
学位申請論文

京都大学教養部 助手

木 哲 介

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Quenching of Intrinsic Luminescence by Coloration
in KI and RbI

~ Tetsusuke HAYASHI

Physics Laboratory, School of General Education,
Kyoto University, Kyoto 606

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Sýnopsis

The luminescence intensity measured under uv-excitation at low temperatures in X-irradiated KI and RbI has been compared with that measured in uncolored crystals. Reduction in the luminescence efficiency caused by the coloration is due to non-radiative annihilation of free excitons or conduction electron-hole pairs through interaction with V-centers. Change in the degree of the luminescence quenching observed with the increase in the photon energy of excitation from the exciton band region toward the interband transition region can be ascribed to the difference between the relaxation probability of excitons and of electron-hole pairs to the self-trapped exciton state. The temperature dependence of the luminescence quenching suggests that a part of free excitons relax to the self-trapped exciton state through thermal activation process

§1. Introduction

Intrinsic luminescence of alkali halide crystals observed with high efficiency at low temperatures is known to be due to radiative decay of so-called self-trapped excitons (STE)¹⁾ The STE is the state resulting from lattice relaxation of an exciton or from recombination of a free electron with a self-trapped hole (V_k -center). As for the self-trapping process of excitons, the possibility that a potential barrier exists between the free exciton state and the STE state was pointed out theoretically by Toyozawa²⁾ and Sumi et al.³⁾ Experimentally, Fontana et al.⁴⁾ explained the temperature dependence of the luminescence efficiency of NaI excited in the exciton band in terms of an exciton state separated from the STE state by a potential barrier. The similar idea was also used by Nishimura⁵⁾ in order to explain the intrinsic luminescence of KI, and by Lushchik et al.⁶⁾ in order to interpret the mechanism of the host-sensitized luminescence in some impurity-doped alkali halides. Recently, Kuusmann et al.⁷⁾ observed a weak emission in alkali iodide crystals which lies close to the fundamental absorption edge, and they assumed it to come from a metastable free exciton state which lies above the STE state. These experimental works provide interesting informations about the possibility of a stable free exciton state in alkali halide, but are rather hypothetical.

It has been reported by several authors⁸⁻¹⁰⁾ that coloration with irradiation of uv-light or ionizing radiation induces quenching of luminescence in some alkali halide crystals.

Especially, Goldstein⁸⁾ observed this phenomenon in the intrinsic luminescence of KI excited in the exciton absorption band at 77 K, and concluded that nonradiative annihilation of the excited state occurs with a probability proportional to the concentration of F-centers produced by irradiation. A similar experiment was performed by the present author and his coworkers¹¹⁾ at various temperatures from 10 K to 100 K, and it was suggested that temperature dependence of the luminescence quenching reveals that of the self-trapping probability of excitons. These experiments were, however, restricted to excitation in the exciton absorption band region.

The present paper is concerned with quenching of the intrinsic luminescence by coloration in KI and RbI. Ratio of the luminescence intensity in colored crystals to that in uncolored crystals is studied under excitation with the light of various energies in the fundamental absorption band. Dependence of the ratio on temperature as well as on the concentration of defects is investigated. It is shown that behavior of the luminescence quenching obtained under excitation in the inter-band transition region is different from that obtained under excitation in the first exciton absorption band. The results are analyzed on the assumption that the luminescence intensity is proportional to the creation probability of the STE state, and the lattice relaxation process of excitons or free electron-hole pairs is discussed.

§2. Experimental Procedures

Single crystals of KI were obtained from the Harshaw Chemical Co., while single crystals of RbI were grown in our laboratory by the Kyropoulos technique. Crystals of about 1mm thick were cleaved and mounted on a copper holder of a cryostat. Temperature of specimen was monitored by a thermocouple attached to the copper holder. By rotation of the sample holder the crystal was able to face toward various windows of the cryostat for X-ray irradiation, for optical absorption measurements and for luminescence excitation.

Excitation of luminescence was made with uv-light from a Seya-Namioka type vacuum monochromator in conjunction with a deuterium lamp through a LiF window of the cryostat. The band-pass was kept below 8 Å. The luminescence was detected at a right angle to the exciting light with an HTV-R636 photomultiplier tube through a Nikon G250 grating monochromator or through appropriate filters.

Coloration of crystals was made with X-rays from a tungsten target tube operated at 40 kV and 20 mA through a beryllium window of the cryostat. The coloration was always carried out near room temperature after the luminescence in the uncolored crystal had been measured at low temperatures. Before the crystal was colored, and after each step of coloration, transmittance of light at various wavelengths was measured at 78 K in order to estimate the density of produced defects. The coloration did not occur uniformly into the crystal, but it seemed to occur apparently deep enough compared with the penetration depth of the monochromatic uv-light used for luminescence excitation.

The number of F-center was estimated approximately from the optical density of F-absorption band with use of Smakula's formula. The irradiation with uv-light during luminescence studies brought negligible effect on the defect concentration.

§3. Experimental Results

Under irradiation with uv-light in the fundamental absorption region, KI luminesces with high efficiency below ~ 100 K. Two emission bands are mainly induced at 3.31 eV and 4.10 eV under excitation in the interband transition region, while an emission band at 3.05 eV is induced especially below ~ 30 K in addition to the 3.31 eV emission under excitation in the first exciton absorption band region.¹²⁾ These emission bands have been confirmed to originate from excited states of a self-trapped exciton.^{13,14)} After the crystal was colored, reduction in the intensity of each of these emission bands was observed at whole temperature range below ~ 100 K. No appreciable change was observed in the peak position or in the shape of these bands. The total intensity of these emission bands is used hereafter as the intensity of the STE luminescence. Excitation spectra of the total luminescence measured at 5 K for an uncolored KI (I_0) and for the colored one (I) are shown by solid curves in Fig. 1. Here the reflection loss at the crystal surface has not been corrected. The number of F-center in the colored crystal was estimated to be $1.72 \times 10^{16}/\text{cm}^3$. Ratio of the luminescence intensity I/I_0 is also shown by the broken curve in Fig. 1. In the upper part of the figure the

one-photon absorption spectrum at 10 K measured by Teegarden and Baldini¹⁵⁾ and the two-photon absorption spectrum at 4.2 K as well as the location of the band gap energy E_g given by Hopfield and Worlock¹⁶⁾ are shown for comparison. The value of the ratio I/I_0 increases with raising the photon energy of excitation from ~6.2 eV which is in the neighborhood of the threshold for two-photon absorption.

Intrinsic luminescence of RbI observed below ~70 K consists of three emission bands at 3.95 eV, 3.14 eV and 2.30 eV¹⁷⁾ The 3.95 eV emission is absent in the case of excitation in the first exciton absorption region. Excitation spectra of the total luminescence measured at 5 K for an uncolored RbI (I_0) and for the colored one (I) are shown by solid curves in Fig. 2 The broken curve shows the ratio of the intensity I/I_0 . Here the number of F-center in the colored crystal was $1.18 \times 10^{16}/\text{cm}^3$ One- and two-photon absorption spectra and the band gap energy^{15,18)} are also shown for comparison. The increase in the value of I/I_0 appears above ~6.1 eV which is in the neighborhood of the two-photon absorption edge, and it is also nearly coincident with the threshold for intrinsic photoconductivity measured by Huggett and Teegarden.¹⁹⁾ It is clear in both KI and RbI that the luminescence is strongly quenched by coloration especially when the luminescence is excited around the first exciton absorption band.

It should be noted that, in the case of excitation above ~6.2 eV for KI or above ~6.1 eV for RbI, the luminescence intensity rises somewhat slowly to a stationary value when the

exciting light is turned on first after coloration. The time period for reaching the stationary value becomes longer as the exciting light intensity decreases. After the luminescence intensity reaches once the saturated value, the response is instantaneous when the exciting light is turned off or turned on again. Such initial slow build-up of the luminescence intensity in the colored crystal does not appear under excitation in the exciton band region. It is suggested from observation of the similar effect in impurity-doped alkali halides^{20,21)} that the slow build-up of the luminescence is due to trapping of free electrons and holes by some centers created with X-ray irradiation which may be associated with traces of impurities. In the present paper, the luminescence intensity of colored crystals I is always given by the saturated value.

The total luminescence intensity in an uncolored KI is nearly constant below ~ 80 K,¹²⁾ while it depends strongly on temperature in colored crystals especially if the excitation is made in the first exciton absorption band region. Typical examples are shown in Fig. 3. The temperature dependence obtained under excitation with 6.01 eV photons is shown in (a), and that obtained under excitation with 6.37 eV photons is shown in (b). The number of F-center in the colored crystal was $3.64 \times 10^{16}/\text{cm}^3$ in both (a) and (b). Figure 4 shows the temperature dependence of the intensity ratio I/I_0 obtained from values in Fig. 3. Closed circles and open circles correspond to values obtained for 6.01 eV excitation and for 6.37 eV excitation, respectively. The excitation at 6.01 eV^(is) typical of

the excitation in the exciton band. It was confirmed that the temperature dependence of I/I_0 is quite similar to the result for 6.01 eV excitation as far as the excitation is made in the high energy side of the first exciton absorption band,* while the temperature dependence becomes weak as the photon energy of excitation is raised toward the interband transition region beyond the first exciton band. The excitation at 6.37 eV is typical of the interband excitation.

* The study of the luminescence quenching under excitation in the low energy side of the first exciton band is omitted in the present work. Because a strong absorption change is induced by coloration in this energy region which is called γ -band,²²⁾ and it brings unavoidable complications.

In the case of RbI the total luminescence intensity in an uncolored crystal is nearly constant below ~60 K, while the intensity in the colored one varies with temperature in a similar way to the case of KI. Temperature dependence of the intensity ratio I/I_0 obtained under excitation in the exciton band (5.98 eV) and under excitation above the band gap energy (6.26 eV) is shown by closed circles and open circles, respectively, in Fig. 5. Here the number of F-center in the colored crystal is $1.85 \times 10^{16} / \text{cm}^3$. It appears in both KI and RbI that the strong temperature dependence of the luminescence quenching is characteristic of exciton excitation.

Quenching of luminescence by coloration depends also on the

density of defects. Coloration with X-rays at room temperature produces mainly F-center and V-center.²²⁾ Since optical absorption measurements proved that V-centers are produced in numbers proportional to F-centers, the number of F-center is used here as a measure of the defect concentration. The quantities $[(I_0/I) - 1]$ derived from values of I/I_0 obtained at various temperatures under excitation with 6.01 eV photons are plotted as a function of F-center density in Fig. 6. It is seen that values of $[(I_0/I) - 1]$ depend linearly on F-center density, so that the ratio I/I_0 can be given by

$$I/I_0 = \frac{1}{1 + \alpha \cdot n_F}, \quad (1)$$

here n_F is the F-center density and α is a constant. In Fig. 7 values of $[(I_0/I) - 1]$ obtained at 5 K are plotted against F-center density in various cases of different values of excitation energy (6.01 eV, 6.26 eV, 6.37 eV and 6.56 eV). The values obtained for 6.26 eV, 6.37 eV and 6.56 eV excitation are somewhat apart from the linear relation in contrast to the case of 6.01 eV excitation as shown by dashed curves. Similar results were obtained for RbI. The relation (1) was confirmed to be applicable when the excitation was made in the exciton band region of RbI below ~60 K.

In addition to experiments on the luminescence intensity as shown above, decay time of the 3.31 eV emission and of the 3.05 eV emission in colored KI was measured at 5 K and 78 K with use of a pulsed light from a spark source. All of observed

decay curves were single exponential in time range longer than ~ 1 μ sec. after pulsed excitation. No appreciable change was found to be induced by coloration in the decay time of each emission even when the intensity of each emission was reduced by coloration down to 1/10. The decay time of the 3.31 eV emission at 78 K was ~ 1.8 μ sec. irrespective of the photon energy of excitation between 5.8 eV and 6.5 eV. At 5 K the decay time of the 3.05 eV emission excited with 6.0 eV photons and that of the 3.31 eV emission excited with 6.5 eV photons were ~ 6.0 μ sec. and ~ 5.5 μ sec., respectively. These values are equal to those measured in the uncolored crystal, and are nearly coincident with those reported by other authors.^{14,23)}

§4. Analysis and Discussion

Irradiation with X-rays near room temperature produces not only F-centers but also V-centers. Goldstein⁸⁾ found that no luminescence quenching occurs in the case of electrolytically colored KI, and concluded that the defect which plays a role on the luminescence quenching is not an isolated F-center. The present author and his coworkers²⁴⁾ reported the luminescence quenching in KI induced by uv-coloration below ~ 200 K, and showed that the quenching phenomenon is concerned with the presence of vacancies or V-centers. As vacancies are not stable at room temperature, it must be true in the present case that the presence of V-centers is responsible for the luminescence quenching.

Optical absorption band due to V-centers appears from ~ 2.5

eV to ~ 5.0 eV with no well-defined structure in both KI and RbI. Most of intrinsic emission bands appear in the energy region of the V-absorption band of each crystal, so that one might consider that the observed reduction in the luminescence intensity is due to reabsorption of emitted photons by V-centers. However, if it were true, the remarkable dependence of the luminescence intensity in colored crystals on the excitation energy or on temperature should not be expected. As the optical density of V-absorption band measured in the present experiment is very small (less than ~ 0.05), the reabsorption effect is considered to be negligible.

Decay time measurements performed for KI supports the inference that the nonradiative annihilation of intrinsic excitation occurs not through relaxed excited states from which photons are emitted. If it were the case, the decay time of the emission should be reduced by coloration. As mentioned in the preceding section, behaviors of the luminescence quenching by coloration depend on whether the excitation is made in the exciton band or in the interband transition region. This also suggests that the annihilation process occurs not through the STE state but in an early stage before the self-trapping.

4.1 Excitation in the exciton band

Let us first discuss the luminescence quenching observed under excitation in the exciton absorption region with use of a simple exciton decay model. Namely, let us assume that a part of excitons are annihilated nonradiatively with a proba-

bility $1/\tau_D$ due to the interaction with V-centers; let $1/\tau_{st}$ be the probability of exciton relaxation to STE states; and let $1/\tau_{nr}$ be that of other nonradiative decays. If we assume that the total intensity of the STE luminescence is proportional to $1/\tau_{st}$, the intensity I can be expressed by

$$I \propto \frac{1/\tau_{st}}{1/\tau_{st} + 1/\tau_D + 1/\tau_{nr}} \quad (2)$$

In addition, assuming that $1/\tau_D$ is proportional to the concentration of V-centers, and so that to the concentration of F-centers ($1/\tau_D = c \cdot n_F$), and that $\tau_{nr} \gg \tau_{st}$, we obtain the relation

$$I/I_0 = \frac{1}{1 + c \cdot \tau_{st} \cdot n_F} \quad (3)$$

This is the similar relation as that derived by Goldstein⁸⁾ in order to explain his experimental results obtained at 77 K. The relation (3) coincides with the experimental relation (1) if we put $c \cdot \tau_{st}$ to α .

Temperature dependence of the value of $c \cdot \tau_{st}$ can be estimated from that of measured I/I_0 according to the relation (3). The obtained values are plotted as a function of $1/T$ in Fig. 8 (a) and (b). The temperature dependence of $c \cdot \tau_{st}$ for each crystal can be approximated by the relation

$$c \cdot \tau_{st} = [A + B \cdot \exp(-\Delta E/kT)]^{-1} \quad (4)$$

Solid curves in Fig. 8 (a) and (b) represent this relation, where constants A, B and ΔE are chosen so as to fit the experimental points of each crystal. Values of energy ΔE chosen here are 18.0 ± 0.5 meV, and 12.2 ± 0.5 meV for KI and for RbI, respectively.

Assuming that the probability $1/\tau_D$ ($= c \cdot n_F$) is independent of temperature, one can attribute the temperature dependence of $c \cdot \tau_{st}$ to that of τ_{st} , and so that the self-trapping probability $1/\tau_{st}$ can be represented from the relation (4) as the form

$$1/\tau_{st} = 1/\tau_{st}^0 + v \exp(-\Delta E/kT). \quad (5)$$

This relation implies that a metastable free exciton state exists above STE states separated by a potential barrier ΔE . Relaxation of excitons to STE states occurs on the one hand with a probability $1/\tau_{st}^0$ which is independent of temperature, and on the other hand occurs thermally over the potential barrier. This conclusion is consistent with hypotheses given by other authors²⁻⁷⁾ mentioned in §1. It can be tentatively said that the temperature dependence of the luminescence quenching is governed by that of self-trapping probability of excitons. A part of excitons are trapped and annihilated near V-center sites efficiently at lower temperatures because they stay at the metastable state. This leads to the strong reduction in the STE luminescence. On the other hand, the increase in the self-trapping probability with the rise of temperature gives rise to the recovery of the luminescence intensity.

The above conclusion, which is based on the assumption

that the probability $1/\tau_D$ ($= c \cdot n_F$) is independent of temperature, is well supported by a previous paper on the edge emission in KI and RbI.²⁵⁾ It was found that the temperature dependence of the intensity of the edge emission is quite similar to that of $c \cdot \tau_{st}$ shown in Fig. 8. The efficiency of the edge emission was approximately given by τ_{st}/τ_r , where τ_r is the radiative lifetime of a free exciton. The good agreement between the temperature dependence of τ_{st}/τ_r and of $c \cdot \tau_{st}$ suggests that the temperature dependence of both quantities comes from that of τ_{st} as given by the relation (5).

4.2 Excitation in the interband transition region

As the photon energy of excitation is raised toward the interband transition region, the intensity ratio I/I_0 measured at 5 K increases as shown in Figs. 1 and 2 as well as the temperature dependence of I/I_0 becomes weak as shown in Figs. 4 and 5. Let us assume that pairs of conduction electrons and holes created with optical excitation relax to STE states through the following two processes. (1) The holes are localized to form self-trapped holes (V_k -centers) and then recombination of electrons with self-trapped holes produces STE states. (2) A part of electron-hole pairs immediately create the lowest free excitons which is followed by the relaxation to STE states. The luminescence intensity I_0 is then given by

$$I_0 = I_0^{re} + I_0^{ex}$$

Here I_0^{re} is the intensity of the luminescence induced through

the former process mentioned above, and I_0^{ex} is that induced through the latter process. The luminescence intensity I in the colored crystal is given by

$$I = I^{\text{re}} + I^{\text{ex}}.$$

From the two equations we obtain

$$I/I_0 = \gamma \cdot (I^{\text{ex}}/I_0^{\text{ex}}) + (1-\gamma)(I^{\text{re}}/I_0^{\text{re}}), \quad (6)$$

where

$$\gamma = I_0^{\text{ex}}/I_0,$$

which corresponds to the production rate of the lowest free exciton.

The temperature dependence of I/I_0 obtained under excitation in the interband transition region can be explained in terms of the relation (6) on the assumption that the fraction γ is independent of temperature and that the value of $I^{\text{ex}}/I_0^{\text{ex}}$ depends on temperature in the same way as the intensity ratio I/I_0 obtained in the case of exciton excitation. In the case of KI the solid curve in Fig. 4 along experimental points for 6.37 eV excitation is the calculated curve with use of the relation (6). Here values of I/I_0 obtained for 6.01 eV excitation have been used as the values of $I^{\text{ex}}/I_0^{\text{ex}}$. The ratio $I^{\text{re}}/I_0^{\text{re}}$ has been tentatively assumed to be independent of temperature. The value of γ and that of $I^{\text{re}}/I_0^{\text{re}}$ chosen so as to fit the experimental points are 0.21 and 0.23, respectively, in this case. The similar analysis was able to perform in the

case of RbI. The solid curve along experimental points for 6.26 eV excitation in Fig. 5 is the calculated curve from the relation (6). The value of γ estimated is 0.15.

Several data of the temperature dependence of I/I_0 were obtained in KI under various excitation near and above the band gap energy. These results were all possible to be analyzed in the same way as above with use of the relation (6) based on the assumption that the ratio $I^{\text{re}}/I_0^{\text{re}}$ is independent of temperature. Values of γ chosen in these analyses are presented in Table I for several cases of different excitation energy. It is likely that the production rate of the lowest free exciton is less than 1/10 of the total excitation when the excitation is made far above the band gap energy. Nouailhat et al.²⁶⁾ reported that the quenching of luminescence by coloration is nearly independent of temperature in KI under excitation with electron beam. This seems consistent with the present result. The dependence of the fraction γ on excitation energy obtained here is consistent with the observation²⁷⁾ that the efficiency of the edge luminescence stimulated with uv-light decreases sharply with raising the photon energy of excitation from the exciton band region toward the interband transition region.

The value of $[(I_0/I) - 1]$ obtained under excitation in the interband transition region appears not to be linear against F-center concentration as shown in Fig. 7. Such a fact also can be expected from the relation (6). It is clear that the ratio I/I_0 does not coincide with the relation (1) if the concentration dependence of $I^{\text{re}}/I_0^{\text{re}}$ is different from that of

$$I^{\text{ex}}/I_0^{\text{ex}}.$$

It is concluded that pairs of conduction electrons and holes created initially with excitation in the interband transition region relax to STE states through the two processes mentioned above. A large part of them recombine to form the STE directly without passing through the lowest (metastable) free exciton state. The thermal activation process is not involved in this process, and the probability of the recombination must be independent of temperature in contrast to the case of exciton excitation. Before the relaxation to form the STE through either of the two processes, a part of excitons or of pairs of conduction electrons and holes are annihilated nonradiatively due to the interaction with V-centers, which leads to the quenching of luminescence. Differences in the behavior of the luminescence quenching observed with changing the photon energy of excitation from the first exciton absorption band into the interband transition region can be attributed to the change in the production rate of the lowest free exciton. It should be noted that there exist absorption bands due to higher states of excitons in the interband transition region. However, no appreciable increase appeared in the production rate of the lowest free exciton γ under excitation in these bands. The creation of higher states of excitons is not able to be distinguished from the creation of free electron-hole pairs.

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Figure Captions

Fig. 1 Excitation spectra of the total luminescence measured at 5 K in an uncolored KI (I_0) and in the colored one (I). Broken curve shows the spectrum of the intensity ratio I/I_0 . The F-center concentration in the colored crystal is $1.72 \times 10^{16}/\text{cm}^3$. Arrows indicate the band-pass of the exciting monochromator. One- and two-photon absorption spectra and the location of the band gap energy E_g are shown for comparison in the upper part.

Fig. 2. Excitation spectra of the total luminescence measured at 5 K in an uncolored RbI (I_0) and in the colored one (I). Broken curve shows the spectrum of the intensity ratio I/I_0 . The F-center concentration in the colored crystal is $1.18 \times 10^{16}/\text{cm}^3$. Arrows indicate the band-pass of the exciting monochromator. One- and two-photon absorption spectra and the location of the band gap energy E_g are shown for comparison in the upper part.

Fig. 3. Temperature dependence of the total luminescence intensity in an uncolored KI (I_0) and in the colored one (I) measured under excitation with 6.01 eV photons (a) and with 6.37 eV photons (b). The F-center concentration in the colored crystal is $3.64 \times 10^{16}/\text{cm}^3$ for both (a) and (b).

- Fig. 4. Temperature dependence of the intensity ratio I/I_0 obtained from values in Fig. 3 for 6.01 eV excitation (closed circles) and for 6.37 eV excitation (open circles). The solid curve along experimental points for 6.37 eV excitation is the calculated one in the text.
- Fig. 5. Temperature dependence of the ratio of the luminescence intensity in a colored RbI to that in the uncolored one (I/I_0) obtained under excitation with 5.98 eV photons (closed circles) and with 6.26 eV photons (open circles). The F-center concentration in the colored crystal is $1.85 \times 10^{16}/\text{cm}^3$ in both cases. The solid curve along experimental points for 6.26 eV excitation is the calculated one in the text
- Fig. 6. Plot of quantities $[(I_0/I) - 1]$ obtained in KI as a function of the F-center concentration. The luminescence intensity, I_0 and I , was measured under 6.01 eV excitation at various temperatures given in the figure.
- Fig. 7 Plot of quantities $[(I_0/I) - 1]$ obtained in KI at 5 K as a function of the F-center concentration. Photon energies of the luminescence excitation are given in the figure
- Fig. 8 The value of $c \cdot \tau_{st}$ plotted as a function of $1/T$ obtained for KI (a) and for RbI (b) The solid curve in each figure represents the calculated curve in the text.

Table I. The production rate of the lowest free exciton γ for various excitation energies in KI.

photon energy of excitation	6.26 eV	6.37 eV	6.56 eV	7.00 eV
γ	0.46	0.21	0.13	0.08

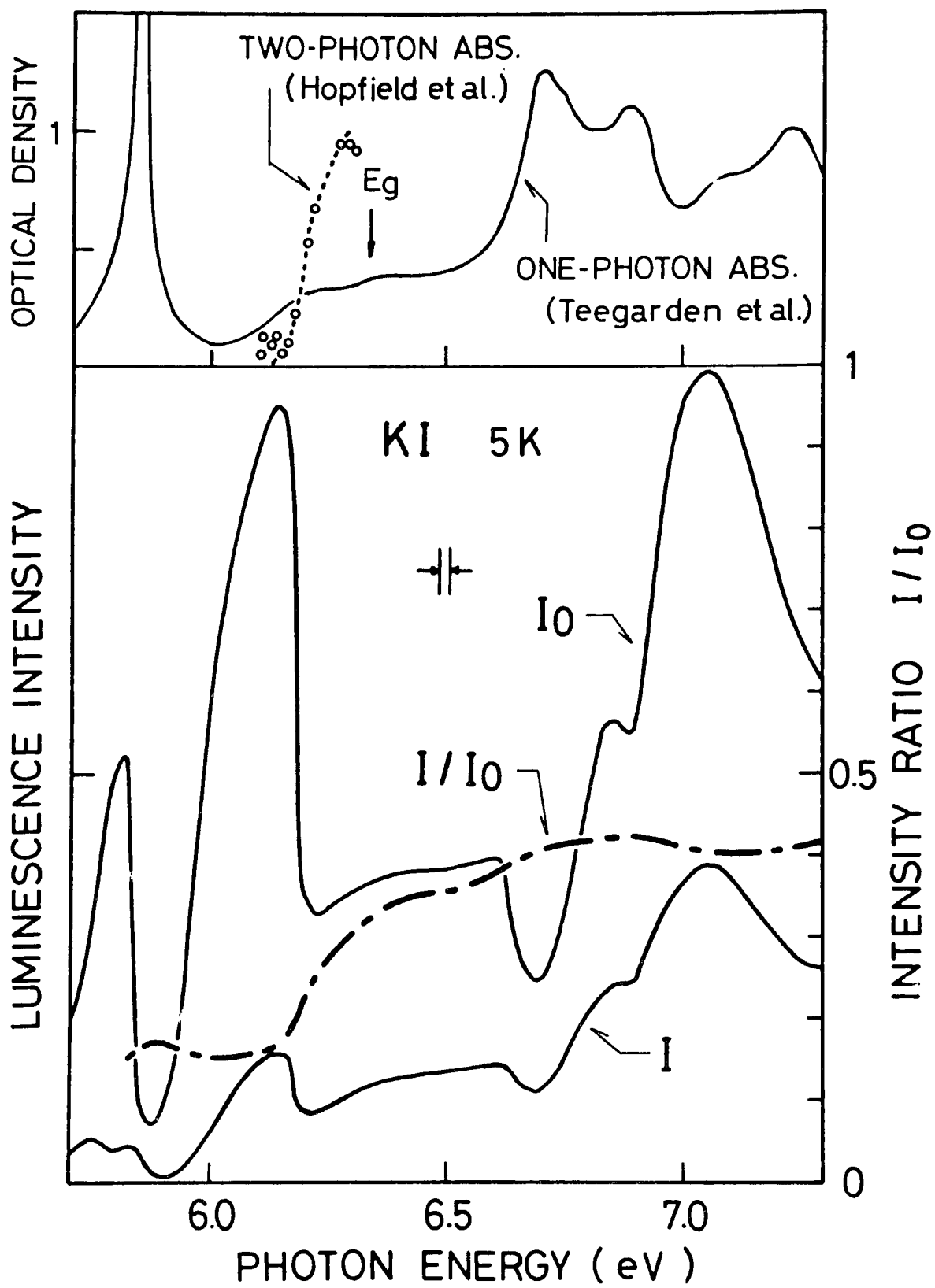


Fig. 1

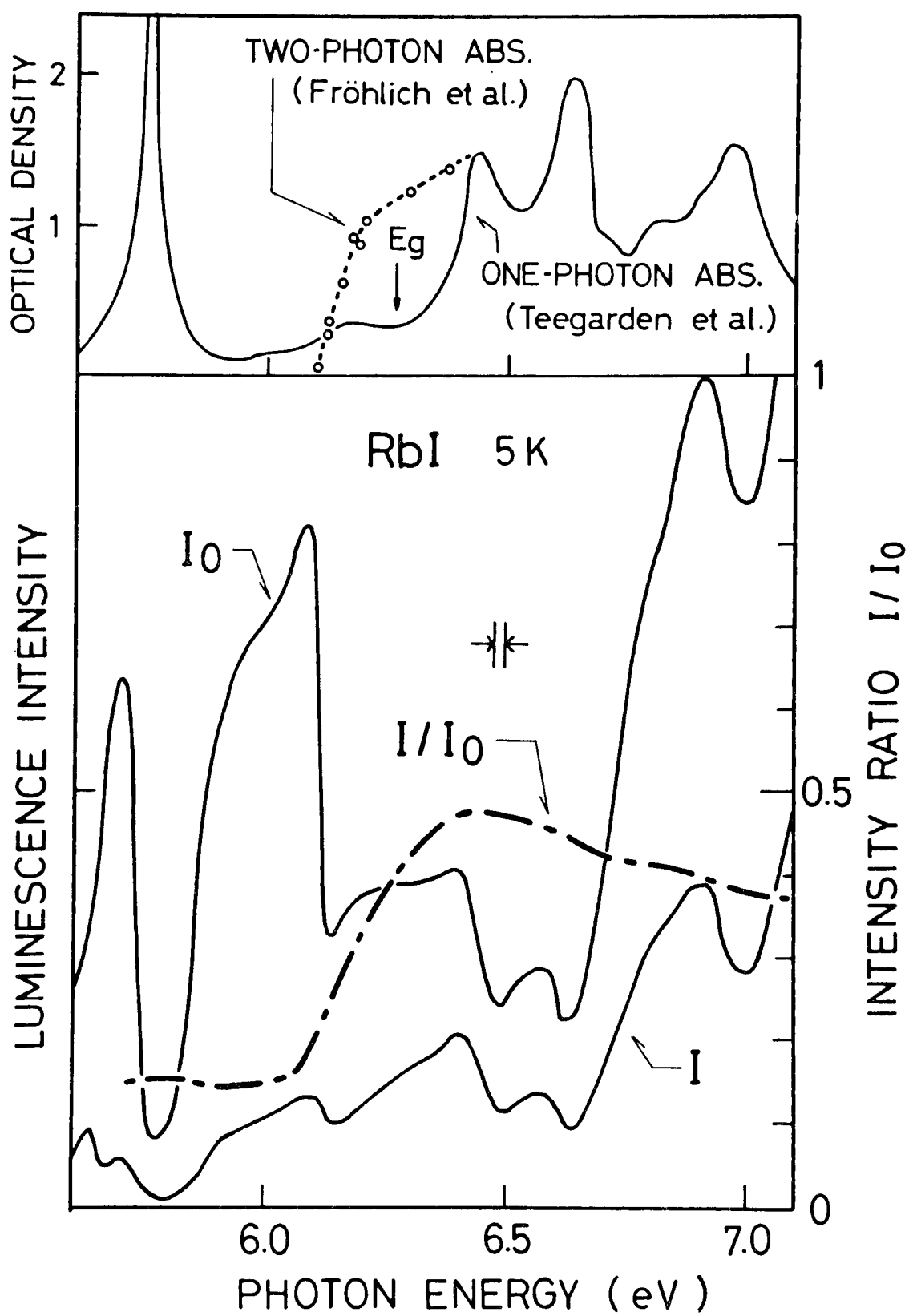


Fig. 2

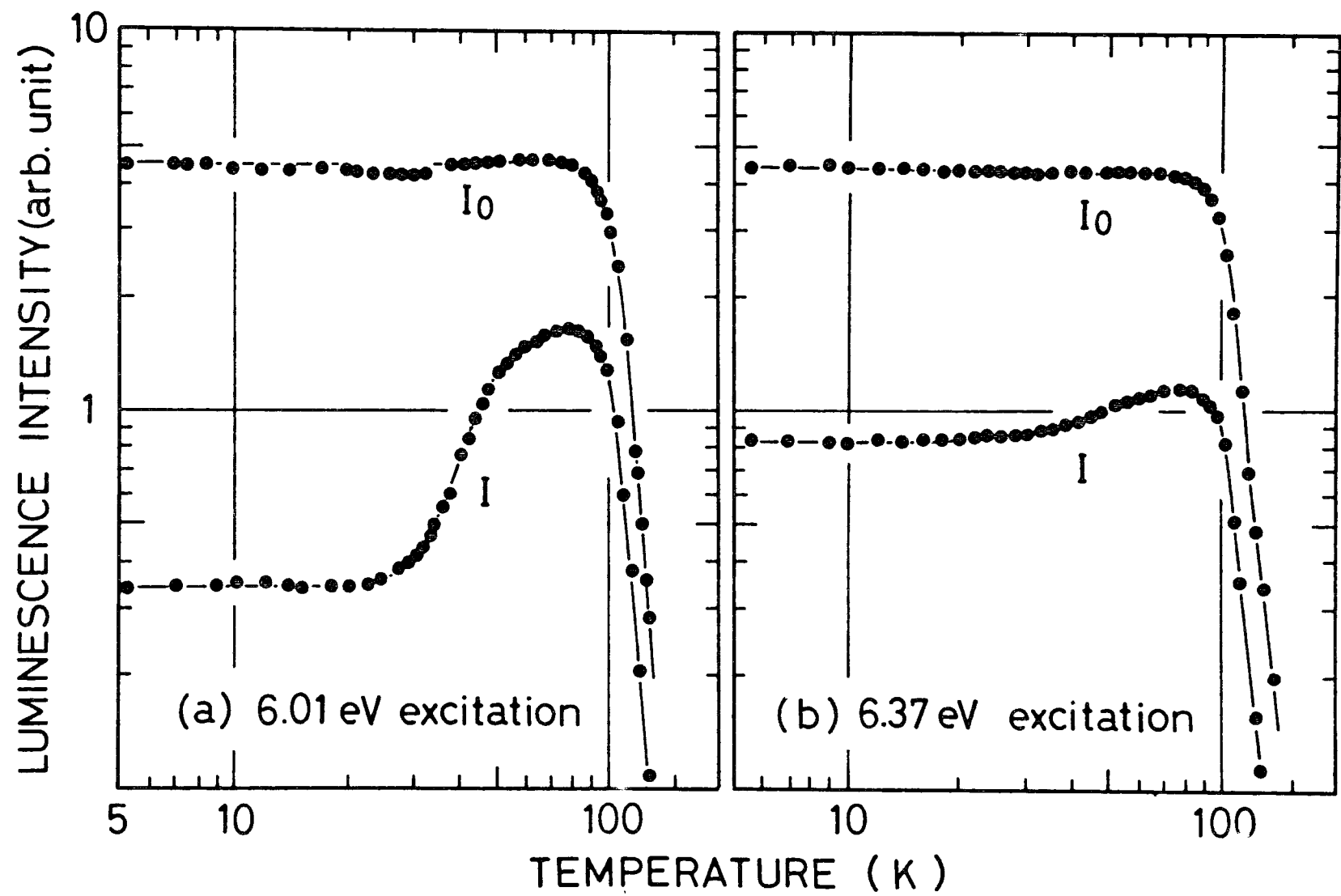


Fig.3

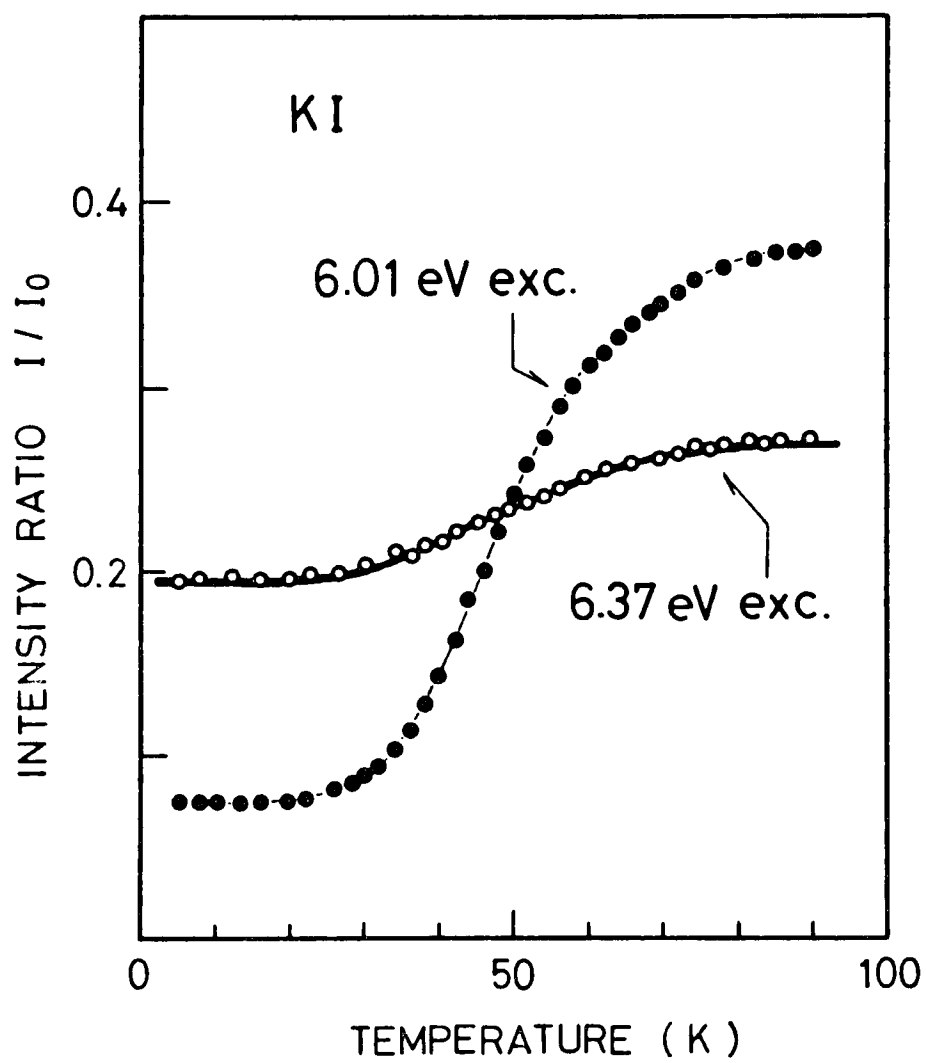


Fig.4

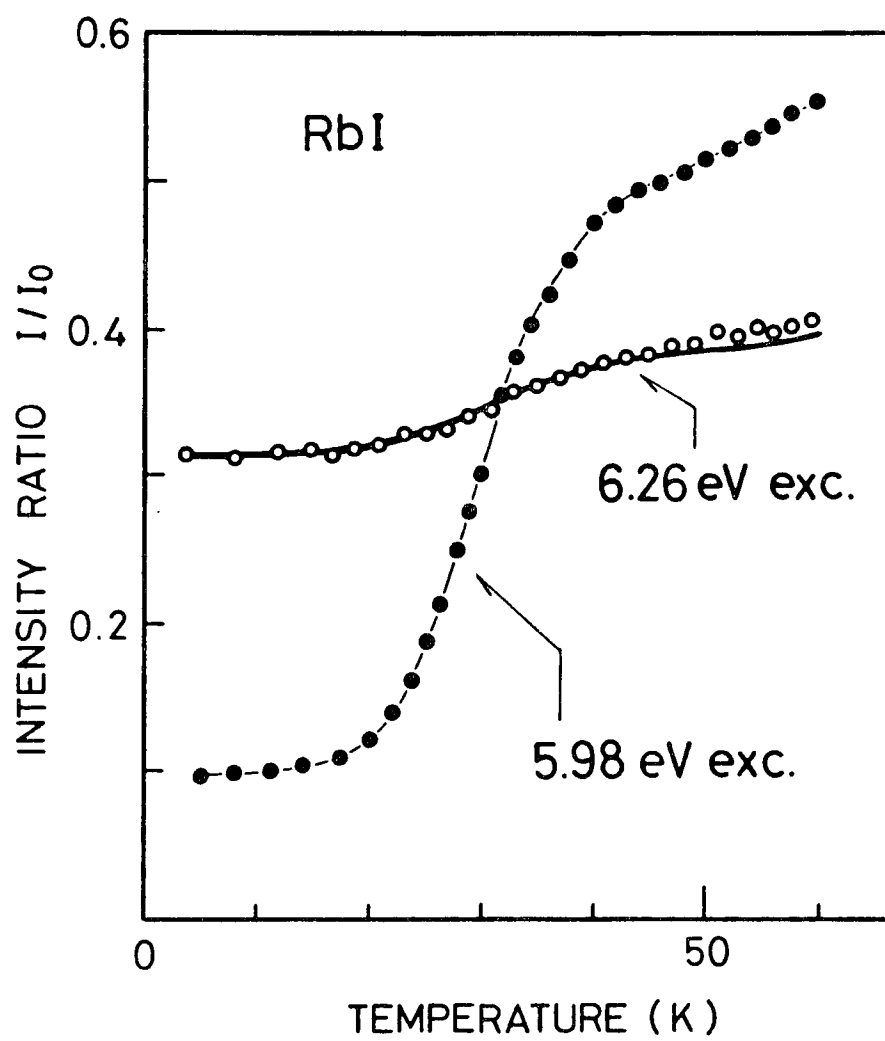


Fig. 5

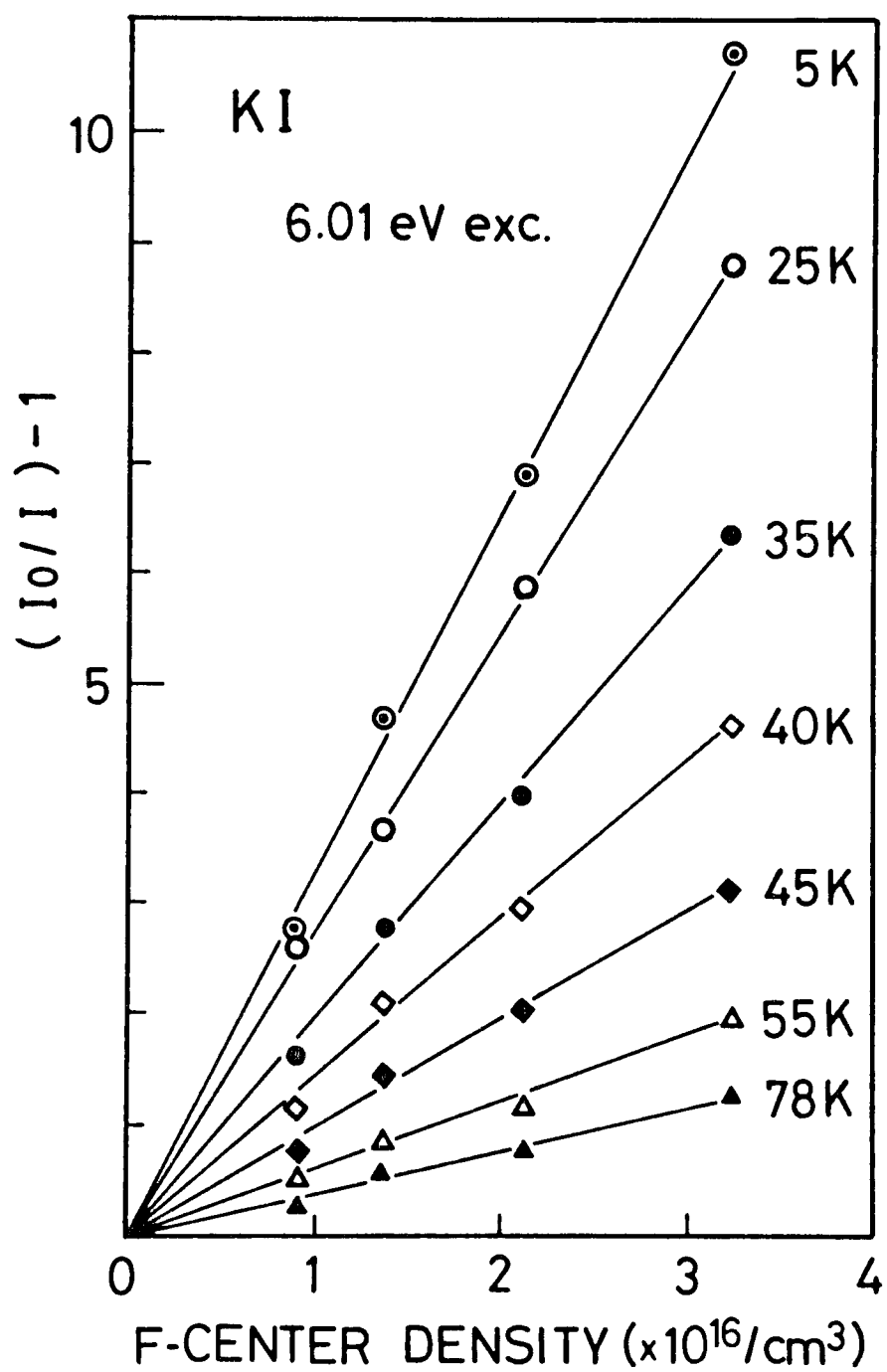


Fig. 6

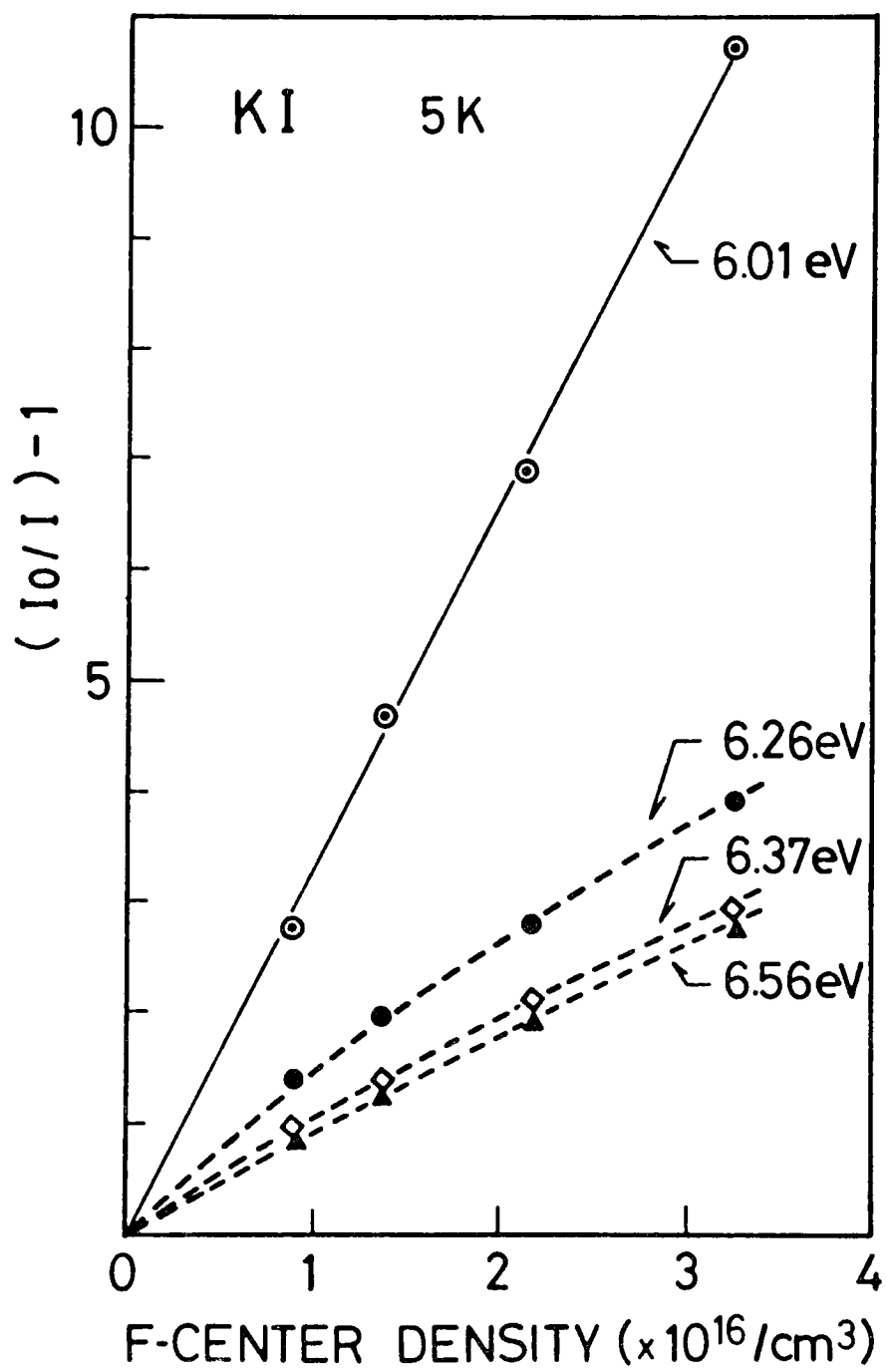


Fig. 7

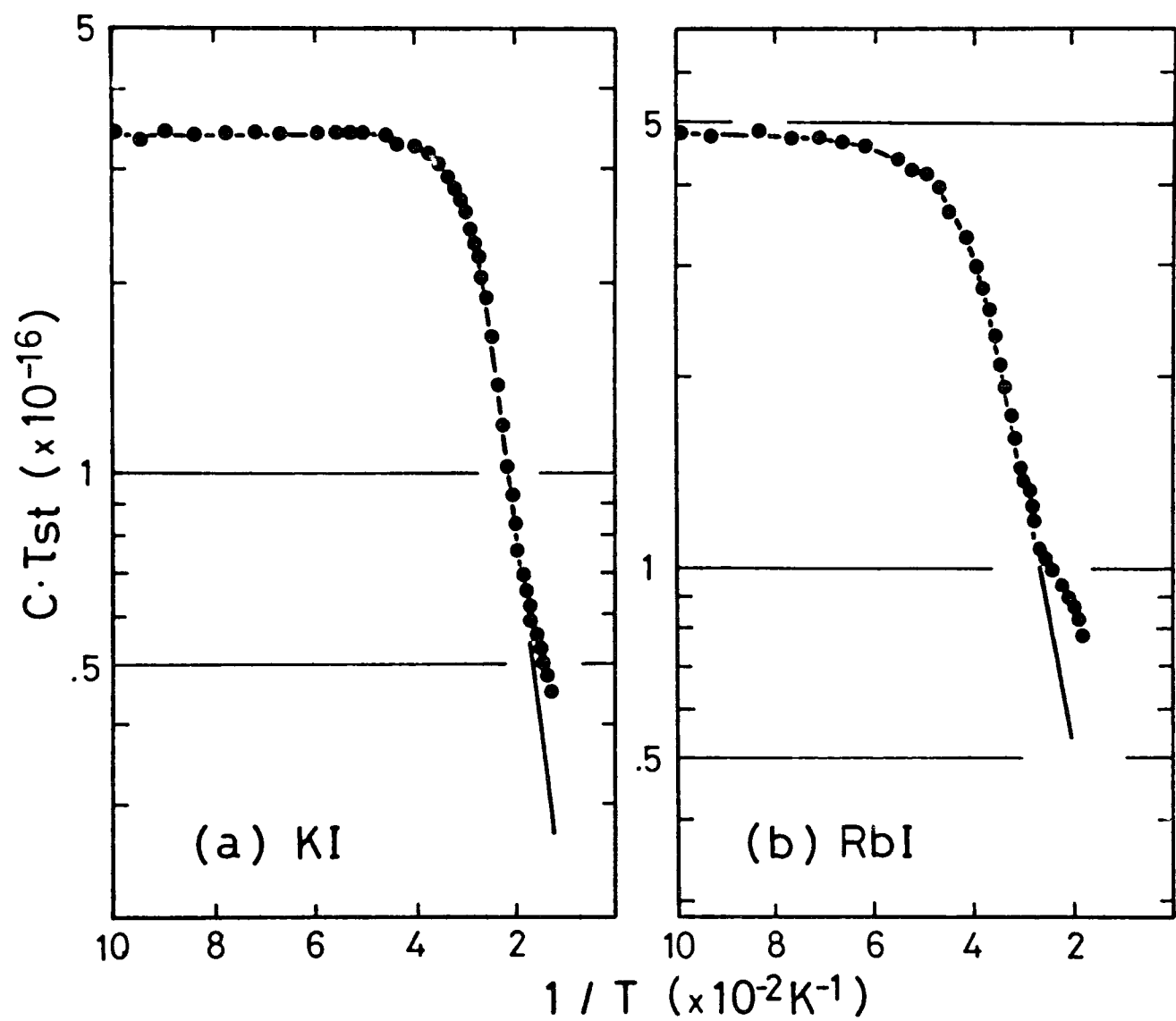


Fig. 8