \xi) + \\
\int_{(-x+y+\xi)/2}^{y} (aG)_x(x - y + Y, Y, \xi)dY + \\
\int_{(x+y-\xi)/2}^{y} (aG)_x(x - Y + Y, Y, \xi)dY,
\]

\[
G_{I,y} = \frac{1}{2}aG((x - y + \xi)/2, (-x + y + \xi)/2, \xi)
- \frac{1}{2}aG((x + y + \xi)/2, (x + y - \xi)/2, \xi) + \\
2aG(x, y, \xi) + \\
\int_{(-x+y+\xi)/2}^{y} (aG)_y(x - y + Y, Y, \xi)dY + \\
\int_{(-x+y-\xi)/2}^{y} (aG)_y(x - Y + Y, Y, \xi)dY,
\]

\[
G_{III,x} = \frac{1}{2}BG((x + y + \xi)/2, (x + y - \xi)/2, \xi) + \\
- \frac{1}{2}BG((x - y + \xi)/2, (-x + y + \xi)/2, \xi) + \\
\int_{(x+y-\xi)/2}^{y} (BG)_x(x + y - Y, Y, \xi)dY + \\
\int_{(-x+y+\xi)/2}^{y} (BG)_x(x - y + Y, Y, \xi)dY,
\]

\[
G_{III,\xi} = \frac{1}{2}BG((x + y + \xi)/2, (x + y - \xi)/2, \xi) + \\
+ \frac{1}{2}BG((x - y + \xi)/2, (-x + y + \xi)/2, \xi) + \\
\int_{(x+y-\xi)/2}^{y} BG_{\xi}(x + y - Y, Y, \xi)dY + \\
\int_{(-x+y+\xi)/2}^{y} BG_{\xi}(x - y + Y, Y, \xi)dY,
\]
\[ G_{II,y} = -\frac{1}{2}BG((x + y + \xi)/2, (x + y - \xi)/2, \xi) + \]
\[ \frac{1}{2}BG((x - y + \xi)/2, (-x + y + \xi)/2, \xi) + \int_{(x+y-\xi)/2}^{y} (BG)_{x}(x + y - Y, Y, \xi) dY + \int_{(-x+y+\xi)/2}^{y} (BG)_{x}(x - y + Y, Y, \xi) dY; \]
\[ G_{III,x} = \int_{(y-\xi)/2}^{y} bG(x + Y, Y, \xi) dY - \int_{(x+y+\xi)/2}^{y} bG(x - Y, Y, \xi) dY, \]
\[ G_{III,\xi} = \int_{0}^{(\epsilon+\xi)/2} bG(\xi + Y, Y, \xi) dY + \int_{0}^{(-x+y+\xi)/2} bG(\xi - Y, Y, \xi) dY + \int \int_{D(y,\xi)} bG(x, Y, \xi) dxdY, \]
\[ G_{III,y} = \int_{(x+y-\xi)/2}^{y} bG(x + y - Y, Y, \xi) dY + \int_{(-x+y+\xi)/2}^{y} bG(x - y + Y, Y, \xi) dY; \]

and the derivatives of \( G_{IVk} \) are similar to \( G_{II} \) and the derivatives of \( G_{IVk} \) are similar to \( G_{III} \). Then it is easy to see inductively that

\[ |G_{x}^{n+1} - G_{x}^{n}| + |G_{\xi}^{n+1} - G_{\xi}^{n}| + |G_{y}^{n+1} - G_{y}^{n}| \leq \frac{M^{n}y^{n}}{n!}. \]

Thus the limit \( G \) is differentiable. In a similar manner we see

\[ |G_{xx}^{n+1} - G_{xx}^{n}| + |G_{\xi\xi}^{n+1} - G_{\xi\xi}^{n}| + |G_{yy}^{n+1} - G_{yy}^{n}| \leq \frac{M^{n}y^{n}}{(n-1)!}. \]

Thus \( G \) is twice continuously differentiable. In a similar manner we see that \( G \) is \( N + 2 \)-times continuously differentiable. The rough estimates stated in the propositions is obvious since the coefficients are all of \( O(1/c^{2}) \). QED.

The solution \( \eta_{N-k} \) enjoys an integral representation

\[ \eta_{N-k} = \int_{x-y}^{x+y} K_{N-k}(x, y, \xi)\phi(\xi) d\xi, \]

where

\[ K_{N-k}(x, y, \xi) = J K_{N-k+1}(x, y, \xi) = J^{k}G(x, y, \xi). \]

So the solution \( \eta \) of the relativistic Euler-Poisson-Darboux equation is given by

\[ \eta(x, y) = \int_{x-y}^{x+y} K(x, y, \xi)\phi(\xi) d\xi, \]
where

\[ K(x, y, \xi) = J^N G(x, y, \xi). \]

By induction we see

\[ J^k G(x, y, \xi) = \frac{2^N N!}{2^k k!} (y^2 - (x - \xi)^2)^k (1 + O(y/c^2)). \]

Thus we have

**Proposition 11** There is a kernel \( K(x, y, \xi) \) which is of \( C^{N+2} \)-class in 

\(|x| < \infty, 0 \leq y, x - y \leq \xi \leq x + y \) such that

\[ \eta(x, y) = \int_{x-y}^{x+y} K(x, y, \xi) \phi(\xi) d\xi \]

gives a solution of the relativistic Euler-Poisson-Darboux equation for any smooth \( \phi \). Moreover

\[ K(x, y, \xi) = (y^2 - (x - \xi)^2)^N (1 + O(y/c^2)). \]

But in order to apply this integration formula, the generalized Darboux formula, to the study of the relativistic Euler equation, more detailed estimates of the remainder are necessary.

**Proposition 12** We have

\[ G_y = O(y/c^2). \]

Proof. Since \( a = O(y/c^2) \), it is clear that \( G_{I,y} = O(y/c^2) \). Next we see

\[ G_{II,y} = -B((x+y+\xi)/2, (x+y-\xi)/2) + B((x-y+\xi)/2, (-x+y+\xi)/2) + O(y/c^2). \]

On the other hand we can write

\[ B = \frac{1}{c^2} B_0(x) + O(y^2/c^2) \]

and

\[ \frac{x+y+\xi}{2} = x + \frac{y+Z}{2}, \quad \frac{x-y+\xi}{2} = x + \frac{-y+Z}{2}, \quad Z = \xi - x. \]

Therefore we see \( G_{II,y} = O(y/c^2) \). It is clear that \( G_{III,y} = O(y/c^2) \) and \( G_{IV,k,y}, G_{V,k,y} = O(y^2/c^2) \). QED.

**Proposition 13** We have

\[ G = 2^N N! + \frac{1}{c^2} C_0(x, c)(\xi - x) + O(y^2/c^2), \]

where \( C_0(x, c) \) is a function of the form

\[ \left[ \frac{x^2}{c^2} \right]_0 + \frac{x}{c^2} \left[ \frac{x^2}{c^2} \right]_0. \]
Proof. It is clear that $G_I = O(y^2/c^2)$ since $a = O(y/c^2)$. Next we see

$$G_{II} = 2^N N! \int_{(x+y-\xi)/2}^{y} B(x+y-Y, Y) dY - 2^N N! \int_{(-x+y+\xi)/2}^{y} B(x-y+Y, Y) dY + O(y^2/c^2),$$

since $G = 2^N N! + O(y/c^2)$. If we write

$$B = \frac{1}{c^2} B_0(x) + O(y^2/c^2), \quad Z = \xi - x$$

then we see

$$\int_{(x+y-\xi)/2}^{y} B(x+y-Y, Y) dY - \int_{(-x+y+\xi)/2}^{y} B(x-y+Y, Y) dY =$$

$$= \frac{1}{c^2} (\int_{x}^{x+\xi/2} B_0(s) ds - \int_{x+\xi/2}^{x+y} B_0(s) ds) + O(y^2/c^2) =$$

$$= \frac{1}{c^2} B_0(x) Z + O(y^2/c^2).$$

Note $|Z| \leq y$. It is clear that $G_{III}, G_{IV, k}, G_{Vk} = O(y^2/c^2)$. QED.

**Proposition 14** We have

$$(G_x + G_\xi) = O(y/c^2).$$

Proof. First we see

$$G_{I,x} + G_{I,\xi} = \int_{(-x+y+\xi)/2}^{y} ((aG)_x + aG_\xi)(x - y + Y, Y, \xi) dY +$$

$$+ \int_{(x+y-\xi)/2}^{y} ((aG)_x + aG_\xi)(x + y - Y, Y, \xi) dY$$

$$= O(y^2/c^2),$$

since $a, a_x = O(y/c^2)$. Next we see

$$G_{III,x} + G_{III,\xi} = \int_{(x+y-\xi)/2}^{y} ((BG)_x + BG_\xi)(x - y - Y, Y, \xi) dY +$$

$$- \int_{(-x+y+\xi)/2}^{y} ((BG)_x + BG_\xi)(x + y - Y, Y, \xi) dY$$

$$= O(y/c^2).$$

It is clear that $G_{III, x}, G_{III, \xi}, G_{Vk, x}, G_{Vk, \xi} = O(y^2/c^2)$. $G_{IV, k, x} + G_{IV, k, \xi}$ is estimated in a similar manner as $G_{III, x} + G_{III, \xi}$. QED.

**Proposition 15** We have

$$(G_x + G_\xi)_y = O(y/c^2).$$
Proof. First we see

\[(G_{I,x} + G_{I,\xi})_{y} = 2((aG)_{x} + aG_{\xi})(x, y, \xi) + \]

\[- \frac{1}{2}((aG)_{x} + aG_{\xi})((x - y + \xi)/2, (x + y - \xi)/2, \xi) + \]

\[- \frac{1}{2}((aG)_{x} + aG_{\xi})((x + y + \xi)/2, (x + y - \xi)/2, \xi) + \]

\[- \int_{(-x+y+\xi)/2}^{y} ((aG)_{x} + aG_{\xi})_{x}(x - y + Y, Y, \xi) dY + \]

\[+ \int_{((x+y-\xi)/2}^{y} ((aG)_{x} + aG_{\xi})_{x}(x + y - Y, Y, \xi) dY \]

\[= O(y/c^2), \]

since \( a, a_{x} = O(y/c^2) \). Next we see

\[(G_{II,x} + G_{II,\xi})_{y} = - \frac{1}{2}((BG)_{x} + BG_{\xi})((x + y + \xi)/2, (x + y - \xi)/2, \xi) + \]

\[+ \frac{1}{2}((BG)_{x} + BG_{\xi})((x - y + \xi)/2, (-x + y + \xi)/2, \xi) + \]

\[+ \int_{(x+y-\xi)/2}^{y} ((BG)_{x} + BG_{\xi})_{x}(x + y - Y, Y, \xi) dY + \]

\[+ \int_{(-x+y+\xi)/2}^{y} ((BG)_{x} + BG_{\xi})_{x}(x - y + Y, Y, \xi) dY \]

\[= 2^{N-1}N!B_{x}((x - y + \xi)/2, (x + y - \xi)/2) \]

\[+ 2^{N-1}N!B_{x}((x + y + \xi)/2, (x + y - \xi)/2) + \]

\[O(y/c^2), \]

since \( G = 2^{N}N! + O(y/c^2) \) and \( G_{x} + G_{\xi} = O(y/c^2) \). But

\[B_{x} = \frac{1}{c^{2}}B_{0}'(x) + O(y^{2}/c^{2}) \]

and

\[B_{x}((x - y + \xi)/2, (-x + y + \xi)/2) - B_{x}((x + y + \xi)/2, (x + y - \xi)/2) = \]

\[= \frac{1}{c^{2}}B_{0}'(x)(-y) + O(y^{2}/c^{2}) \]

\[= O(y/c^2). \]

It is clear that

\[(G_{III,x} + G_{III,\xi})_{y} = \int_{(x+y-\xi)/2}^{y} ((bG)_{x} + bG_{\xi})(x + y - Y, Y, \xi) dY + \]

\[+ \int_{(-x+y+\xi)/2}^{y} ((bG)_{x} + bG_{\xi})(x - y + Y, Y, \xi) dY \]

\[= O(y/c^2). \]
Similarly we can estimate \((G_{IVk,x} + G_{IVk,\xi})_y\), \((G_{Vk,x} + G_{Vk,\xi})_y\) bearing in mind that \((JG)_x + (JG)_\xi = J(G_x + G_\xi)\). QED.

**Proposition 16** We have

\[ G_x + G_\xi = \frac{1}{c^2} C_1(x,c)(\xi - x) + O(y^2/c^2), \]

where \(C_1(x,c)\) is a function of the form

\[ \left[ \frac{x^2}{c^2} \right]_0 + \frac{x}{c^2} \left[ \frac{x^2}{c^2} \right]_0. \]

**Proof.** We already observed that \(G_{Ix} + G_{I\xi} = O(y^2/c^2)\). Next we look at

\[
G_{II,x} + G_{II,\xi} = \int_{(x+y-\xi)/2}^{y} ((BG)_x + BG_\xi)(x + y - Y, Y, \xi) dY + \\
- \int_{(-x+y+\xi)/2}^{y} ((BG)_x + BG_\xi)(x - y + Y, Y, \xi) dY \\
= 2^N N! \int_{(x+y-\xi)/2}^{y} B_x(x + y - Y, Y) dY - 2^N N! \int_{(-x+y+\xi)/2}^{y} B_x(x - y + Y, Y) dY + \\
+ O(y^2/c^2),
\]

since \(G = 2 + O(y/c^2)\) and \(G_x + G_\xi = O(y/c^3)\). Bearing in mind that \(B_y = O(y/c^2)\), we see

\[
\int_{(x+y-\xi)/2}^{y} B_x(x + y - Y, Y) dY - \int_{(-x+y+\xi)/2}^{y} B_x(x - y + Y, Y) dY = \\
= - \int_{(x+y-\xi)/2}^{y} (-B_x + B_y)(x + y - Y, Y) dY - \int_{(-x+y+\xi)/2}^{y} (B_x + B_y)(x - y + Y, Y) dY + \\
+ O(y^2/c^2) \\
= -2B(x,y) + B((x + y + \xi)/2, (x - y - \xi)/2) + \\
+ B((x - y + \xi)/2, (x + y + \xi)/2) + O(y^2/c^2) \\
= \frac{1}{c^2} (-2B_0(x) + B_0(x + \frac{y + Z}{2}) + O(y^2/c^2) \\
= \frac{1}{c^2} B_0'(x) + o(y^2/c^2).
\]

Next we look at

\[
G_{III,x} + G_{III,\xi} = \int_{(x+y-\xi)/2}^{y} bG(x + y - Y, Y, \xi) dY - \int_{(-x+y+\xi)/2}^{y} bG(x - y + Y, Y, \xi) dY + \\
+ \int_{(x+y-\xi)/2}^{y} bG(\xi + Y, Y, \xi) dY - \int_{(-x+y+\xi)/2}^{y} bG(\xi - Y, Y, \xi) dY + \\
+ \int \int_{D(x,y,\xi)} bG(X,Y,\xi) dX dY.
\]
Putting
\[ b(x, y) = \frac{1}{c^2}b_0(x) + O(y^2/c^2), \]
we see
\[
G_{III,x} + G_{III,\xi} = 2^N N! \left( \int_{x}^{x+Z} b_0(s)ds - \int_{x+Z}^{x+2Z} b_0(s)ds \right) + O(y^2/c^2)
\]
\[ = \frac{2^N N!}{c^2}b_0(x) \left( \frac{y+Z}{2} - \frac{y-Z}{2} + \frac{y-Z}{2} - \frac{y+Z}{2} \right) + O(y^2/c^2) \]
\[ = O(y^2/c^2). \]

\[ G_{IVk,x} + G_{IVk,\xi} \]
can be estimated in a similar manner as \( G_{II,x} + G_{III,\xi} \).
Finally \( G_{V_k,x}, G_{V_k,\xi} = O(y^3/c^2) \) since \( J^k G = O(y^3/c^2) \) for \( k \geq 1 \). QED.

**Proposition 17** We have
\[ (G_x + G_\xi)_x + (G_x + G_\xi)_\xi = O(y/c^2). \]

Proof. First we see
\[
(G_{I,x} + G_{I,\xi})_x + (G_{I,x} + G_{I,\xi})_\xi = \frac{1}{2} \left( \int_{x+y-\xi/2}^{x+y+\xi/2} ((aG)_x + 2(aG_\xi)_x + aG_{\xi\xi})dY \right) + \frac{1}{2} \left( \int_{x+y-\xi/2}^{x+y+\xi/2} (aG)_x + 2(aG_\xi)_x + aG_{\xi\xi} \right) \left( x - y + Y, Y, \xi \right) dY
\]
\[ = O(y^2/c^2), \]
since \( a, a_x, a_{xx} = O(y/c^2) \). Next
\[
(G_{II,x} + G_{II,\xi})_x + (G_{II,x} + G_{II,\xi})_\xi = \frac{1}{2} \left( \int_{x+y-\xi/2}^{x+y+\xi/2} ((BG)_x + 2(BG_\xi)_x + (BG)_{\xi\xi})dY \right) + \frac{1}{2} \left( \int_{x+y-\xi/2}^{x+y+\xi/2} ((BG)_x + 2(BG_\xi)_x + (BG)_{\xi\xi}) \right) \left( x - y + Y, Y, \xi \right) dY
\]
\[ = O(y/c^2). \]

It is easy to see
\[ (G_{III,x} + G_{III,\xi})_x + (G_{III,x} + G_{III,\xi})_\xi = O(y/c^2). \]
The estimates of \( G_{IVk} \) and \( G_{V_k} \) can be seen similarly. QED.
**Proposition 18** We have

\[(G_x + G_\xi)_x + (G_x + G_\xi)_\xi = \frac{1}{c^2} C_2(x, c)(\xi - x) + O(y^2/c^2),\]

where \(C_2(x, c)\) is a function of the form

\[\left[\frac{x^2}{c^2}\right]_0 + \frac{x}{c^2}\left[\frac{x^2}{c^2}\right]_0.\]

**Proof.** We already observed that

\[(G_{I,x} + G_{I,\xi})_x + (G_{I,x} + G_{I,\xi})_\xi = O(y^2/c^2).\]

Next, bearing in mind that \(G_x + G_\xi = O(y/c^2)\) and \((G_x + G_\xi)_x + (G_x + G_\xi)_\xi = O(y/c^2),\) we see

\[(G_{II,x} + G_{II,\xi})_x + (G_{II,x} + G_{II,\xi})_\xi =
\int_{(x+y}^{y} - \xi/2) (B_{xx}(x + y, Y) + B_{x}(x - y, Y, \xi)) dY +
\int_{(+}^{y} - x+y/2) (B_{xx}(x - y, Y) dY =
2^{N} N! \int_{x+y}^{y} (-\xi/2) B_{xx}(x+y, Y, Y, \xi) dY - 2^{N} N! \int_{(+}^{y}-x+y/2 B_{xx}(x-y, Y, Y) dY +
O(y^2/c^2).\]

The same discussion to that of the proof of Proposition 16 can be applied by replacing \(B\) by \(B_x\). Let us look at \((G_{III,x} + G_{III,\xi})_x + (G_{III,x} + G_{III,\xi})_\xi\).

Note that

\[(bG)_x + bG_\xi = b_0 G + b(G_x + G_\xi) = 2^{N} N! b_0 + O(y/c^2),\]

\[b G = 2^{N} N! b + O(y/c^2).\]

Applying the discussion of the proof of Proposition 16 by replacing \(b\) by \(b_x\), we see

\[(G_{III,x} + G_{III,\xi})_x + (G_{III,x} + G_{III,\xi})_\xi =
2^{N} N! \int_{x+z}^{x+Z} b_0(s) ds - \int_{x+z}^{x+Z} b_0(s) ds +
O(y^2/c^2)
= -2^{N} N! b_0(x) Z + O(y^2/c^2).\]

The estimates of \(G_{IV,k}, G_{V,k}\) are parallel. QED.
Proposition 19 We have

$$G_{\xi} = \frac{1}{c^2} C_3(x, c) + O(y/c^2).$$

Proof. It is sufficient to note that

$$G_{II,\xi} = 2^{N-1}N!(B((x+y+\xi)/2, (x+y-\xi)/2) + B((x-y+\xi)/2, (-x+y+\xi)/2)) +$$

$$+ O(y/c^2)$$

$$= \frac{2^{N-1}N!}{c^2} (B_0(x+y+\xi/2) + B_0(-x+y+\xi/2)) + O(y/c^2)$$

$$= \frac{2^N N!}{c^2} B_0(x) + O(y/c^2).$$

QED.

Proposition 20 We have

$$(G_x + G_{\xi})_{\xi} = \frac{1}{c^2} C_4(x, c) + O(y/c^2).$$

Proof. We see

$$(G_{I,x} + G_{I,\xi})_{\xi} = O(y/c^2)$$

by $a, a_x = O(y/c^2)$. Next we see

$$(G_{II,x} + G_{II,\xi})_{\xi} =$$

$$= 2^{N-1}N!(B_x((x+y+\xi)/2, (x+y-\xi)/2) +$$

$$+ B_x((x-y+\xi)/2, (-x+y+\xi)/2)) + O(y/c^2)$$

$$= \frac{2^N N!}{c^2} B'_0(x) + O(y/c^2).$$

And we see

$$(G_{III,x} + G_{III,\xi})_{\xi} =$$

$$= 2^N N!b((x-y+\xi)/2, (-x+y+\xi)/2) + O(y/c^2)$$

$$= \frac{2^N N!}{c^2} b_0(x) + O(y/c^2).$$

Other terms can be estimated similarly. QED.

7 Estimates of the derivatives of entropies

Let us consider the entropy $\eta$ generated by $\phi$ of $C^3$-class, that is,

$$\eta(x, y) = \int_{x-y}^{x+y} K(x, y, \xi) \phi(\xi) d\xi.$$
In this section we will find estimates of the derivatives of $\eta$ with respect to $E, F$. As auxiliary variables we introduce

$$R = y^{2N+1}, \quad M = xy^{2N+1}. \quad (7.1)$$

We are going to prove the following

**Proposition 21** We have

$$\frac{\partial \eta}{\partial M} = 2^{2N+1} \int_{0}^{1} (s-s^{2})^{N} D\phi(x + (2s-1)y) ds + O(y^{2}/c^{2}), \quad (7.2)$$

$$\frac{\partial \eta}{\partial R} = 2^{2N+1} \int_{0}^{1} (s-s^{2})^{N} \phi ds +$$

$$2^{2N+1} \int_{0}^{1} (s-s^{2})^{N} \left(-x + \frac{y}{2N+1}(2s-1)\right) D\phi ds + O(y^{2}/c^{2}), \quad (7.3)$$

$$\frac{\partial^{2} \eta}{\partial M^{2}} = 2^{2N+1} y^{-2N-1} \int_{0}^{1} (s-s^{2})^{N} D^{2}\phi ds + O(y^{-2N+1}/c^{2}), \quad (7.4)$$

$$\frac{\partial^{2} \eta}{\partial R \partial M} = 2^{2N+1} y^{-2N-1} \int_{0}^{1} (s-s^{2})^{N} \left(-x + \frac{y}{2N+1}(2s-1)\right) D^{2}\phi ds +$$

$$O(y^{-2N+1}/c^{2}), \quad (7.5)$$

$$\frac{\partial^{2} \eta}{\partial R^{2}} = 2^{2N+1} y^{-2N-1} \int_{0}^{1} (s-s^{2})^{N} \left((-x + \frac{y}{2N+1}(2s-1))^{2} + \frac{4}{(2N+1)^{2}} s(1-s)y^{2}\right) D^{2}\phi(x + (2s-1)y) ds + O(y^{-1}/c^{2}). \quad (7.6)$$

Proof. We write

$$\eta = 2R^{2N+1} \int_{0}^{1} K\left(\frac{M}{R}, \frac{x}{R^{2N+1}}, \frac{M}{R} + (2s-1)\right) D\phi\left(\frac{M}{R} + (2s-1)\right) ds. \quad (7.7)$$

Differentiating $\eta$ with respect to $M$, we have

$$\frac{\partial \eta}{\partial M} = (1) + (2), \quad (1) = 2R^{2N+1} \int_{0}^{1} (K_{x} + K_{\xi})(x, y, x + (2s-1)y) \phi(x + (2s-1)y) ds,$$

$$\quad (2) = 2R^{2N+1} \int_{0}^{1} K(x, y, x + (2s-1)y) D\phi(x + (2s-1)y) ds.$$

Since $K(x, y, \xi) = J^{N}G(x, y, \xi)$, i.e.

$$K(x, y, \xi) = \int_{|x-\xi|}^{y} Y_{N} \int_{|x-\xi|}^{Y_{N}} Y_{N-1} \cdots \int_{|x-\xi|}^{Y_{2}} Y_{1} G(x, Y_{1}, \xi) dY_{1} \cdots dY_{N},$$
by Proposition 16 we see

\[
(K_x + K_\xi)(x, y, x + (2s - 1)y) = \int_{2s-1}^{y} Y_N \int_{2s-1}^{y} Y_{N-1} \cdots \int_{2s-1}^{y} Y_1(G_x + G_\xi)(x, y, x + (2s - 1)y)dy_1 \cdots dy_N
\]

\[
= \frac{C_1(x, c)}{2^N N!c^2} y^{2N+1}(2s-1)(1 - (2s - 1)2)^N + O(y^{2N+2}/c^2)
\]

\[
= -\frac{2^N C_1(x, c)}{(N+1)!c^2} \frac{d}{ds}(s - s^2)^N + O(y^{2N+2}/c^2).
\]

Therefore by integration by part we get

\[
(1) = R^{\frac{-2N}{2N+1}} y^{2N+2} \frac{2^{N+1} C_1(x, c)}{(N+1)!c^2} \int_{0}^{1} (s(1-s))^N \phi \phi(x + (2s - 1)y)ds + O(y^2/c^2)
\]

\[
= O(y^2/c^2).
\]

By Proposition 13 we see

\[
K(x, y, \xi) = \int_{|x-\xi|}^{y} Y_N \int_{|x-\xi|}^{y} Y_{N-1} \cdots \int_{|x-\xi|}^{y} Y_1 G(x, y, \xi)dy_1 \cdots dy_N,
\]

\[
= 2^{2N}(s - s^2)^N y^{2N} + \frac{2^N C_0(x, c)}{N!c^2} (2s - 1)(s - s^2)^N y^{2N+1} + O(y^{2N+2}/c^2).
\]

Therefore by integration by parts we get

\[
(2) = 2^{2N+1} R^{\frac{-2N}{2N+1}} y^{2N} \int_{0}^{1} (s(1-s))^N \phi \phi(x + (2s - 1)y)ds
\]

\[
+ R^{\frac{-2N}{2N+1}} O(y^{2N+2}/c^2).
\]

Thus we have (7.2). Next we show (7.3). We have

\[
\frac{\partial n}{\partial R} = (3) + (4) + (5),
\]

\[
(3) = \frac{2}{2N+1} R^{\frac{-2N}{2N+1}} \int_{0}^{1} K(x, y, x + (2s - 1)y)\phi(x + (2s - 1)y)ds,
\]

\[
(4) = 2 R^{\frac{-2N}{2N+1}} \int_{0}^{1} (-x(K_x + K_\xi)) + \frac{1}{2N+1} y(K_y + (2s - 1)K_\xi)) \times
\]

\[
\phi(x + (2s - 1)y)ds,
\]

\[
(5) = 2 R^{\frac{-2N}{2N+1}} \int_{0}^{1} K(x, y, x + (2s - 1)y)(-x + \frac{y}{2N+1} (2s - 1))D\phi(...)ds.
\]

By Proposition 13 we get

\[
(3) = \frac{2^{2N+1}}{2N+1} \int_{0}^{1} (s - s^2)^N \phi(...)ds + O(y^2/c^2).
\]

As for (4) we use Proposition 16 and

\[
K_y + (2s - 1)K_\xi = yJ^{N-1}G - (2s - 1)(\xi - x)G(x, |x - \xi|, \xi)J^{N-1}1 + (2s - 1)J^N G_\xi
\]
\[
\begin{aligned}
&\quad = 2^{2N+1} N (s - s^2)^N y^{2N-1} + \frac{2^{N-1} C_0(x, c)}{(N - 1)! c^2} (2s - 1)(s - s^2)^N y^{2N} + \\
&\quad + \frac{2^N C_3(x, c)}{N! c^2} (2s - 1)(s - s^2)^N y^{2N} + O(y^{2N+1}/c^2)
\end{aligned}
\]  

(See Proposition 19). Then by integration by parts we have
\[
(4) = \frac{2^{2N+2} N}{2N + 1} \int_0^1 (s - s^2)^N \phi(...) ds + O(y^2/c^2).
\]

As (2) we get
\[
(5) = 2^{2N+1} \int_0^1 (s - s^2)^N (-x + \frac{y}{2N + 1}(2s - 1)) \phi(...) ds + O(y^2/c^2).
\]

Thus we get (7.3).

Next we show (7.4). We have
\[
\frac{\partial^2 \eta}{\partial M^2} = (6) + (7) + (8),
\]

\[
(6) = 2 R^{-2N+1} \int_0^1 ((K_x + K_\xi)_x + (K_x + K_\xi)_\xi)(x, y, \ldots) \times \phi(...) ds,
\]

\[
(7) = 4 R^{-2N+1} \int_0^1 (K_x + K_\xi)(x, y, \ldots) D\phi(...) ds,
\]

\[
(8) = 2 R^{-2N+1} \int_0^1 K(x, y, \ldots) D^2\phi(...) ds.
\]

By Proposition 18 we have
\[
((K_x + K_\xi)_x + (K_x + K_\xi)_\xi)(x, y, x + (2s - 1)y) = \\
\quad = \frac{2^N C_2(x, c)}{N! c^2} (s - s^2)^N (2s - 1)y^{2N+1} + O(y^{2N+2}/c^2).
\]

Thus by integration by parts we get
\[
(6) = O(y^{-2N+1}/c^2).
\]

By the same discussion as (1) we see (7) = O(y^{-2N+1}/c^2). By the same discussion as (2) we see
\[
(8) = 2^{2N+1} y^{-2N-1} \int_0^1 (s - s^2)^N D^2\phi(...) ds + O(y^{-2N+1}/c^2).
\]

Thus we get (7.4).

Next we show (7.5). We see
\[
\frac{\partial^2 \eta}{\partial M \partial R} = (9) + (10) + (11) + (12) + (13) + (14),
\]
\begin{align*}
(9) &= -\frac{4N}{2N + 1} R^{-\frac{4N-1}{2N+1}} \int_0^1 (K_x + K_\xi)(x, y, \ldots) \phi(\ldots) ds,
(10) &= 2R^{-\frac{4N-1}{2N+1}} \int_0^1 (-x(K_x + K_\xi)x + (K_x + K_\xi)\xi) + \\
&\quad + \frac{y}{2N + 1}((K_x + K_\xi)y + (2s - 1)(K_x + K_\xi)\phi(\ldots)ds,
(11) &= 2R^{\frac{4N}{2N+1}} \int_0^1 (K_x + K_\xi)(-x + \frac{y}{2N + 1}(2s - 1))D\phi ds,
(12) &= -\frac{4N}{2N + 1} R^{-\frac{4N-1}{2N+1}} \int_0^1 KD\phi ds,
(13) &= 2R^{-\frac{4N-1}{2N+1}} \int_0^1 (-x(K_x + K_\xi) + \frac{y}{2N + 1}(K_y + (2s - 1)K_\xi))D\phi ds
(14) &= 2R^{-\frac{4N-1}{2N+1}} \int_0^1 K(-x + \frac{y}{2N + 1}(2s - 1))D^2\phi ds.
\end{align*}

We already know that (9) = \(O(y^{-2N+1}/c^2}\). (Recall (1).) Next we look at (10). The first term is \(O(y^{-2N+1}/c^2}\). (Recall (6)). By Proposition 16 and 20 we see

\( (K_x + K_\xi)_y + (2s - 1)(K_x + K_\xi)_\xi = \)
\[
= \frac{2^{N-1}C_1(x, c)}{(N-1)c^2} y^{2N}(2s - 1)(s - s^2)^N + \\
+ \frac{2^N C_4(x, c)}{N!c^2} y^{2N}(2s - 1)(s - s^2)^N + \\
- \frac{2^{N-1}C_1(x, c)}{(N-1)c^2} y^{2N}(2s - 1)^3(s - s^2)^{N-1} + O(y^{2N+1}/c^2).}
\]

By integration by parts we see (10) = \(O(y^{-2N+1}/c^2}\). We already know (11) = \(O(y^{-2N+1}/c^2}\). Clearly

\[
(12) = -\frac{2^{2N+2}N}{2N + 1} y^{-2N-1} \int_0^1 (s - s^2)^N D\phi ds + O(y^{-2N+1}/c^2).}
\]

We see

\[
(13) = O(y^{-2N+1}/c^2) + \frac{2}{2N + 1} R^{-\frac{4N}{2N+1}} \int_0^1 (K_y + (2s - 1)K_\xi)D\phi ds.
\]

As (4) we have

\[
(13) = \frac{2^{2N+2}N}{2N + 1} y^{-2N-1} \int_0^1 (s - s^2)^N D\phi ds + O(y^{-2N+1}/c^2).
\]

Finally we see

\[
(14) = 2^{2N+1}y^{-2N-1} \int_0^1 (s - s^2)^N (-x + (2s - 1)\frac{y}{2N + 1})D^2\phi ds + O(y^{-2N+1}/c^2).}
\]
Recall (5). Summing up we get (7.5).

Next we show (7.6).

\[ \frac{\partial^2 \eta}{\partial R^2} = \frac{\partial}{\partial R} (3) + \frac{\partial}{\partial R} (4) + \frac{\partial}{\partial R} (5), \]

\[ \frac{\partial}{\partial R} (3) = (15) + (16) + (17), \]

\[ (15) = -\frac{4N}{(2N+1)^2} R^{2N-1} \int_0^1 K \phi ds, \]

\[ (16) = \frac{2}{2N+1} R^{2N-1} \int_0^1 (-x(K_x + K_\xi) + \frac{y}{2N+1}(K_y + (2s-1)K_\xi)) \phi ds, \]

\[ (17) = \frac{2}{2N+1} R^{2N-1} \int_0^1 K(-x + \frac{y}{2N+1}(2s-1)) \phi ds, \]

\[ \frac{\partial}{\partial R} (4) = (18) + (19) + (20), \]

\[ (18) = -\frac{4N}{2N+1} R^{2N-1} \int_0^1 (-x(K_x + K_\xi) + \frac{y}{2N+1}(K_y + (2s-1)K_\xi)) \phi ds, \]

\[ (19) = 2R^{2N-1} \int_0^1 K'' \phi ds, \]

where

\[ K'' = x(K_x + K_\xi) + x^2((K_x + K_\xi)x + (K_x + K_\xi)\xi) + \]
\[ + \frac{y}{2N+1}((K_x + K_\xi)y + (2s-1)(K_x + K_\xi)\xi) + \]
\[ - \frac{xy}{2N+1}((K_x + K_\xi)y + (2s-1)(K_x + K_\xi)\xi) + \]
\[ + \frac{y}{2N+1}((K_x + K_\xi)y + (2s-1)(K_x + K_\xi)\xi) + \]
\[ + \frac{y^2}{2N+1}((K_x + K_\xi)y + (2s-1)(K_x + K_\xi)\xi) + \]
\[ + \frac{y}{2N+1}((K_y + (2s-1)K_\xi)y + (2s-1)(K_y + (2s-1)K_\xi)\xi) + \]
\[ + \frac{y}{(2N+1)^2}(K_y + (2s-1)K_\xi), \]

\[ (20) = 2R^{2N-1} \int_0^1 (-x(K_x + K_\xi) + \frac{y}{2N+1}(K_y + (2s-1)K_\xi)) \times (-x + \frac{y}{2N+1}(2s-1)) \phi ds, \]

\[ \frac{\partial}{\partial R} (5) = (21) + (22) + (23) + (24), \]

\[ (21) = -\frac{4N}{2N+1} R^{2N-1} \int_0^1 K(-x + \frac{y}{2N+1}(2s-1)) \phi ds, \]

\[ (22) = 2R^{2N-1} \int_0^1 (-x(K_x + K_\xi) + \frac{y}{2N+1}(K_y + (2s-1)K_\xi)) \times (-x + \frac{y}{2N+1}(2s-1)) \phi ds, \]

\[ (23) = 2R^{2N-1} \int_0^1 K(x + \frac{y}{(2N+1)^2}(2s-1)) \phi ds, \]
\begin{align*}
(24) & = 2R^{\frac{-4N+1}{2N+1}} \int_0^1 K(-x + \frac{y}{2N + 1}(2s - 1))^2 D^2 \phi ds. \\
\text{First we see} & \\
(15) &= -\frac{2^{2N+2}N}{(2N+1)^2} y^{-2N-1} \int_0^1 (s - s^2)^N \phi ds + O(y^{-2N+1}/c^2), \\
(16) &= \frac{2^{2N+2}N}{(2N+1)^2} y^{-2N-1} \int_0^1 (s - s^2)^N \phi ds + O(y^{-2N+1}/c^2), \\
(17) &= \frac{2^{2N+1}}{2N + 1} y^{-2N-1} \int_0^1 (-x + \frac{y}{2N + 1}(2s - 1)) D\phi ds + O(y^{-2N+1}/c^2). \\
\text{Thus we have} & \\
\frac{\partial}{\partial R}(3) &= \frac{2^{2N+1}}{2N + 1} y^{-2N-1} \int_0^1 (-x + \frac{y}{2N + 1}(2s - 1)) D\phi ds + O(y^{-2N+1}/c^2). \\
\text{Since (18) is similar to (16), we have} & \\
(18) &= -\frac{2^{2N+3}N^2}{(2N+1)^2} y^{-2N-1} \int_0^1 (s - s^2)^N \phi ds + O(y^{-2N+1}/c^2). \\
\text{Next let us look at (19). We already know} & \\
2R^{\frac{-4N+1}{2N+1}} y \int_0^1 x(K_x + K_{\xi}) \phi ds &= O(y^{-2N+1}/c^2), \\
2R^{\frac{-4N+1}{2N+1}} y \int_0^1 x^2((K_x + K_{\xi})_x + (K_x + K_{\xi})_{\xi}) \phi ds &= O(y^{-2N+1}/c^2). \\
\text{Recalling (10), we see} & \\
\frac{2}{2N + 1} R^{\frac{-4N+1}{2N+1}} y \int_0^1 ((K_x + K_{\xi})_y + (2s - 1)(K_x + K_{\xi})_{\xi}) \phi ds &= O(y^{-2N+1}/c^2), \\
\frac{2}{2N + 1} R^{\frac{-4N+1}{2N+1}} xy \int_0^1 ((K_x + K_{\xi})_y + (2s - 1)(K_x + K_{\xi})_{\xi}) \phi ds &= O(y^{-2N+1}/c^2). \\
\text{When } N = 1, \text{ we have} & \\
(K_y + (2s - 1)K_{\xi})_y + (2s - 1)(K_y + (2s - 1)K_{\xi})_y &= 8(s - s^2) + \frac{C_3}{c^2}(2s - 1)y - \frac{C_0}{c^2}(2s - 1)^2y - \\
&- \frac{2C_3}{c^2}(2s - 1)y + O(y^2/c^2). \\
\text{When } N \geq 2, \text{ there are bounded functions } F_j(x, c) \text{ such that} & \\
(K_y + (2s - 1)K_{\xi})_y + (2s - 1)(K_y + (2s - 1)K_{\xi})_y &= \\
&= 2^{2N+1} N(2N - 1)(s - s^2)^N y^{2N-2} + \frac{F_1(x, c)}{c^2}(2s - 1)(s - s^2)^N y^{2N-1} + \\
&+ \frac{F_2(x, c)}{c^2}(2s - 1)(s - s^2)^N y^{2N-1} + \frac{F_3(x, c)}{c^2}(2s - 1)^3(s - s^2)^N y^{2N-1} + \\
&+ \frac{F_4(x, c)}{c^2}(2s - 1)^3(s - s^2)^N y^{2N-1} + \frac{F_5(x, c)}{c^2}(2s - 1)^5(s - s^2)^N y^{2N-1} + \\
&+ O(y^{2N}/c^2). 
\end{align*}
Thus we see

\[ 2R^{-4N} \int_0^1 ((K_y + (2s-1)K_\xi)_y + (2s-1)(K_y + (2s-1)K_\xi)_y) \phi ds \]

\[ = \frac{2^{2N+2}N}{(2N+1)^2} \int_0^1 (s-s^2)^N \phi ds + O(y^{-2N+1}/c^2). \]

We have

\[ \frac{2}{(2N+1)^2} \int_0^1 (\xi + (2s-1)(2s-1)^{2N+1}) \phi ds \]

\[ = \frac{2^{2N+2}N}{(2N+1)^2} \int_0^1 (s-s^2)^N \phi ds + O(y^{-2N+1}/c^2). \]

Therefore

\[ (19) = \frac{2^{2N+2}N}{(2N+1)^2} \int_0^1 (s-s^2)^N \phi ds + O(y^{-2N+1}/c^2). \]

We see

\[ \frac{2}{(2N+1)^2} \int_0^1 (\xi + (2s-1)(2s-1)^{2N+1}) \phi ds \]

\[ = \frac{2^{2N+2}N}{(2N+1)^2} \int_0^1 (s-s^2)^N \phi ds + o(y^{-2N+1}/c^2). \]

Therefore

\[ \frac{\partial}{\partial R} (4) = \frac{2^{2N+2}N}{(2N+1)^2} \int_0^1 (s-s^2)^N (x + \frac{y}{(2N+1)^2}) \phi ds + O(y^{-2N+1}/c^2). \]

Next we see

\[ \frac{\partial}{\partial R} (5) = \frac{2^{2N+2}N}{(2N+1)^2} \int_0^1 (s-s^2)^N (x + \frac{y}{(2N+1)^2}) \phi ds + O(y^{-2N+1}/c^2). \]

Therefore we get

\[ \frac{\partial}{\partial R} (5) = \frac{2^{2N+2}N}{(2N+1)^2} \int_0^1 (s-s^2)^N (x + \frac{y}{(2N+1)^2}) \phi ds + O(y^{-2N+1}/c^2). \]

Summing up, we have

\[ \frac{\partial^2 \eta}{\partial R^2} = \frac{2^{2N+2}N}{(2N+1)^2} \int_0^1 (s-s^2)^N (x + \frac{y}{(2N+1)^2}) \phi ds. \]
Thus we get (7.6). QED.

Let us recall the standard entropy $\eta^*$. This is generated by
\[
\phi^*(x) = A'c^2\left(\frac{1}{1-u^2/c^2} - \frac{1}{\sqrt{1-u^2/c^2}}\right),
\]
where
\[
A' = (2N + 1)^{-2N} ((2N + 1)/(2N + 3)A)^{\frac{2N+1}{2}} (2N-1)!!/2N+1N!
\]
We note that
\[
D^2\phi^*(x) = A'(1 + \frac{u^2/c^2}{1-u^2/c^2})(2 - \sqrt{1-u^2/c^2}) \geq A'.
\]
We are going to show that the Hessian $D_U^2\eta^*$ dominates any $D_U^2\eta$.

**Proposition 22** For each $\phi$ fixed in $C^3$ we have on each compact subset of $
\{\rho \geq 0\}$
\[
|\langle \xi|D_U^2\eta^*\xi\rangle| \leq C|\xi|D_U^2\eta^*|\xi|,
\]
provided that $c$ is sufficiently large.

By the assumption we have
\[
R = y^{2N+1} = K\rho(1 + \rho^{\frac{2N+1}{4}}/c^2)_1,
\]
\[
\frac{dR}{d\rho} = K + \rho^{\frac{2N+1}{4}}/c^2)_1,
\]
\[
\frac{d^2R}{d\rho^2} = \frac{1}{c^2} [\rho^{\frac{2N+1}{4}}/c^2]_0,
\]
where $K = ((2N + 3)(2N + 1)A)^{\frac{2N+1}{2}}$. Using these, we have
\[
\frac{\partial R}{\partial E} = \frac{dR}{d\rho} \frac{1 + u^2/c^2}{d\rho - Pr^2/c^3}
\]
\[ \begin{align*}
\frac{\partial R}{\partial F} &= -\frac{\partial R}{\partial \rho} \frac{2u/c^2}{1 - \rho u^2/c^4}, \\
\frac{\partial M}{\partial E} &= -\frac{R}{\rho + P/c^2} \frac{1 + P'/c^4}{1 - P'u^2/c^4} + \frac{dR}{d\rho} \frac{1 + u^2/c^2}{1 - P'u^2/c^4}, \\
\frac{\partial R}{\partial F} &= K(1 - 2xu/c^2) + O(y^2/c^2). 
\end{align*} \]

Differentiating once more, we see
\[ \begin{align*}
\frac{\partial^2 R}{\partial E^2} &= -\frac{K^2}{y^{2N+1}} \frac{2u}{c^2} \frac{1 + u^2/c^2}{c^2} + O(y^{-2N+1}/c^2), \\
\frac{\partial^2 M}{\partial E^2} &= \frac{K^2}{y^{2N+1}} u(-2u^2/c^2 - 2ux(1 - u^2/c^2)/c^2) + O(y^{-2N+1}/c^2), \\
\frac{\partial^2 R}{\partial E^2} &= \frac{K^2}{y^{2N+1}} \frac{2u}{c^2} (1 - u^2/c^2) + O(y^{-2N+1}/c^2), \\
\frac{\partial^2 M}{\partial E^2} &= \frac{K^2}{y^{2N+1}} \frac{2u}{c^2} + O(y^{-2N+1}/c^2). 
\end{align*} \]

The chain rule gives
\[ \begin{align*}
\frac{\partial^2 \eta}{\partial E^2} &= \left( \frac{\partial R}{\partial E} \right)^2 \frac{\partial^2 \eta}{\partial R^2} + 2 \frac{\partial R}{\partial E} \frac{\partial M}{\partial E} \frac{\partial^2 \eta}{\partial R \partial M} + \left( \frac{\partial M}{\partial E} \right)^2 \frac{\partial^2 \eta}{\partial M^2} + \frac{\partial^2 R}{\partial E^2} \frac{\partial \eta}{\partial R} + \frac{\partial^2 M}{\partial E^2} \frac{\partial \eta}{\partial M}. 
\end{align*} \]

and so on. Inserting (7.7) and (7.8) into (7.9), and using Proposition 21, we have
\[ \begin{align*}
(\xi|D^2\eta|\xi) &= \frac{2^{2N+1}K^2}{y^{2N+1}} \int_0^1 (s - s^3)^N Z[x] D^2 \phi ds + \frac{2K^2}{y^{2N+1}} \frac{1}{c^2} \frac{1 + u^2/c^2}{c^2} (u\xi_0 - \xi_1) \frac{\partial \eta}{\partial R} + \\
&\quad - \frac{2K^2}{y^{2N+1}} \frac{1}{c^2} \frac{1 + u^2/c^2}{c^2} (u + x(1 - u^2/c^2))(u\xi_0 - \xi_1) \frac{\partial \eta}{\partial M} + \frac{\partial^2 \eta}{\partial E^2} \frac{\partial \eta}{\partial R} + \frac{\partial^2 \eta}{\partial E^2} \frac{\partial \eta}{\partial M}, 
\end{align*} \]

where
\[ \begin{align*}
Z[\xi] &= Z_{00} \xi_0^2 + 2Z_{01} \xi_0 \xi_1 + Z_{11} \xi_1^2, 
\end{align*} \]
\[ Z_{00} = (1 + u^2/c^2)^2((-x + \frac{y}{2N+1}(2s-1))^2 + \frac{4}{(2N+1)^2}s(1-s)y^2) + 2(1 + u^2/c^2)(-u + x(1 + u^2/c^2))(-x + \frac{y}{2N+1}(2s-1)) + (-u + x(1 + u^2/c^2))^2, \]
\[ Z_{01} = -2(1 + u^2/c^2)u/c^2((-x + \frac{y}{2N+1}(2s-1))^2 + \frac{4}{(2N+1)^2}s(1-s)y^2) + (1 + 3u^2/c^2 - 4x(1 + u^2/c^2)/u^2)(-x + \frac{y}{2N+1}(2s-1)) + (-u + x(1 + u^2/c^2))(1 - 2xu/c^2), \]
\[ Z_{11} = \frac{4u^2}{c^4}((-x + \frac{y}{2N+1}(2s-1))^2 + \frac{4}{(2N+1)^2}s(1-s)y^2) + \frac{4u}{c^2}(1 - 2xu/c^2)(-x + \frac{y}{2N+1}(2s-1)) + (1 - 2xu/c^2)^2. \]

It can be shown that
\[ Z[\xi] \geq \kappa s(1-s)y^2, \]
where \( \kappa \) is a positive constant depending on the compact subset of \( \{ \rho \geq 0 \} \).

In fact we see
\[ Z_{00}Z_{11} - Z_{01}^2 = (1 - u^2/c^2) \frac{4}{(2N+1)^2}s(1-s)y^2. \]

On the other hand, we can estimate
\[ \left| \frac{2K^2}{y^{2N+1}c^2}(1 - u^2/c^2) \frac{\partial \eta}{\partial R} \right| \leq \frac{\epsilon}{y^{2N+1}}, \]
\[ \left| \frac{2K^2}{y^{2N+1}c^2}(u + x(1 - u^2/c^2)) \frac{\partial \eta}{\partial M} \right| \leq \frac{\epsilon}{y^{2N+1}}, \]
where \( \epsilon = K'/c^2 \). Let us introduce the parameters
\[ \zeta_0 = \xi_0, \quad \zeta_1 = \xi_1 - u\xi_0. \]

Then we have
\[ Z[\xi] = Q_{00}\zeta_0^2 + 2Q_{01}\zeta_0\zeta_1 + Q_{11}\zeta_1^2, \]
and
\[ Q_{00} = Q_{00}^{(1)}(x)(2s-1)y + Q_{00}^{(2)}(x,s)y^2, \]
\[ Q_{01} = Q_{01}^{(1)}(x)(2s-1)y + Q_{01}^{(2)}(x,s)y^2, \]
\[ Q_{11} = Z_{11} = 1 + O(1/c^2) > 0. \]

Therefore if \( |D^2\phi| \leq C \), we see
\[ |(\xi[D^2\eta]\xi)| \leq \frac{2^{2N+1}K^2C}{y^{2N+1}} \int_0^1 (s - s^2)^N Z[\xi] ds \]
\[ + \frac{12\epsilon}{y^{2N+1}} \int_0^1 (s-s^2)^N \zeta^2 \, ds + O(y^{-2N+1}/c^2) \]
\[ \leq \frac{2^{2N+1}K^2C}{y^{2N+1}} \int_0^1 (s-s^2)^N (Q_{11}(1+\epsilon')\zeta^2 + 2Q_{01}\zeta_0\zeta_1 + Q_{00}\zeta_0^2) \, ds \]
\[ + O(y^{-2N+1}/c^2). \]

But since \( Q_{00} = Q_{01} = 0 \), \( \int_0^1 (s-s^2)^N (2s-1) \, ds = 0 \), we see
\[ \int_0^1 (s-s^2)^N (-2\epsilon'Q_{01}\zeta_0\zeta_1 - \epsilon Q''\zeta_0^2) \, ds = O(y^{-2N+1}/c^2). \]

Therefore we get
\[ |(\xi|D_0^2\eta|\xi)| \leq \frac{2^{2N+1}K^2C(1+\epsilon')}{y^{2N+1}} \int_0^1 (s-s^2)^N Z(\xi) \, ds + O(y^{-2N+1}/c^2). \]

Similarly, if \( D^2\phi^* \geq \mu \), we have
\[ (\xi|D_0^2\eta^*|\xi) \geq \frac{2^{2N+1}K^2\mu(1-\epsilon')} {y^{2N+1}} \int_0^1 (s-s^2)^N Z(\xi) \, ds + O(y^{-2N+1}/c^2). \]

Thus we get
\[ |(\xi|D_0^2\eta|\xi)| \leq \frac{C(1+\epsilon)}{\mu(1-\epsilon')} (\xi|D_0^2\eta^*|\xi) + O(y^{-2N+1}/c^2). \]

But we know
\[ (\xi|D_0^2\eta^*|\xi) \geq \kappa |\xi|^2 y^{-2N+1}. \]

Hence if \( c \) is sufficiently large we get the required estimate. QED.

As for the first derivatives, the following conclusion is now clear.

**Proposition 23** On each compact subset of \( \{\rho \geq 0\} \), we have
\[ |\frac{\partial \eta}{\partial E}| + |\frac{\partial \eta}{\partial F}| \leq C. \]

## 8 Usefull entropies

Let us consider an entropy \( \eta \) generated by \( \phi \), that is,
\[ \eta(x, y) = \int_{x-y}^{x+y} K(x, y, \xi) \phi(\xi) \, d\xi. \] (8.1)

The corresponding entropy flux \( q \) is given by integrating the differential equations
\[ \frac{\partial q}{\partial w} = \lambda_2 \frac{\partial \eta}{\partial w}, \quad \frac{\partial q}{\partial z} = \lambda_1 \frac{\partial \eta}{\partial z}. \]
We can solve these equations as

\[ q = \lambda_2 \eta - \int_z^w \frac{\partial \lambda_2}{\partial w} \eta \, dw \]
\[ = \lambda_1 \eta + \int_z^w \frac{\partial \lambda_1}{\partial z} \eta \, dz. \]

Thus we get the formula

\[ q(x, y) = \int_{x-y}^{x+y} L(x, y, \xi) \phi(\xi) \, d\xi, \quad (8.2) \]

where

\[ L(x, y, \xi) = \lambda_1 K(x, y, \xi) + L_1(x, y, \xi) \]
\[ = \lambda_2 K(x, y, \xi) + L_2(x, y, \xi), \]
\[ L_1(x, y, \xi) = 2 \int_{(x+y-\xi)/2}^{y} \mu_1(x + y - Y, Y) K(x + y - Y, Y, \xi) \, dY, \]
\[ L_2(x, y, \xi) = -2 \int_{(-x+y+\xi)/2}^{y} \mu_2(x - y + Y, Y) K(x - y + Y, Y, \xi) \, dY, \]
\[ \mu_1(x, y) = \frac{\partial \lambda_1}{\partial z} = \frac{1 - u^2/c^2}{2(1 - \sqrt{P}u/c^2)} \left( 1 - \frac{P'}{c^2} + \frac{(\rho + P/c^2)P'}{2P} \right), \]
\[ \mu_2(x, y) = \frac{\partial \lambda_2}{\partial w} = \frac{1 - u^2/c^2}{2(1 + \sqrt{P}u/c^2)} \left( 1 - \frac{P'}{c^2} + \frac{(\rho + P/c^2)P'}{2P} \right), \]

In this section we will construct various kinds of usefull entropies.

1) Let us put

\[ \eta_k^1(x, y) = \int_{x-y}^{x+y} K(x, y, \xi) k^{N+1} e^{k\xi} \, d\xi, \]
\[ \eta_k^2(x, y) = \int_{x-y}^{x+y} K(x, y, \xi) k^{N+1} e^{-k\xi} \, d\xi. \]

**Proposition 24** If \( 1/c^2 \) is sufficiently small, we have

\[ \eta_k^1 > 0, \quad \eta_k^2 > 0 \quad \text{for} \ y > 0, \quad (8.3) \]
\[ \eta_k^1 = 2^N N! y^N (1 + O(y/c^2)) e^{k(x+y)} (1 + O(1/k)), \]
\[ \eta_k^2 = 2^N N! y^N (1 + O(y/c^2)) e^{-k(x-y)} (1 + O(1/k)). \quad (8.4) \]
uniformly on each compact subset of \( \{ y > 0 \} \). Moreover

\[
q_k^1 = \eta_k^1(\lambda_2 + O(1/k)), \\
q_k^2 = \eta_k^2(\lambda_1 + O(1/k))
\]

uniformly on each compact subset of \( \{ y \geq 0 \} \) and

\[
\eta_k^2 q_k^1 - \eta_k q_k^2 = (2^N N!)^2 y^{2(N-1)}(\frac{1}{2N+1} + O(1/c^2))e^{2ky}(y + O(1/k))^3. \tag{8.6}
\]

Proof. Since \( K = (1 + O(y/c^2))(y^2 - (x - \xi)^2)^N \), we see

\[
\eta_k^1 = (1 + O(y/c^2)) \int_{x-y}^{x+y} (y^2 - (x - \xi)^2)^N k^{N+1} e^{k\xi} d\xi
\]

where

\[
f(r) = r^{N+1}e^{-r} \int_{0}^{1} (s(1-s))^N e^{2rs} ds
\]

It is easy to see

\[
e^{-r}f(r) = 2^{-(N+1)N!} + O(1/r)
\]

This implies (8.4). We note

\[
\eta^1 = (1 + O(1/c^2))2^N N!y^{N-1}e^{k(y+y)}(y + O(1/k)) \\
\eta^2 = (1 + O(1/c^2))2^N N!y^{N-1}e^{-k(y-y)}(y + O(1/k))
\]

uniformly on \( \{ y \geq 0 \} \). Let us consider the flux. We have

\[
L_2(x, y, \xi) = -2 \int_{(-x+y+\xi)/2}^{y} \mu_2(x - y + Y, Y)K(x - y + Y, Y, \xi)dY
\]

\[
= -2(\frac{N}{2N+1} + O(1/c^2)) \int_{(-x+y+\xi)/2}^{y} (Y^2 - (x - y + Y - \xi)^2)^N dY
\]

\[
= -((\frac{N}{2N+1}) + O(1/c^2))(y - x + \xi)^N (y + x - \xi)^{N+1},
\]

\[
q^1 - \lambda_2 \eta^1 = -((\frac{N}{2N+1}) + O(1/c^2)) \int_{x-y}^{x+y} (y - x + \xi)^N (y + x - \xi)^{N+1} k^{N+1} e^{k\xi} d\xi.
\]

But

\[
0 \leq \int_{x-y}^{x+y} (y - x + \xi)^N (y + x - \xi)^{N+1} k^{N+1} e^{k\xi} d\xi
\]

\[
= (N + 1)k^N \int_{x-y}^{x+y} (y^2 - (x - \xi)^2)^N e^{k\xi} d\xi
\]
\[ - Nk^N \int_{x-y}^{x+y} (y - x + \xi)^{N-1} (y + x - \xi)^{N+1} e^{k\xi} d\xi \]
\[ \leq (N + 1) \frac{1}{k} \int_{x-y}^{x+y} (y^2 - (x - \xi)^2)^{N} e^{k\xi} d\xi. \]

Thus
\[ q^1 - \lambda_2 \eta^1 = O(1/k) \eta^1. \]

Since
\[ \lambda_2 - \lambda_1 = \frac{\sqrt{P'(1-u^2/c^2)}}{1-P'u^2/c^4} = \left( \frac{1}{2N+1} + O(1/C^2) \right) y, \]
we have
\[ \eta^2 q^1 - \eta^1 q^2 = \eta^1 \eta^2 \left( \left( \frac{1}{2N+1} + o(1/C^2) \right) y + o(1/k) \right). \]

This implies (8.6). QED.

2) Let \( \psi \) be a function in \( C_0^\infty(-1,1) \) such that \( \psi \geq 0, \int \psi = 1 \). We put
\[ \phi_n^3(x) = \psi_n(x) = n \psi(n(x-a)), \]
\[ \phi_n^4(x) = -D \psi_n(x), \]
\[ \eta_n^3(x, y) = \int_{x-y}^{x+y} K(x, y, \xi) \phi_n^3(\xi) d\xi, \]
\[ \eta_n^4(x, y) = \int_{x-y}^{x+y} K(x, y, \xi) \phi_n^4(\xi) d\xi, \]
\[ \eta^3(x, y) = K(x, y, a) X, \]
\[ \eta^4(x, y) = K_\xi(x, y, a) X, \]
\[ q^3(x, y) = L(x, y, a) X, \]
\[ q^4(x, y) = L_\xi(x, y, a) X, \]
\[ X = \begin{cases} 1 & (x - y < a < x + y) \\ \frac{1}{2} & (|x - a| = y) \\ 0 & (|x - a| > y). \end{cases} \]

**Proposition 25** As \( n \to \infty \), we have
\[ \eta_n^3 \to \eta^3, \quad q_n^3 \to q^3, \quad \eta_n^4 \to \eta^4, \quad q_n^4 \to q^4. \]

Moreover
\[ |\eta_n^3| \leq My^{2N}, \quad |q_n^3| \leq My^{2N} (|x| + y), \quad (8.7) \]
\[ |\eta_n^4| \leq My^{2N-1}, \quad |q_n^4| \leq My^{2N-1} (|x| + y), \quad (8.8) \]
\[ \eta^3 q^4 - \eta^4 q^3 = \frac{N}{(2N+1)(N+1)} \left( 1 + O(1/c^2) \right) (y^2 - (x - a)^2)^{2N} \quad (8.9) \]
Proof. We note

\[ K_{\xi} = - (\xi - x) G(x, |\xi - x|, \xi) \frac{1}{2^{N-1}(N-1)!} (y^2 - (x - \xi)^2)^{N-1} + J^N G_{\xi} \]

\[ = (2N(x - \xi) + O(1/c^2)(\xi - x)^2)(y^2 - (x - \xi)^2)^{N-1} + O(1/c^2)(y^2 - (x - \xi)^2)^N, \]

\[ L_{1,\xi} = 2 \int_{(x+y-a)/2}^{y} \mu_1(x + y - Y, Y) K(\xi(x + y - Y, \xi))dY. \]

The estimates (8.7), (8.8) can be seen easily. Let us consider

\[ \eta^3q^4 - \eta^4q^3 = (KL_{\xi} - LK_{\xi})(x, y, a). \]

Suppose \( x - a \geq 0 \). Then

\[ \frac{1}{2}(KL_{\xi} - LK_{\xi}) = K \int_{(x+y-a)/2}^{y} \mu_1 K(\xi(x + y - Y, Y, a)dY - \]

\[ = \frac{1}{2}(KL_{\xi} - LK_{\xi}) \int_{(x+y-a)/2}^{y} \mu_1 K(\xi(x + y - Y, Y, a)dY. \]

We note

\[ 0 \leq \frac{x+y-a}{2} \leq x - y + Y - a \leq x - a \leq y. \]

Hence we have

\[ \int_{(x+y-a)/2}^{y} \mu_1 K(\xi(x + y - Y, Y, a)dY \]

\[ = \left( \frac{N}{2N+1} + O(1/c^2) \right) 2N \int_{(x+y-a)/2}^{y} (x + y - Y - a)(y^2 - (x + y - Y - a)^2)^{N-1}dY + \]

\[ + O(1/c^2) \int_{(x+y-a)/2}^{y} (y^2 - (x + y - Y - a)^2)^N dY \]

\[ = \left( \frac{N^2}{2(2N+1)(N+1)} + O(1/c^2) \right) (x + y - a)^{N-1}(-x + y + a)^N \frac{1}{N(N+1)}(y + (2N+1)(x - a)) \]

\[ + O(1/c^2)(y^2 - (x - a)^2)^N. \]

Thus

\[ K \int_{(x+y-a)/2}^{y} \mu_1 KdY \]

\[ = \left( \frac{N}{2(2N+1)(N+1)} + O(1/c^2) \right) (y^2 - (x - a)^2)^{2N-1}(-x + y + a)(y + (2N+1)(x - a)) \]

\[ + O(1/c^2)(y^2 - (x - a)^2)^{2N}. \]

Also we have

\[ K_{\xi} \int_{(x+y-a)/2}^{y} \mu_1 KdY \]

\[ = \left( \frac{N^2}{(2N+1)(N+1)} + O(1/c^2) \right) (x - a)(-x + y + a)(y^2 - (x - a)^2)^{2N-1} \]

\[ + O(1/c^2)(-x + y + a)(y^2 - (x - a)^2)^{2N}. \]
Hence
\[
\frac{1}{2}(KL_{\xi} - LK_{\xi}) = \left( \frac{N}{2(2N+1)(N+1)} + O(1/c^2) \right)(y^2 - (x-a)^2)^{2N}.
\]

Here we have used
\[
0 \leq (x-a)(y-(x-a)) \leq y^2 - (x-a)^2,
\]
\[
0 \leq (y-x+a)(y+(2N+1)(x-a)) \leq (2N+1)(y^2 - (x-a)^2)
\]
provided that \(0 \leq x-a \leq y\). When \(x-a \leq 0\), we can discuss in a similar manner by using \(L_2\). QED.

3) Let \(\Phi\) be a function in \(C_0^\infty(-1,1)\) such that \(\int \Phi = 0\) and the support \(\text{supp}\Phi\) is \([-1 + \alpha, 1 + \alpha]\), where \(\alpha\) is a small positive number. We put
\[
\psi_n(x) = n\Phi(n(x-a)),
\]
\[
\eta^5_n(x, y) = \int_{x-y}^{x+y} K(x, y, \xi) D^{N+1} \psi_n(\xi) d\xi,
\]
\[
q^5_n(x, y) = \int_{x-y}^{x+y} L(x, y, \xi) D^{N+1} \psi_n(\xi) d\xi;
\]
\[
\hat{\Phi}(x) = \frac{d}{dx}(x \int_{-1}^{x} \Phi),
\]
\[
\hat{\psi}_n(x) = n\hat{\Phi}(n(x-a)),
\]
\[
\eta^6_n(x, y) = \int_{x-y}^{x+y} K(x, y, \xi) D^{N+1} \hat{\psi}_n(\xi) d\xi,
\]
\[
q^6_n(x, y) = \int_{x-y}^{x+y} L(x, y, \xi) D^{N+1} \hat{\psi}_n(\xi) d\xi;
\]
\[
B^3_n = \eta^3_n - \eta^5_n q^3,
\]
\[
B^4_n = \eta^4_n q^5 - \eta^6_n q^4,
\]
\[
B_n = \eta^5_n q^6 - \eta^6_n q^5.
\]

Let us divide the domain \(\Sigma = \{-B \leq x-y \leq x+y \leq B\}\) into the following 5 parts.

\[
S_0 = \left\{ -\frac{1}{n} < x+y-a \leq \frac{1}{n}, -\frac{1}{n} \leq x-y-a < -\frac{1}{n} \right\} \cap \Sigma,
\]
\[
S_1 = \left\{ \frac{1}{n} < x+y-a, x-y-a < -\frac{1}{n} \right\} \cap \Sigma,
\]
\[
S_L = \left\{ -\frac{1}{n} < x+y-a \leq -\frac{1}{n}, x-y-a < -\frac{1}{n} \right\} \cap \Sigma,
\]
\[
S_R = \left\{ \frac{1}{n} < x+y-a, -\frac{1}{n} \leq x-y-a < -\frac{1}{n} \right\} \cap \Sigma,
\]
\[
S = \Sigma - (S_0 \cup S_1 \cup S_L \cup S_R).
\]
Proposition 26 We have

\[ |B_{n}^{3}| \leq M/n, \quad |B_{n}^{4}| \leq M \]  

on \( \Sigma \), and

\[ |B_{n}| \leq M/n \]  

on \( S_{0} \cup S_{1} \cup S \). Moreover, on \( S_{L} \), we have

\[ B_{n} = ny^{2N}A_{1} + y^{N}A_{2} + A_{3}, \]  

where

\[ A_{1} = \left( \frac{N(2^{N}N!)^{2}}{2N+1} + O(1/c^{2}) \right) \left( \int_{-1}^{n} \Phi(x+y-a) \right)^{2}, \]

\[ |A_{2}| \leq M \left( \int_{-1}^{n} |\Phi| + |\Phi(n(x+y-a))| \right), \]

\[ |A_{3}| \leq \frac{M}{n}. \]

On \( S_{R} \), we have

\[ B_{n} = ny^{2N}C_{1} + y^{N}C_{2} + C_{3}, \]

\[ C_{1} = \left( \frac{N(2^{N}N!)^{2}}{2N+1} + O(1/c^{2}) \right) \left( \int_{-1}^{n} \Phi(x-y-a) \right)^{2}, \]

\[ |C_{2}| \leq M \left( \int_{-1}^{n} |\Phi| + |\Phi(n(x-y-a))| \right), \]

\[ |C_{3}| \leq \frac{M}{n}. \]

Proof. For the simplicity, we write \( \eta_{n} = \eta_{n}^{5}, q_{n} = q_{n}^{5}, \hat{\eta}_{n} = \eta_{n}^{6}, \hat{q}_{n} = q_{n}^{6} \).

It is easy to see inductively that, for \( G_{j} = J^{j}G = K_{N-j} \), we have

\[ \partial_{\xi}^{p} G_{j} = J \partial_{\xi}^{p} G_{j-1} \]

for \( j \geq p + 1 \) and

\[ \partial_{\xi}^{p} G_{p} = (-1)^{p}(\xi - x)^{p} G(x, |\xi - x|, \xi) + J \partial_{\xi}^{p} G_{p-1}. \]

Therefore

\[ \partial_{\xi}^{p} K = \partial_{\xi}^{p} G_{N}(x, y, \xi) = 0 \]

for \( p \leq N - 1 \) and \( y = |x - \xi| \). Thus by integration by parts we have

\[ \eta_{n} = (-1)^{N} \partial_{\xi}^{N} K(x, y, x + y) \psi_{n}(x + y) + \]

\[ - (-1)^{N} \partial_{\xi}^{N} K(x, y, x - y) \psi_{n}(x - y) + \]

\[ + F_{n}^{1}(x, y), \]

\[ F_{n}^{1}(x, y) = (-1)^{N+1} \int_{x-y}^{x+y} \partial_{\xi}^{N+1} K(x, y, \xi) \psi_{n}(\xi) d\xi. \]
We see
\[ \partial_{\xi}^{p} L_{2}(x, y, \xi) = -2 \int_{(-x+y+\xi)/2}^{y} \mu_{2} \partial_{\xi} K(x - y + Y, Y, \xi) dY \]
for \( p \leq N - 1 \). Therefore
\[ \partial_{\xi}^{p} L_{2}(x, y, x + y) = \partial_{\xi}^{p} L_{2}(x, y, x - y) = 0 \]
for \( p \leq N - 1 \). Moreover we see
\[ \partial_{\xi}^{N} L_{2}(x, y, x+y) = \partial_{\xi}^{N} L_{2}(x, y, x-y) = 0. \]

Therefore by integration by parts we have
\[
\sigma_{n}(x, y) = q_{n}(x, y) - \lambda_{2} \eta_{n}(x, y) = \begin{cases} 
(\xi-x)^{N}G_{n}(x, |x-\xi|, \xi) + J\partial_{\xi}^{N}G_{N-1}, 
\end{cases}
\]
It is easy to see inductively that
\[ \partial_{\xi}^{p+1} G_{p}(x, y, \xi) = (-1)^{p} \frac{p(p+1)}{2} (\xi-x)^{p-1} G(x, |x-\xi|, \xi) + \sum \partial_{\xi}^{p} G_{p-1}, \]
where \( H_{p} = O(1/c^{2}) \). Therefore
\[ \partial_{\xi}^{N+1} K(x, y, \xi) = (-1)^{N} \frac{N(N+1)}{2} (\xi-x)^{N-1} G(x, |x-\xi|, \xi) + \sum \partial_{\xi}^{N} G_{N-1}. \]
2) Suppose \((x, y) \in S_0\). Then we see

\[
\eta^3 = K(x, y, a) \\
= O\left((y^2 - (x - a)^2)^N\right) \\
= O(n^{-2N}),
\]

\[
\eta^4 = K_\xi(x, y, a) \\
= O\left(|x - a|(y^2 - (x - a)^2)^{N-1}\right) + O((y^2 - (x - a)^2)^N) \\
= O(n^{-2N+1}),
\]

\[
\sigma^3 = L_2(x, y, a) \\
= -2 \int_{-y}^{y} \mu_2 K(x - y + Y, Y, a) dY \\
= O(n^{-2N-1}),
\]

\[
\sigma^4 = L_{2, \xi}(x, y, a) \\
= -2 \int_{-y}^{y} \mu_2 K_\xi(x - y + Y, Y, a) dY \\
= O(n^{-2N}).
\]

Since \(y = O(1/n)\) and \(\psi_n = O(n)\), we see

\[
(-1)^N \partial_\xi^N K(x, y, x+y)\psi_n(x+y) + \\
- (-1)^N \partial_\xi^N K(x, y, x-y)\psi_n(x-y) = \\
= O(n^{-N+1}).
\]

Since \(F_n^1 = O(1)\), we have \(\eta_n = O(1)\). We see

\[
\partial_\xi^N L_2(x, y, x-y) = -2 \int_{0}^{y} \mu_2 \partial_\xi^N K(x-y+Y, Y, x-y) dY = O(n^{-N-1}).
\]

Therefore

\[
-(-1)^N \partial_\xi^N L_2(x, y, x-y)\psi_n(x-y) = O(n^{-N}).
\]

Since

\[
\partial_\xi^{N+1} L_2(x, y, \xi) = \mu_2 \partial_\xi^N K((x - y + \xi)/2, (-x + y + \xi)/2, \xi) + \\
- 2 \int_{-y+y+\xi}^{y} \partial_\xi^{N+1} K(x - y + Y, Y, \xi) dY \\
= O((-x + y + \xi)^N + O(x + y - \xi),
\]

we see

\[
F_n^2(x, y) = (-1)^{N+1} \int_{x-y}^{x+y} \partial_\xi^{N+1} L_2(x, y, \xi)\psi_n(\xi) d\xi \\
= O(n^{-1}).
\]
Hence $\sigma_n = O(n^{-1})$. Therefore

$$
B_n^3 = \eta^3 \sigma_n - \eta_n \sigma^3 = O(n^{-2N-1}),
$$

$$
B_n^4 = \eta^4 \sigma_n - \eta_n \sigma^4 = O(n^{-2N}),
$$

$$
B_n = \eta_n \hat{\sigma}_n - \hat{\eta}_n \sigma_n = O(n^{-1}).
$$

3) Suppose $(x, y) \in S_1$, where $x + y > a + \frac{1}{n}$ and $x - y < a - \frac{1}{n}$. Then $\psi_n(x + y) = \psi_n(x - y) = \hat{\psi}_n(x - y) = 0$. So, $\eta_n = F_n^1, \sigma_n = F_n^2$, and so on. But

$$
F_n^1(x, y) = (-1)^{N+1} \int_{x-y}^{x+y} \partial^{N+1}_\xi K(x, y, \xi) \psi_n(\xi) d\xi
$$

= $(-1)^{N+1} \int_{-1}^{1} (\partial^{N+1}_\xi K(x, y, a + \frac{s}{n}) - \partial^{N+1}_\xi K(x, y, a)) \Phi(s) ds$

= $O(1/n)$

since $\int \Phi = 0$ and $\partial^{N+1}_\xi K$ is Lipschitz continuous. Same estimates hold for $F_n^2, \hat{F}_n^1, \hat{F}_n^2$.

Thus

$$
B_n^3 = \eta^3 \sigma_n - \eta_n \sigma = O(n^{-2N-1}),
$$

$$
B_n^4 = \eta^4 \sigma_n - \eta_n \sigma = O(n^{-2N-1}).
$$

4) Suppose $(x, y) \in S_L$, where $|x + y - a| \leq 1/n$. It is easy to see $\eta^3 = O(n^{-N}), \eta^4 = O(n^{-N+1}), \sigma^3 = O(n^{-N-1}), \sigma^4 = O(n^{-N})$. Since $n(x - y - a) < -1$, we have $\psi_n(x - y) = 0$. Thus $\eta_n = O(n), \sigma_n = F_n^2 = O(1)$.

Therefore

$$
B_n^3 = \eta^3 \sigma_n - \eta_n \sigma = O(n^{-N}),
$$

$$
B_n^4 = \eta^4 \sigma_n - \eta_n \sigma = O(n^{-N}).
$$

Let us estimate $B_n = \eta_n \hat{\sigma}_n - \hat{\eta}_n \sigma_n$. Since

$$
\partial^{N+1}_\xi K = (-1)^N \frac{N(N+1)}{2} (\xi - x)^{N-1} G(x, |x|, \xi) +
$$

$$
+ (\xi - x)^N H_N(x, \xi) + J \partial^{N}_{\xi} G_{N-1},
$$

we have

$$
F_n^1 = (-1)^{N+1} \int_{x-y}^{x+y} \partial^{N+1}_\xi K(x, y, \xi) \psi_n(\xi) d\xi =
$$

$$
= (-1)^{N+1} ((-1)^N \frac{N(N+1)}{2} 2^N N!(a-x)^{N-1} + F'(x, a)) \int_{-1}^{1} \Phi +
$$

$$
+ O(1/n) =
$$

$$
= (-1)^{N+1} \frac{N(N+1)}{2} 2^N N!(a-x)^{N-1} + F'(x, a)) \int_{-1}^{1} \Phi +
$$

$$
+ O(1/n),
$$
where \( F' = O(1/c^2)|x - a|^N, F'' = O(1/c^2) \). On the other hand

\[
\partial_x^N K(x, y, x + y) = (-1)^N y^N G(x, y, x + y).
\]

Hence

\[
\eta_n = n y^N G(x, y, x + y) \Phi(n(x + y - a)) + \frac{N(N + 1)}{2} 2^N N! y^{N-1} \int_{-1}^{n(x+y-a)} \Phi + O(1/n).
\]

Since

\[
\partial_x^{N+1} L_2(x, y, \xi) = \mu_2 \partial_x^N K((x - y + \xi)/2, (-x + y + \xi)/2, \xi) + 2 \int_{(-x+y+\xi)/2}^{y} \mu_2 \partial_x^{N+1} K(x - y + Y, Y, \xi) dY = \left( \frac{N}{2N+1} + O(1/c^2) \right)(-1)^N \frac{-x + y + \xi}{2}^N \times G((x + y + \xi)/2, (-x + y + \xi)/2, \xi) + O(x + y - \xi),
\]

we see

\[
\sigma_n = F_n^2 = (-1)^{N+1} \int_{x-y}^{x+y} \partial_x^{N+1} L_2(x, y, \xi) \psi_n(\xi) d\xi = -\frac{N}{2N+1} 2^N N! y^N \int_{-1}^{n(x+y-a)} \Phi + O(1/n),
\]

where \( L' = O(1/c^2) \). Here we have used

\[
\left( \frac{-x + y + a}{2} \right)^N = (y - \frac{x + y - a}{2})^N = y^N + O(1/n).
\]

Similar estimates hold for \( \hat{\eta}_n, \hat{\sigma}_n \). Thus

\[
B_n = n y^2 N A_1 + y^N A_2 + A_3,
\]

where

\[
A_1 = -G \frac{N}{2N+1} 2^N N!(1 + L') \Phi(\beta) \int_{-1}^{\beta} \Phi + \frac{N}{2N+1} 2^N N!(1 + L') \Phi(\beta) \int_{-1}^{\beta} \Phi = \frac{N}{2N+1} 2^N N! G(1 + L') (\int_{-1}^{\beta} \Phi)^2,
\]

\[
\beta = n(x + y - a).
\]
The estimates on $S_R$ can be obtained in a similar manner considering $\sigma^3, \sigma^4, \bar{\sigma}_n$. QED.

If we put

$$\bar{B}_n^3 = \eta^3\eta_n^6 - \eta_n^6q^3,$$
$$\bar{B}_n^4 = \eta^4\eta_n^6 - \eta_n^6q^4,$$

then the same estimates hold.

9 Compactness of $\eta_t + q_x$

Let us consider an entropy $\eta$ generated by $\phi$ through the generalized Darboux formula and its flux $q$. In this section we will prove

**Lemma 1** Let $U^\Delta$ be the approximate solutions constructed in Section 4. Then $\eta(U^\Delta)_t + q(U^\Delta)_x$ lies in a compact subset of $H^{-1}_{loc}(\Omega)$, $\Omega$ being a bounded open subset of $\{t \geq 0\}$.

Proof. Let $\Phi$ be a test function and we consider

$$J = \int \int (\eta(U^\Delta)_t + q(U^\Delta)_x)\Phi \, dx \, dt$$
$$= N + L + \Sigma,$$
$$N = - \int \eta(U^\Delta(+0, x)\Phi(0, x)dx,$$
$$L = \sum_n \int [\eta(U^\Delta(t, x)]_{t=\Delta t-n}^{\Delta t-n}\Phi(n\Delta t, x)dx,$$
$$\Sigma = \int \sum_{\text{shock}} (\sigma[\eta] - [q])\Phi \, dt.$$

Since $U^\Delta$ is bounded, we see

$$|N| \leq M\|\Phi\|_{C}.$$

Let us look at $L$. We see

$$L = L_1 + L_2,$$
$$L_1 = \sum_{j, n} \Phi(n\Delta t, (2j + 1)\Delta x) \int_{2j\Delta x}^{(2j+2)\Delta x} [\eta(U^\Delta)]_{t=n\Delta t+0}^{\Delta t+0} \, dx,$$
$$L_2 = \sum_{j, n} \int_{2j\Delta x}^{(2j+2)\Delta x} (\Phi(n\Delta t, x) - \Phi(n\Delta t, (2j + 1)\Delta x) \times$$
$$\times [\eta(U^\Delta)]_{t=n\Delta t+0}^{\Delta t+0} \, dx.$$
We note

\[
[\eta(U^\Delta)]_{t=n}^{t=0} = D_U \eta(U^\Delta(n \Delta t + 0, x))[U^\Delta]
+ \int_0^1 (1 - \theta) (D_U^2 \eta(U^\Delta(n \Delta t + 0) + \theta[KU^\Delta])[U^\Delta]) d\theta.
\]

and

\[
\int_{2j \Delta x}^{(2j+2)\Delta x} [U^\Delta] dx = 0
\]

by the scheme. Therefore

\[
|L_1| \leq M ||\Phi||_C \sum_{j,n} \int_j^1 (1 - \theta)|F(\theta, \eta)| d\theta dx,
\]

where

\[
F(\theta, \eta) = ([U^\Delta] D_U^2 \eta(U^\Delta(n \Delta t + 0) + \theta[KU^\Delta])[U^\Delta]).
\]

By Proposition 22 we know \(|F(\theta, \eta)| \leq MF(\theta, \eta^*)\). But in the proof of Proposition 7 we know

\[
\sum_{j,n} \int_0^1 (1 - \theta)F(\theta, \eta^*) d\theta dx \leq C.
\]

Thus we know

\[
|L_1| \leq M ||\Phi||_C.
\]

In the proof of Proposition 7 we know

\[
\sum_{j,n} \int_{2j \Delta x}^{(2j+2)\Delta x} ||U^\Delta||^2 dx \leq C.
\]

Therefore

\[
|L_2| \leq 2^\alpha ||\Phi||_C^\alpha \sum_n \int (\Delta x)^\alpha ||\eta(U^\Delta)|| dx
\]

\[
\leq 2^{\alpha-1} ||\Phi||_C^\alpha \sum_n \int ((\Delta x)^{\alpha+\frac{1}{2}} + (\Delta x)^{\alpha-\frac{1}{2}} ||\eta(U^\Delta)||^2) dx
\]

\[
\leq M ||\Phi||_C^\alpha ((\Delta x)^{\alpha-\frac{1}{2}} + (\Delta x)^{\alpha-\frac{1}{2}} \sum \int ||U^\Delta||^2 dx
\]

\[
\leq M'(\Delta x)^{\alpha-\frac{1}{2}} ||\Phi||_C^\alpha,
\]

where we use the boundedness of \(D_U \eta\) and \(n = O(1/(\Delta x))\). Next we look at \(\Sigma\). Along the shock we have

\[
\sigma[\eta(U)] = [q(U)]
\]

\[
= \int_{\rho_L}^{\rho_H} (-\frac{d\sigma}{d\rho}) \int_0^1 \theta(U - U_L)D_U^2 \eta(U_L + \theta(U - U_L))(U - U_L) d\theta d\rho.
\]
This implies  

$$|\sigma[\eta] - [q]| \leq M(\sigma[\eta^*] - [q^*]).$$

But we know  

$$\int \sum_{\text{shock}} (\sigma[\eta^*] - [q^*]) dt \leq C$$

in the proof of Proposition 7. Therefore  

$$|\Sigma| \leq M||\Phi||_C.$$  

Summing up, we know the compactness. QED.

10 Convergence of approximate solutions

We consider the approximate solutions $U^\Delta$ constructed in Section 4. Since 

$U^\Delta$ is bounded, there is a sequence $U^{\Delta_n}$ and a family of Young measures 

$\nu_{t,x}$ such that $\text{supp} \nu_{t,x} \subset \Sigma = \Sigma_B$ and for any continuous function $f$  

$$f(U^{\Delta_n}(t,x)) \rightarrow \bar{f} = <\nu_{t,x}, f>$$

in $L^\infty$ weak star topology. By Lemma 1, we can apply the compensated compactness theory, and we can assume  

$$(\eta q' - \eta'q)(U^{\Delta_n}) \rightarrow <\nu, q > <\nu, q' > - <\nu, \eta' > <\nu, q >$$

in $L^\infty$ weak star. Here $\eta, q; \eta', q'$ are arbitrary Darboux entropy pairs. Thus we have

Lemma 2 For any pairs $(\eta, q), (\eta', q')$ of Darboux entropies-entropy flux, the identity  

$$<\nu, \eta q' - \eta'q > = <\nu, \eta > <\nu, q' > - <\nu, \eta' > <\nu, q >$$

holds a.e. $(t,x)$, where $\nu = \nu_{t,x}$.  

Since entropies we will use are countably many, we can assume that the above identity holds outside a null set which is common to all $\eta$. We fix $(t,x)$ at which the identity holds, and we write $\nu = \nu_{t,x}$. Of course $\text{supp} \nu \subset \Sigma$. Suppose that $\text{supp} \nu \cap \{\rho > 0\} \neq \emptyset$. Let $\Sigma_0$ be the smallest triangle $\{z_0 \leq z \leq w \leq w_0\}$ such that $\text{supp} \nu \cap \{\rho > 0\} \subset \Sigma_0$. Let us denote by $P_0$ the state $(w_0, z_0)$. It will be verified that $\nu = \delta_{P_0}$. (the Dirac measure). First we show

Proposition 27

$$P_0 \in \text{supp} \nu.$$
Proof. Suppose $P_0 \not\in \text{supp.} \nu$. Since $\Sigma_0$ is the smallest triangle containing $\text{supp.} \nu \cap \{ \rho > 0 \}$, $w = w_0$ and $z = z_0$ intersect with $\text{supp.} \nu \cap \{ \rho > 0 \}$. On neighborhoods of these intersection points we have

$$\eta^1 \geq \frac{1}{M} e^{k(w_0 - \epsilon)},$$
$$\eta^2 \geq \frac{1}{M} e^{-k(z_0 + \epsilon)}.$$  

(See Proposition 24). Since $\nu, \eta^1, \eta^2$ are nonnegative, we see

$$<\nu, \eta^1> \geq \frac{1}{M} e^{k(w_0 - \epsilon)},$$
$$<\nu, \eta^2> \geq \frac{1}{M} e^{-k(z_0 + \epsilon)}.$$  

Since $P_0 \not\in \text{supp.} \nu$, we have

$$<\nu, \eta^2 q^1 - \eta^1 q^2> \leq M e^{k(z_0 - w_0)}.$$

Taking $2\epsilon < \delta$, we have

$$\left| \frac{<\nu, q^1>}{<\nu, \eta^1>} - \frac{<\nu, q^2>}{<\nu, \eta^2>} \right| = \left| \frac{<\nu, \eta^2 q^1 - \eta^1 q^2>}{<\nu, \eta^1><\nu, \eta^2>} \right| \leq M e^{-k(\delta - 2\epsilon)} \rightarrow 0$$

as $k \rightarrow \infty$. Let $\beta$ be a sufficiently small positive number, and we put

$$\Sigma_2 = \{ z_0 \leq z \leq w < w_0 - \beta \}$$
$$\Sigma_3 = \{ z_0 \leq z \leq w \leq w_0, w_0 - \beta \leq w \}.$$  

Then

$$\eta^1 e^{-kw} = (1 + O(1/c^2))2N!y^{N-1}(y + O(1/k))$$

is bounded on $\Sigma_0$ and we have

$$<\nu|_{\Sigma_2}, \eta^1> \leq M e^{k(w_0 - \beta)}.$$  

Taking $\epsilon = \beta/2$, we know

$$\frac{<\nu|_{\Sigma_2}, \eta^1>}{<\nu, \eta^1>} \leq M e^{-\beta k/2} \rightarrow 0.$$  

Since $\partial \lambda_2/\partial w > 0$, we know

$$\lambda_2(w, z) \geq \lambda_2(w_0 - \beta, z_0)$$
on $\Sigma_3$. Therefore we have

\[
\frac{<\nu,q_1^1>}{<\nu,\eta_1^1>} = \frac{<\nu|_{\Sigma_2}\eta_2^1\lambda_2>}{<\nu,\eta_1^1>} + \frac{<\nu|_{\Sigma_3}\eta_2^1\lambda_2>}{<\nu,\eta_1^1>} + O(1/k)
\geq o(1) + \lambda_2(w_0 - \beta, z_0)
\]

Similarly we see

\[
\frac{<\nu,q_2^2>}{<\nu,\eta_2^2>} \leq o(1) + \lambda_1(w_0, z_0 + \beta).
\]

Therefore we have

\[
\lambda_2(w_0 - \beta, z_0) - \lambda_1(w_0, z_0 + \beta) \leq 0 + o(1).
\]

Passing to the limit, we know

\[
\lambda_2(w_0, z_0) \leq \lambda_1(w_0, z_0).
\]

But this means $P_0 \in \{\rho = 0\}$, a contradiction. QED.

Let us fix $a$ such that $z_0 < a < w_0$. We have

\[
<\nu, B_n^3> = <\nu, \eta_5^5> <\nu, q_n^6> - <\nu, q_n^3>,
\]

\[
<\nu, B_n^4> = <\nu, \eta_5^6> <\nu, q_n^6> - <\nu, q_n^4>,
\]

\[
<\nu, \eta_5^4 q_4^4 - \eta_5^2 q_2^3> = <\nu, \eta_5^4> <\nu, q_4^4> - <\nu, \eta_5^2> <\nu, q_2^3>,
\]

\[
<\nu, B_n> = <\nu, \eta_5^6> <\nu, q_n^6> - <\nu, \eta_5^6> <\nu, q_n^5>.
\]

From (8.8) we know

\[
<\nu, \eta_5^4 q_4^4 - \eta_5^2 q_2^3> > 0
\]

and from (8.10) we know

\[
<\nu, B_n^3> \to 0
\]

Using these we can prove the following propositions. Proofs can be found in Chen et al [2].

**Proposition 28** As $n \to \infty$, $<\nu, \eta_5^5>, <\nu, q_n^5>, <\nu, q_n^6>, <\nu, q_n^6>$ are bounded.

**Proposition 29** As $n \to \infty$, we have $<\nu, B_n> \to 0$.

Now, taking

\[
\Phi_0(x) = \begin{cases} 
  e^{-\frac{1}{1-x^2}} & \text{if } |x| < 1 \\
  0 & \text{if } |x| \geq 1
\end{cases}
\]
we put
\[
\Phi(x) = \frac{1}{\beta}(\Phi_0(\frac{x+\beta}{\beta}) - \Phi_0(\frac{x-\beta}{\beta}))
\]
for the generating function of $\eta_5$. Here $\beta = (1 - \alpha)/2$. We put
\[
S_+ = \{z \leq w, |w - a| \leq \frac{1 - 3\alpha}{n}\},
\]
\[
S_- = \{z \leq w, |z - a| \leq \frac{1 - 3\alpha}{n}\}.
\]

**Proposition 30** As $n \to \infty$, we have
\[
<\nu|_{S_+}, ny^{2N}> + <\nu|_{S_-}, ny^{2N}> \to 0.
\]

**Proof.** Put $S'_L = S_+ \cap S_L, S'_R = S_- \cap S_R$. It is sufficient to prove that
\[
<\nu|_{S'_L}, ny^{2N}> + <\nu|_{S'_R}, ny^{2N}> \to 0.
\]
From (8.11) we have
\[
<\nu|_{S_+}, ny^{2N}A_1 + y^N A_2> + <\nu|_{S_-}, ny^{2N}C_1 + y^N C_2> \to 0.
\]
Note
\[
A_1 = \left(\frac{N(2^N N!)}{2^N + 1} + O(\frac{1}{c})\right)(\int_{-1}^{n(x+y-a)} \Phi)^2 \geq \frac{1}{M_0} > 0
\]
on $S'_L$. Put
\[
E_n = \{0 \leq y \leq (\frac{1}{n})^\mu\},
\]
where $\mu$ is a positive parameter. Then $|y^N A_2| \leq M(1/n)^\mu N = o(1)$ on $S_L \cap E_n$ and $|y^N A_2| \leq Mny^{2N}(1/n)^{1-\mu N}$ on $S_L - E_n$. Choose $d_n \searrow 0$ such that
\[
\int_{-1+\alpha}^{-\alpha-d_n} \Phi = -\int_{1-\alpha-d_n}^{1-\alpha} \Phi \geq (1/n)^{\mu_0}.
\]
Then
\[
(\int_{-1}^{H} \Phi)^2 \geq (1/n)^{2\mu_0}
\]
for $|H| \leq 1 - \alpha - d_n$, and
\[
|\Phi(H)| + |\int_{-1}^{H} \Phi| = o(1)
\]
for $1 - \alpha - d_n \leq |H| \leq 1$. Put
\[
S'_n = S_L \cap \{|w - a| \leq \frac{1 - \alpha - d_n}{n}\}.
\]
Then $S'_L \subset S'_n \subset S_L$ and
\[
|y^N A_2| = o(1)$
on $S_L - S_+^n$ and
\[ ny^{2N} A_1 + y^N A_2 \geq ny^{2N} \left( \frac{1}{M} \left(1/n\right)^{2\mu_0} - M \left(1/n\right)^{1-\mu N} \right) \geq 0 \]
on $S_+^n - E_n$. Here we take $0 < 2\mu_0 < 1 - \mu N$. Then
\[ <\nu|_{S_L}, ny^{2N} A_1 + y^N A_2 > \geq <\nu|_{S_L \cap E_n}, ny^{2N} A_1 > + \frac{1}{M_0} <\nu|_{S_L \cap E_n}, ny^{2N} > + o(1) \]
\[ + <\nu|_{S_L - E_n}, ny^{2N} A_1 > + \frac{1}{M_0} <\nu|_{S_L - E_n}, ny^{2N} > + o(1) \]}
\[ \geq \frac{1}{2M_0} <\nu|_{S_L}, ny^{2N} > + o(1). \]

Similarly we know
\[ <\nu|_{S_R}, ny^{2N} C_1 + y^N C_2 > \geq \frac{1}{2M_0} <\nu|_{S_R}, ny^{2N} > + o(1) \]
. Thus we see
\[ <\nu|_{S_L}, ny^{2N} > + <\nu|_{S_R}, ny^{2N} > \to 0. \]
\[ \text{QED.} \]

**Proposition 31** We have
\[ \nu|_{\{\rho > 0\}} = \delta_{P_0}. \]

Proof. Proposition 30 says that the projections $P_w \tilde{\nu}, P_z \tilde{\nu}$ of the measure $\tilde{\nu} = y^{2N} \nu$ admits the Lebesgue lower derivatives which vanish at any $a$. Therefore we can claim that
\[ \text{supp.} \nu \cap \{\rho > 0\} = \{P_0\}. \]
Since $\nu$ is a probability measure, we have
$$\nu|_{\{\rho>0\}} = C\delta_{P_0}.$$ But
$$C(\eta^3q^4 - \eta^4q^3) = C^2(\eta^3q^4 - \eta^3q^3)$$
at $P_0$. Hence $C = 1$. QED.

Summing up we get the final

**Theorem 2** For any $M_0$ there is a positive number $\epsilon_0$ such that if the initial data satisfy
$$0 \leq \rho_0(x) \leq M_0, \quad \left| \frac{c}{2} \log \frac{c+u_0(x)}{c-u_0(x)} \right| \leq M_0.$$ and if $1/c^2 \leq \epsilon_0$, then a subsequence of the approximate solutions $U^\Delta$ converges a.e. to a limit $U$ which is a weak solution of the relativistic Euler equation.

### 11 Acknowledgment

The author expresses his sincere thanks to Dr. Tadayoshi Kano (Osaka University) and Pr. Hiroshi Uesaka (Nihon University) for helpful informations on studies of the Euler-Poisson-Darboux equation. The thesis of Dr. Shigeharu Takeno was often referred during the preparation of the manuscript. This work was, in part, financially supported by National Center for Theoretical Sciences Mathematics Division, R. O. C.

### 参考文献


