<table>
<thead>
<tr>
<th>項目</th>
<th>内容</th>
</tr>
</thead>
<tbody>
<tr>
<td>タイトル</td>
<td>Dielectric abnormal relaxation in condensed matter (Session I: Cross-Disciplinary Physics, The 1st Tohwa University International Meeting on Statistical Physics Theories, Experiments and Computer Simulations)</td>
</tr>
<tr>
<td>著者</td>
<td>Mashimo, Satoru</td>
</tr>
<tr>
<td>引用</td>
<td>物性研究 (1996), 66(3): 372-375</td>
</tr>
<tr>
<td>発行日</td>
<td>1996-06-20</td>
</tr>
<tr>
<td>URL</td>
<td><a href="http://hdl.handle.net/2433/95844">http://hdl.handle.net/2433/95844</a></td>
</tr>
<tr>
<td>タイプ</td>
<td>Departmental Bulletin Paper</td>
</tr>
<tr>
<td>テキストバージョン</td>
<td>publisher</td>
</tr>
</tbody>
</table>

京都大学
Dielectric abnormal relaxation in condensed matter
Satoru Mashimo
Department of Physics, Tokai University
Hiratsuka-shi, Kanagawa 259-12, Japan

It has been shown by many experimental data of polymers that the dielectric response
function is of a stretched exponential, \( \phi(t) = \exp \left[ -\left( t/\tau^* \right)^{\beta} \right] \). The parameter \( \beta \) is
indicated to express a measure of the molecular interaction by the coupling model.
Experimental data on \( \beta \) support this. However deviation of the experimental complex
permittivity from that calculated is common for polymers. The stretched exponential function
explained the complex permittivity only in a frequency range \( \log f / \tau^* \leq 1.5 \). The deviation is in
many cases independent of the frequency. No satisfactorily explanation has been given to this
phenomenon. Below conventional glass transition temperature a relaxation peak of Debye
type can be seen in an ultra-low frequency region.

§ 1. Introduction

The theory of dielectric relaxation gives the complex permittivity \( \varepsilon^* \) as
\[
\frac{\varepsilon^*-\varepsilon_\infty}{\varepsilon_0-\varepsilon_\infty} = \int_0^\infty \exp(i\omega t) \left[ -\frac{d\phi(t)}{dt} \right] \, dt, \tag{1}
\]
where \( \varepsilon_0 \) and \( \varepsilon_\infty \) are the permittivity at the low- and high-frequency limits, respectively, and
\( \phi(t) \) is the dielectric response function. The dielectric relaxation has been discussed mostly
in terms of the complex permittivity. Empirical representations of the Cole-Cole and
Davidson-Cole have been used quite often for the discussion. Havriliak and Negami (HN)
presented the following equation which involves both representations:
\[
\frac{\varepsilon^*-\varepsilon_0}{\varepsilon_0-\varepsilon_\infty} = \frac{1}{1+\left[ (i\omega\tau_{HN})^\beta \right]^{\alpha}}, \quad 0<\alpha,\beta\leq1 \tag{2}
\]
where \( \tau_{HN} \) is a relaxation time, and \( \alpha \) and \( \beta \) are shape parameters dominating the frequency
dependence of the complex permittivity. This corresponds to the Cole-Cole equation when
\( \alpha=1 \) and the Davidson-Cole equation if \( \beta=1 \). It is established that Eq. (2) gives a good
explanation of the complex permittivity of polymers. In spite of the great success of Eq. (2),
the physical meanings of the parameters \( \alpha \) and \( \beta \) are still in despite.

Previously Williams and Watts found empirically a response function and it was later
pointed out that the function is the Kohlrausch frictional exponential function,
\[
\phi(t) = \exp \left[ -\left( t/\tau^* \right)^{\beta} \right], \quad 0<\beta,\tau^* \tag{3}
\]
\( \tau^* \) is a relaxation time.

Particularly interesting theories have been presented recently for the Kohlrausch
function. Ngai et al. have developed a coupling model for condensed matter, where
molecular orientation undergoes certainly an interaction with its complex environment and the
environment fluctuates during the orientation. This brings naturally a time-dependent rate
constant and $\beta_k$ is a measure of the interaction. Amorphous polymers are good examples to
examine the validity of the theory.

Bendler et al.\textsuperscript{14} have also derived the Kohlrausch function by applying the
Montroll-Weiss random walk to a defect diffusion model.\textsuperscript{15} It was shown that $\varepsilon'(\omega)$ given
by the Kohlrausch function is described naturally from the Levy distributions.

Palmer et al.\textsuperscript{16} have recently developed a hierarchically constrained 'dynamics for glassy
relaxation and derived the Kohlrausch function under a certain postulate for Ising spins. Over
a very wide time range of $\tau_0 < t < \tau_m$, where $\tau_0$ is a microscopic time ($\sim 10^{-14}$s) and $\tau_m$ is
the ergodic time, the response function is of the Kohlrausch form if temperature is not much
higher than $T_e$. However, it crosses over to pure exponential behavior of $\exp(-t/\tau_m)$ for
$t > \tau_m$. The Kohlrausch function is also derived by the model coupling theory.\textsuperscript{17}

§ 2. Dielectric relaxation described by the Kohlrausch function

Micro-Brownian motion of polymer chain is one of the most important subjects in
polymer physics. Relaxation phenomena such as viscoelastic relaxation, dielectric relaxation,
and NMR relaxation are dominated more or less by such motions. Recent studies on the
micro-Brownian motion showed that the response function is the Kohlrausch function for a
variety of polymers.

As an example, dielectric dispersion and absorption curves of poly(vinyl
acetate)(PVAc)\textsuperscript{19} are shown in Fig.1. Bath curves are well explained by Eqs. (1) and (4) in
frequency region of $f < 10^2 f_m$, where $f_m$ is the frequency of the maximal absorption. However in
the high frequency region of $f > 10^2 f_m$, both curves deviate definitely from the calculated ones.
In solution of PVAc in benzene too, the same trend can be seen as is shown in Fig.2. It is
quite interesting that the experimental observation is always larger than that calculated in the high
frequency region. The difference is almost independent of the frequency. If it comes
from another relaxation, such a relaxation is quite abnormal and cannot be interpreted at all.
This feature may be universal for all polymers. However it can be observed only by the
dielectric relaxation. Other relaxations such as viscoelastic relaxation have never shown such
an abnormal behavior observed by the dielectric measurement.

§ 3. Kohlrausch function and coupling model

The coupling model have shown that $\beta_k$ is a parameter representing a degree of
coupling strength brought from the molecular environment. Usually in amorphous polymer or
polymer solution, $\beta_k$ increases with temperature or decrease with the increase of polymer
concentration.\textsuperscript{10,17}

Aqueous solution of poly(vinyl methyl ether)(PVME) has a lower critical solution
temperature(LCST). This means that intra- and inter- chain interaction increases as the
temperature increases. In Fig.3, the parameter $\beta_k$ is plotted against temperature. It
decreases with the increase of temperature. This offers on evidence of the validity of the
coupling model.

The Kohlrausch function is essentially applicable to dielectric relaxation of polymer.
Nevertheless a sufficient explanation of the deviation in the high frequency region is required.

§ 4. Ergodic relaxation

Palmer et al.\textsuperscript{16} have suggested existence of pure exponential decay for glassy materials,
that is, the ergodic relaxation at very long times. The relaxation time $\tau_m$ estimated loosely
from the present result of PVAc at 26.85°C is $4.4 \times 10^2$ s and very far away from the Kohlrausch region of $t \leq 3 \times 10^4$ s. According to the theory, the current of the Kohlrausch form crosses over to pure exponential type at $t \sim \tau_m$. It should be noted that behavior of the Kohlrausch form at long times is not much different from the pure exponential behavior. Therefore it will be very difficult to distinguish these two decay.

Reference
Fig. 1. Comparison of dielectric dispersion and absorption curves obtained experimentally at 53°C with those calculated by the Kohlrausch function.

Fig. 2. Comparison of PVAc (50wt%) in benzene at 25°C with that calculated by the Kohlrausch function.

Fig. 3. Temperature dependence of coupling parameter $\beta_k$ for poly(vinyl methyl ether) (PVME) in aqueous solution.

Fig. 4. Discharge current of PVA at 26.85°C.
(1) Current calculated by the Kohlrausch function with $\beta_k = 0.470$, $\tau_0 = 2.53 \times 10^3 s$, and $\varepsilon_0 - \varepsilon_\infty = 7.44$.
(2) Current calculated by the Debye equation with $\tau_m = 4.4 \times 10^6 s$, and $\varepsilon_0 - \varepsilon_\infty = 0.51$. The discharge current is proportional to $-d \phi (t)/dt$. 