

Stationary martingale solution for the 2D stochastic Gross-Pitaevskii equation

By

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Abstract

In this short report we give a proof of the existence of a stationary solution to the Gross-Pitaevskii equation in $2d$ driven by a space-time white noise.

§ 1. Introduction

In this short report we give a proof of the existence of a stationary solution to the following Gross-Pitaevskii equation in $2d$ driven by a space-time white noise:

$$(1.1) \quad dX = (\gamma_1 + i\gamma_2)(HX - |X|^2X)dt + \sqrt{2\gamma_1}dW, \quad t > 0, \quad x \in \mathbb{R}^2,$$

where $H = \Delta - |x|^2$, $\gamma_1 > 0$, and $\gamma_2 \in \mathbb{R}$. The unknown function X is a complex valued random field on a probability space $(\Omega, \mathcal{F}, \mathbb{P})$ endowed with a standard filtration $(\mathcal{F}_t)_{t \geq 0}$. This equation is used as a model for Bose-Einstein condensates in the presence of temperature effects. There are some studies in the physics literature using this model [1, 6, 7, 14]. We are interested in this equation from a mathematical point of view, and we have studied the $1d$ case in [2] in particular the properties of the statistical equilibrium.

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As is the case for the stochastic quantization equations [4], the use of renormalization is necessary in order to give a meaning to the solutions of (1.1), as the Gaussian measure generated by the linear equation is only supported in $\mathcal{W}^{-s,q}$, with $s > 0$, $q \geq 2$, and $sq > 2$, where for $1 \leq p \leq +\infty$, $\sigma \in \mathbb{R}$,

$$\mathcal{W}^{\sigma,p}(\mathbb{R}^2) = \{v \in \mathcal{S}'(\mathbb{R}^2), |v|_{\mathcal{W}^{\sigma,p}(\mathbb{R}^2)} := |(-H)^{\sigma/2}v|_{L^p(\mathbb{R}^2)} < +\infty\},$$

denote the Sobolev space associated with the operator H .

Renormalization procedures, using Wick products, have been by now widely used in the context of singular stochastic partial differential equations, in particular for parabolic equations based on gradient flows (see for example [4, 13] for the $2d$ case). The complex Ginzburg-Landau equation driven by space-time white noise, i.e. (1.1) without the harmonic potential, posed on the three-dimensional torus, was studied in [8] and for the two-dimensional torus in [10, 12]. The main difference in our case is the presence of the harmonic potential $|x|^2$.

The proof in this report will not be published anywhere, nore in [3] where the existence of strong global solution for (1.1), i.e., much stronger result is established. In fact, the proof of this report was our first try for ensuring the existence of a solution to (1.1) for *any* dissipation parameter γ_1 : indeed, if γ_1 is sufficiently large we may use simply as in the purely parabolic case [13] an L^p energy estimate to globalize the solution. The small γ_1 case was later solved by using ideas inspired by the bootstrap arguments used in [10] (although the proof in [10] does not directly applies to the present case), and the proof for any γ_1 is written in [3]. However, we think it is interesting to present this first proof in this report, with some precisions about the dependence of estimates of stationary solutions on the dissipation parameters.

§ 2. notation and main result

In what follows, we will use the following notation: Let $\{h_k\}_{k \in \mathbb{N}^2}$ be the orthonormal basis of $L^2(\mathbb{R}^2, \mathbb{R})$, consisting of eigenfunctions of $-H$ with corresponding eigenvalues $\{\lambda_k^2\}_{k \in \mathbb{N}^2}$, i.e. $-Hh_k = \lambda_k^2 h_k$, $\lambda_k^2 = 2|k| + 2$. We take $\{h_k, ih_k\}_{k \in \mathbb{N}^2}$ as a complete orthonormal system in $L^2(\mathbb{R}^2, \mathbb{C})$, and we may write the cylindrical Wiener process in (1.1) as

$$(2.1) \quad W(t, x) = \sum_{k \in \mathbb{N}^2} (\beta_{k,R}(t) + i\beta_{k,I}(t))h_k(x).$$

Here, $(\beta_{k,R}(t))_{t \geq 0}$ and $(\beta_{k,I}(t))_{t \geq 0}$ are sequences of independent real-valued Brownian motions on the stochastic basis $(\Omega, \mathcal{F}, \mathbb{P}, (\mathcal{F}_t)_{t \geq 0})$. The notation $\mathbb{E}$ stands for the expectation with respect to $\mathbb{P}$.

We will use an approximation by finite dimensional objects. To verify the convergence properties of a function series of the form $u = \sum_{k \in \mathbb{N}^2} c_k h_k$, we define, for any $N \in \mathbb{N}$ fixed, for any $p \in [1, \infty]$, and $s \in \mathbb{R}$, a smooth projection operator $S_N : L^2(\mathbb{R}^2, \mathbb{C}) \rightarrow E_N^{\mathbb{C}} := \text{span}\{h_k\}_{|k| \leq N}$ by

$$(2.2) \quad S_N \left[\sum_{k \in \mathbb{N}^2} c_k h_k \right] := \sum_{k \in \mathbb{N}^2} \chi \left[\frac{\lambda_k^2}{\lambda_N^2} \right] c_k h_k = \chi \left[\frac{-H}{\lambda_N^2} \right] \left[\sum_{k \in \mathbb{N}^2} c_k h_k \right],$$

where $\chi \geq 0$ is a cut-off function such that $\chi \in C_0^\infty(-1, 1)$, $\chi = 1$ on $[-\frac{1}{2}, \frac{1}{2}]$. Note that here and in what follows, we denote by λ_N the value $\lambda_{(N,0)}$, for simplicity. The operator S_N , which is self-adjoint and commutes with H , may be extended by duality to any Sobolev space $\mathcal{W}^{s,2}(\mathbb{R}^2; \mathbb{C})$, with $s \in \mathbb{R}$, and thus by Sobolev embeddings, to any space $\mathcal{W}^{s,p}(\mathbb{R}^2; \mathbb{C})$, with $p \geq 1$. A simple modification of Theorem 1.1 of [9] implies that S_N is a bounded operator from L^p to L^p , uniformly in N , for any $p \in [1, \infty]$. We denote the usual spectral projector by

$$\Pi_N \left[\sum_{k \in \mathbb{N}^2} c_k h_k \right] := \sum_{k \in \mathbb{N}^2, |k| \leq N} c_k h_k,$$

which is uniformly bounded only in L^2 .

Let us recall known facts before mentioning precisely our results. Writing the solution of (1.1) as $X = u + Z_\infty^{\gamma_1, \gamma_2}$ with

$$(2.3) \quad Z_\infty^{\gamma_1, \gamma_2}(t) = \sqrt{2\gamma_1} \int_{-\infty}^t e^{(t-\tau)(\gamma_1 + i\gamma_2)H} dW(\tau),$$

which is the stationary solution for the linear stochastic equation

$$(2.4) \quad dZ = (\gamma_1 + i\gamma_2)H Z dt + \sqrt{2\gamma_1} dW,$$

we find out the following random partial differential equation for u :

$$(2.5) \quad \partial_t u = (\gamma_1 + i\gamma_2)(Hu - |u + Z_\infty^{\gamma_1, \gamma_2}|^2(u + Z_\infty^{\gamma_1, \gamma_2})), \quad u(0) = u_0 := X(0) - Z_\infty^{\gamma_1, \gamma_2}(0).$$

We are therefore required to solve this random partial differential equation. However, using standard arguments, it is not difficult to see that the best regularity we may expect for $Z_\infty^{\gamma_1, \gamma_2}$ is almost surely : $Z_\infty^{\gamma_1, \gamma_2} \in \mathcal{W}^{-s,q}(\mathbb{R}^2)$ for $s > 0$, $q \geq 2$, $sq > 2$ as follows.

Lemma 2.1. *Fix any $T > 0$. Let $\gamma_1 > 0, \gamma_2 \in \mathbb{R}, s > 0, q \geq 2, sq > 2$ and $0 < \alpha < \frac{1}{2}[(s - \frac{2}{q}) \wedge 1]$. The stationary solution $Z_\infty^{\gamma_1, \gamma_2}$ of (2.4) has a modification in $C^\alpha([0, T], \mathcal{W}^{-s,q})$. Moreover, there exists a positive constant C_T such that*

$$\mathbb{E} \left[\sup_{t \in [0, T]} |Z_\infty^{\gamma_1, \gamma_2}(t)|_{\mathcal{W}^{-s,q}} \right] \leq C_T.$$

Thus, to give a sense to the nonlinearity in (2.5), we need a renormalization and we will consider in place the renormalized equation of (2.5) :

(2.6)

$$\partial_t u = (\gamma_1 + i\gamma_2)(Hu - :|u + Z_\infty^{\gamma_1, \gamma_2}|^2(u + Z_\infty^{\gamma_1, \gamma_2}):), \quad u(0) = u_0 := X(0) - Z_\infty^{\gamma_1, \gamma_2}(0),$$

where $:$ $:$ in the nonlinear part, using the notation $Z_{R,\infty} = \text{Re}(Z_\infty^{\gamma_1, \gamma_2})$, $Z_{I,\infty} = \text{Im}(Z_\infty^{\gamma_1, \gamma_2})$, $u_R = \text{Re } u$, and $u_I = \text{Im } u$, means

$$(2.7) \quad :|u + Z_\infty^{\gamma_1, \gamma_2}|^2(u + Z_\infty^{\gamma_1, \gamma_2}): = F(u, (:Z_\infty^l:)_{1 \leq l \leq 3}) = F_0 + F_1 + F_2 + F_3$$

with $F_0 = |u|^2 u$, and

$$\begin{aligned} F_1 &= Z_\infty |u|^2 + 2Z_{R,\infty} u_R u + 2Z_{I,\infty} u_I u, \\ F_2 &=: Z_{R,\infty}^2 : (3u_R + iu_I) + : Z_{I,\infty}^2 : (u_R + 3iu_I) + 2 : Z_{R,\infty} Z_{I,\infty} : (u_I + iu_R), \\ F_3 &=: Z_{R,\infty}^3 : + i : Z_{I,\infty}^3 : + : Z_{R,\infty} Z_{I,\infty}^2 : + i : Z_{R,\infty}^2 Z_{I,\infty} :. \end{aligned}$$

Here, for any $k, l \in \mathbb{N}$, the Wick products $:(Z_{R,\infty})^k (Z_{I,\infty})^l:$ are defined as follows.

Recall that the Hermite polynomials $H_n(x)$, $n \in \mathbb{N}$ are defined by

$$(2.8) \quad H_n(x) = \frac{(-1)^n}{\sqrt{n!}} e^{\frac{x^2}{2}} \frac{d^n}{dx^n} (e^{-\frac{x^2}{2}}), \quad n \geq 1$$

and $H_0(x) = 1$.

The notation $:(S_N z)^n:(x)$ for $n \in \mathbb{N}$, $N \in \mathbb{N}$, $x \in \mathbb{R}^2$, with a real-valued centered Gaussian white noise z , means

$$:(S_N z)^n:(x) = \rho_N(x)^n \sqrt{n!} H_n \left[\frac{1}{\rho_N(x)} S_N z(x) \right], \quad x \in \mathbb{R}^2$$

with

$$\rho_N(x) = \left[\sum_{k \in \mathbb{N}^2} \chi^2 \left(\frac{\lambda_k^2}{\lambda_N^2} \right) \frac{1}{\lambda_k^2} (h_k(x))^2 \right]^{\frac{1}{2}}.$$

It is known (see [2]) that the law $\mathcal{L}(Z_\infty^{\gamma_1, \gamma_2})$ equals the complex Gaussian measure $\mu = \mathcal{N}_{\mathbb{C}}(0, 2(-H)^{-1})$.

Proposition 2.2. ([3]) *For any $k, l \in \mathbb{N}$, the sequence $\{:(S_N Z_{R,\infty})^k :: (S_N Z_{I,\infty})^l:\}_{N \in \mathbb{N}}$ is a Cauchy sequence in $L^q(\Omega, \mathcal{W}^{-s,q}(\mathbb{R}^2))$, for $q > 2$, $s > 0$ with $qs > 2$.*

Moreover, defining then, for any $k, l \in \mathbb{N}$, for any fixed t ,

$$:(Z_{R,\infty})^k (Z_{I,\infty})^l: := \lim_{N \rightarrow \infty} : (S_N Z_{R,\infty})^k :: (S_N Z_{I,\infty})^l :, \quad \text{in } L^q(\Omega, \mathcal{W}^{-s,q}(\mathbb{R}^2)),$$

where $s > 0, q > 2$ and $sq > 2$, there exists a constant $M_{s,q,k,l}$ such that

$$(2.9) \quad \mathbb{E} [: (Z_{R,\infty})^k (Z_{I,\infty})^l : |_{\mathcal{W}^{-s,q}}^q] \leq M_{s,q,k,l}.$$

Remark that higher order moments may also be estimated thanks to Nelson formula: Let $s > 0$, $m > q > 2$ and $sq > 2$. Then there is a constant $M_{s,q,k,l,m}$ such that

$$(2.10) \quad \mathbb{E} [| : (Z_{R,\infty})^k (Z_{I,\infty})^l : |_{\mathcal{W}^{-s,q}}^m] \leq M_{s,q,k,l,m}.$$

Note that if we consider the equation (2.6) in terms of X , then

$$(2.11) \quad dX = (\gamma_1 + i\gamma_2)(HX - :|X|^2 X :)dt + \sqrt{2\gamma_1}dW, \quad t > 0, \quad x \in \mathbb{R}^2.$$

In [3], we constructed a measure ρ as a weak limit of the family of finite dimensional Gibbs measure of the form :

$$d\tilde{\rho}_N(y) = \Gamma_N e^{-\tilde{\mathcal{H}}_N(S_N y)} dy, \quad y \in E_N^{\mathbb{C}},$$

where $\Gamma_N^{-1} = \int e^{-\tilde{\mathcal{H}}_N(S_N y)} dy$, and

$$\tilde{\mathcal{H}}_N(y) = \frac{1}{2}|\nabla y|_{L^2}^2 + \frac{1}{2}|xy|_{L^2}^2 + \int_{\mathbb{R}^2} \left[\frac{1}{4}|y(x)|^4 - 2\rho_N^2(x)|y(x)|^2 + 2\rho_N^4(x) \right] dx.$$

Note that

$$\tilde{\mathcal{H}}_N(y) = \frac{1}{2}|\nabla y|_{L^2}^2 + \frac{1}{2}|xy|_{L^2}^2 + \frac{1}{4} \int_{\mathbb{R}^2} :|y(x)|^4 : dx,$$

and $\nabla_y \tilde{\mathcal{H}}_N(y) = -Hy + :|y|^2 y :$.

Proposition 2.3. ([3]) *The family of finite dimensional Gibbs measures $(\tilde{\rho}_N)_N$ is tight in $\mathcal{W}^{-s,q}$ for any $s > 0$, $q \geq 2$ and $sq > 2$.*

Note that $\tilde{\rho}_N$ does not depend on γ_1 or γ_2 . It is an invariant measure for the case of $\gamma_1 = 0$ and $\gamma_2 \neq 0$ or for the case of $\gamma_1 > 0$ and $\gamma_2 = 0$, for the case of $\gamma_1 > 0$ and $\gamma_2 \neq 0$ to the finite dimensional equation:

$$(2.12) \quad dX = (\gamma_1 + i\gamma_2)(HX - S_N(:|S_N X|^2 S_N X :))dt + \sqrt{2\gamma_1}\Pi_N dW, \quad t > 0, \quad x \in \mathbb{R}^2, \quad X(0) \in E_N^{\mathbb{C}}.$$

The global existence of the solution X_N of this finite dimensional equation is ensured by the standard fixed point methods, and an energy estimate as in [2] (see also (2.16) below). Remark that the invariant measure $\tilde{\rho}_N$ is unique if $\gamma_1 > 0$ and $\gamma_2 = 0$. Eq.(2.12) gives the Galerkin approximation of X satisfying (2.11).

Finally we deduce the following main result for the equation (2.11):

Theorem 2.4. *Let $\gamma_1 > 0$ and $\gamma_2 \in \mathbb{R}$, and let $0 < s < 1$, $q > 2$ such that $sq > 2$. Then there exists a stationary martingale solution X of (2.11) having trajectories in $C(\mathbb{R}_+, \mathcal{W}^{-s,q})$, and $\mathcal{L}(X(t)) = \rho$ for all $t \in \mathbb{R}$.*

The proof of Theorem 2.4 will be given in Section 3.

The tightness of the family of measures $(\tilde{\rho}_N)_N$ in Proposition 2.3 was proved in [3] considering the coupled evolution on $E_N^{\mathbb{C}}$ given by

$$(2.13) \quad \begin{cases} \frac{du}{dt} = (\gamma_1 + i\gamma_2) [Hu - S_N(\cdot : |S_N(u+Z)|^2 S_N(u+Z) :)] \\ dZ = (\gamma_1 + i\gamma_2) HZ dt + \sqrt{2\gamma_1} \Pi_N dW, \end{cases}$$

one may easily prove, using e.g. similar estimates as in the proof of Proposition 2.5 below, together with the Gaussianity of Z , and a Krylov-Bogolyubov argument, that (2.13) has an invariant measure ν_N on $E_N^{\mathbb{C}} \times \mathbb{E}_N^{\mathbb{C}}$. Moreover, by uniqueness of the invariant measure of (2.12) in case of $\gamma_2 = 0$, we necessarily have for any bounded continuous function φ on $\mathbb{E}_N^{\mathbb{C}}$:

$$\int_{E_N^{\mathbb{C}}} \varphi(x) \tilde{\rho}_N(dx) = \iint_{E_N^{\mathbb{C}} \times E_N^{\mathbb{C}}} \varphi(u+z) \nu_N(du, dz).$$

Proposition 2.5. *Let $\gamma_1 > 0$ and $\gamma_2 \in \mathbb{R}$. Let $(u_N, Z_N) \in C(\mathbb{R}_+; E_N^{\mathbb{C}} \times E_N^{\mathbb{C}})$ be a stationary solution of (2.13). Then, for any $m > 0$, there is a constant $C_{m,\gamma_1,\gamma_2} > 0$ independent of t and N , such that*

$$(2.14) \quad \mathbb{E}(|(-H)^{\frac{1}{2m}} u_N|_{L^2}^{2m}) \leq C_{m,\gamma_1,\gamma_2}.$$

Remark.

- (1) The constant C_{m,γ_1,γ_2} in the RHS does not depend on γ_1 when $\gamma_2 = 0$, and if $\gamma_2 \neq 0$, then C_{m,γ_1,γ_2} depends on the ratio $\frac{|\gamma_2|}{\gamma_1}$. Let s with $0 < s < 1$ and $q > 2$ such that $sq > 2$. Using this remark, applying Proposition 2.5 with $m = 1$ and $\gamma_2 = 0$, we deduce that for some positive constant C not depending on N, γ_1 , and for any $t \geq 0$,

$$\mathbb{E}(|u_N(t)|_{\mathcal{W}^{-s,q}}^2) \lesssim \mathbb{E}(|u_N(t)|_{L^q}^2) \lesssim \mathbb{E}(|(-H)^{\frac{1}{2}} u_N|_{L^2}^2) \lesssim C,$$

where we have used the embedding $\mathcal{W}^{1,2} \subset L^q$, for any $q < +\infty$. Thus,

$$\begin{aligned} \int_{\mathcal{W}^{-s,q}} |x|_{\mathcal{W}^{-s,q}}^2 \tilde{\rho}_N(dx) &= \iint_{(\mathcal{W}^{-s,q})^2} |u+z|_{\mathcal{W}^{-s,q}}^2 \nu_N(du, dz) \\ &\leq 2\mathbb{E}(|u_N(t)|_{\mathcal{W}^{-s,q}}^2 + |Z_N(t)|_{\mathcal{W}^{-s,q}}^2), \end{aligned}$$

and the right hand side above is bounded indepently of N , γ_1 and t , since the law of Z_N converges to a Gaussian measure μ (which is independent of γ_1) on $\mathcal{W}^{-s,q}$. We see thus that the tightness is independent of the parameters γ_1, γ_2 , which is consistent with the fact that $\tilde{\rho}_N$ is independent of γ_1 and γ_2 .

- (2) Unfortunately, the bound (2.14) does not provide higher moment bounds on the measures $\tilde{\rho}_N$, preventing us to obtain ρ -a.s. initial data global existence of solution by the method in [5].

For the proof of Proposition 2.5, we will use the following interpolation estimates:
Let $\alpha \geq 0$.

$$(2.15) \quad |fg|_{\mathcal{W}^{\alpha,q}} \leq C(|f|_{L^{q_1}} |g|_{\mathcal{W}^{\alpha,\bar{q}_1}} + |f|_{\mathcal{W}^{\alpha,q_2}} |g|_{L^{\bar{q}_2}}),$$

where $1 < q < \infty$, $q_1, q_2 \in (1, \infty]$, $\bar{q}_1, \bar{q}_2 \in [1, \infty)$ with $\frac{1}{q} = \frac{1}{q_1} + \frac{1}{q_2} = \frac{1}{\bar{q}_1} + \frac{1}{\bar{q}_2}$.

Proof of Proposition 2.5. Taking the L^2 -inner product of the first equation in (2.13) with u_N yields

$$(2.16) \quad \begin{aligned} \frac{1}{2} \frac{d}{dt} |u_N(t)|_{L^2}^2 + \gamma_1 |(-H)^{\frac{1}{2}} u_N(t)|_{L^2}^2 + \gamma_1 |S_N u_N(t)|_{L^4}^4 = -\operatorname{Re}(\gamma_1 + i\gamma_2) \int_{\mathbb{R}^2} [F_1(S_N u_N, S_N Z_N) \\ + F_2(S_N u_N, S_N Z_N) + F_3(S_N Z_N)] \overline{S_N u_N}(t) dx. \end{aligned}$$

We first estimate the term containing F_3 in the right hand side above. Thanks to Proposition 2.2, taking $0 < s < 1$ and $q > 2$ such that $sq > 2$, we may bound

$$\left| \int_{\mathbb{R}^2} F_3(S_N Z_N) \overline{S_N u_N} dx \right| \lesssim |F_3(S_N Z_N)|_{\mathcal{W}^{-s,q}} |S_N u_N|_{\mathcal{W}^{s,q'}}$$

with $\frac{1}{q} + \frac{1}{q'} = 1$. Interpolating then $\mathcal{W}^{s,q'}$ between L^r and $\mathcal{W}^{1,2}$, with $\frac{1}{q'} = \frac{s}{2} + \frac{1-s}{r}$, we get

$$|S_N u_N|_{\mathcal{W}^{s,q'}} \lesssim |(-H)^{\frac{1}{2}} S_N u_N|_{L^2}^s |S_N u_N|_{L^r}^{1-s}.$$

On the other hand, noticing that $r \in (1, 2)$, we have for any $v \in \mathcal{W}^{1,2}$:

$$\begin{aligned} \int_{|x| \geq 1} |v(x)|^r dx &\leq \left[\int_{|x| \geq 1} |x|^2 |v(x)|^2 dx \right]^{\frac{r}{2}} \left[\int_{|x| \geq 1} |x|^{-\frac{2r}{2-r}} dx \right]^{\frac{2-r}{2r}} \\ &\lesssim |(-H)^{\frac{1}{2}} v|_{L^2}^r, \end{aligned}$$

so that $|v|_{L^r} \lesssim |(-H)^{\frac{1}{2}} v|_{L^2}$. It follows that

$$(2.17) \quad \begin{aligned} \left| \int_{\mathbb{R}^2} F_3(S_N Z_N) \overline{S_N u_N} dx \right| &\lesssim |F_3(S_N Z_N)|_{\mathcal{W}^{-s,q}} |(-H)^{\frac{1}{2}} S_N u_N|_{L^2} \\ &\leq \frac{4(\gamma_1 + |\gamma_2|)}{\gamma_1} |F_3(S_N Z_N)|_{\mathcal{W}^{-s,q}}^2 \\ &\quad + \frac{\gamma_1}{4(\gamma_1 + |\gamma_2|)} |(-H)^{\frac{1}{2}} S_N u_N|_{L^2}^2. \end{aligned}$$

Next, we consider the term containing F_2 in the right hand side of (2.16). First, by (2.7) and (2.15),

$$\begin{aligned} \left| \int_{\mathbb{R}^2} F_2(S_N u_N, S_N Z_N) \overline{S_N u_N} dx \right| &\lesssim \sum_{l=0}^2 | : S_N Z_{N,R}^l S_N Z_{N,I}^{2-l} : |_{\mathcal{W}^{-s,q}} |(S_N u_N)^2|_{\mathcal{W}^{s,q'}} \\ &\lesssim \sum_{l=0}^2 | : S_N Z_{N,R}^l S_N Z_{N,I}^{2-l} : |_{\mathcal{W}^{-s,q}} |S_N u_N|_{L^4} |S_N u_N|_{\mathcal{W}^{s,p}} \end{aligned}$$

with $\frac{1}{q'} = \frac{1}{4} + \frac{1}{p}$. Note that $p \in (1, 2)$ and we may use the same procedure as before to obtain

$$|S_N u_N|_{\mathcal{W}^{s,p}} \lesssim |(-H)^{\frac{1}{2}} S_N u_N|_{L^2},$$

so that

$$\begin{aligned} &\left| \int_{\mathbb{R}^2} F_2(S_N u_N, S_N Z_N) \overline{S_N u_N} dx \right| \\ (2.18) &\leq \frac{1}{2} \left(\frac{4(\gamma_1 + |\gamma_2|)}{\gamma_1} \right)^3 \sum_{l=0}^2 | : S_N Z_{N,R}^l S_N Z_{N,I}^{2-l} : |_{\mathcal{W}^{-s,q}}^4 + \frac{\gamma_1}{2(\gamma_1 + |\gamma_2|)} |S_N u_N|_{L^4}^4 \\ &\quad + \frac{\gamma_1}{4(\gamma_1 + |\gamma_2|)} |(-H)^{\frac{1}{2}} S_N u_N|_{L^2}^2. \end{aligned}$$

We finally turn to the term containing F_1 in the right hand side of (2.16). We easily get, thanks again to (2.15),

$$\left| \int_{\mathbb{R}^2} F_1(S_N u_N, S_N Z_N) \overline{S_N u_N} dx \right| \lesssim |S_N Z_N|_{\mathcal{W}^{-s,q}} |S_N u_N|_{L^4}^2 |S_N u_N|_{\mathcal{W}^{s,r}}$$

where $r > 2$ is such that $\frac{1}{q'} = \frac{1}{2} + \frac{1}{r}$. Let $m > 2$ with $\frac{1}{r} = \frac{s}{2} + \frac{1-s}{m}$, so that

$$|S_N u_N|_{\mathcal{W}^{s,r}} \lesssim |(-H)^{\frac{1}{2}} S_N u_N|_{L^2}^s |S_N u_N|_{L^m}^{1-s}.$$

If $m > 4$, we interpolate L^m between L^4 and L^{2m} , then use the Sobolev embedding $\mathcal{W}^{1,2} \subset L^{2m}$. If $2 < m \leq 4$, we interpolate L^m between L^2 and L^4 , then use $\mathcal{W}^{1,2} \subset L^2$. In both cases we obtain, using in addition the Poincaré inequality for $(-H)$:

$$|S_N u_N|_{\mathcal{W}^{s,r}} \lesssim |(-H)^{\frac{1}{2}} S_N u_N|_{L^2}^\alpha |S_N u_N|_{L^4}^{1-\alpha}$$

for some constant $\alpha \in (0, 1)$. We deduce that

$$\begin{aligned} &\left| \int_{\mathbb{R}^2} F_1(S_N u_N, S_N Z_N) \overline{S_N u_N} dx \right| \\ &\lesssim |S_N Z_N|_{\mathcal{W}^{-s,q}} |S_N u_N|_{L^4}^{3-\alpha} |(-H)^{\frac{1}{2}} S_N u_N|_{L^2}^\alpha \\ &\leq 2^{-\frac{3-\alpha}{1-\alpha}} \left(\frac{4(\gamma_1 + |\gamma_2|)}{\gamma_1} \right)^{\frac{1}{2-\alpha}(\alpha + \frac{2(3-\alpha)}{1-\alpha})} |S_N Z_N|_{\mathcal{W}^{-s,q}}^{\frac{4}{1-\alpha}} \\ (2.19) &\quad + \frac{\gamma_1}{2(\gamma_1 + |\gamma_2|)} \left[|S_N u_N|_{L^4}^4 + \frac{1}{2} |(-H)^{\frac{1}{2}} S_N u_N|_{L^2}^2 \right], \end{aligned}$$

by Young inequality.

Gathering (2.16)–(2.19), and noticing that $|(-H)^{\frac{1}{2}} S_N u_N|_{L^2} \leq |(-H)^{\frac{1}{2}} u_N|_{L^2}$, leads to

$$\frac{d}{dt} |u_N(t)|_{L^2}^2 + \frac{\gamma_1}{2} |(-H)^{\frac{1}{2}} u_N(t)|_{L^2}^2 \lesssim C_0 \sum_{k+l=1}^3 | : S_N Z_{N,R}^l S_N Z_{N,I}^k : |_{\mathcal{W}^{-s,q}}^{m_{k,l}},$$

with

$$C_0 \leq \gamma_1 \left(\frac{\gamma_1}{4(\gamma_1 + |\gamma_2|)} \right)^{-\kappa}$$

for some integers $m_{k,l}$ and $\kappa = \kappa_\alpha > 0$.

Now, let $m > 0$. We multiply by $|u_N(t)|_{L^2}^{2m-2}$ both sides of the above inequality to get

$$\frac{1}{m} \frac{d}{dt} |u_N(t)|_{L^2}^{2m} + \frac{\gamma_1}{2} |(-H)^{\frac{1}{2}} u_N(t)|_{L^2}^2 |u_N(t)|_{L^2}^{2m-2} \lesssim C_0 |u_N(t)|_{L^2}^{2m-2} \sum_{k+l=1}^3 | : S_N Z_{N,R}^l S_N Z_{N,I}^k : |_{\mathcal{W}^{-s,q}}^{m_{k,l}}.$$

Applying the interpolation inequality

$$|(-H)^{\frac{1}{2m}} u_N|_{L^2} \leq C_m |(-H)^{\frac{1}{2}} u_N|_{L^2}^{\frac{1}{m}} |u_N|_{L^2}^{\frac{m-1}{m}}$$

to the second term in the left hand side, and using Young inequality in the right hand side, we obtain

$$\frac{1}{m} \frac{d}{dt} |u_N(t)|_{L^2}^{2m} + \frac{\gamma_1}{2C_m^{2m}} |(-H)^{\frac{1}{2m}} u_N(t)|_{L^2}^{2m} \lesssim C_0 \varepsilon |u_N(t)|_{L^2}^{2m} + C_0 C_\varepsilon \sum_{k+l=1}^3 | : S_N Z_{N,R}^l S_N Z_{N,I}^k : |_{\mathcal{W}^{-s,q}}^{m'_{k,l}},$$

for any $\varepsilon > 0$ and constants $C_\varepsilon > 0$ and $m'_{k,l} > 0$. We choose $C_0 \varepsilon = \frac{\gamma_1 \lambda_1^2}{4C_m^{2m}}$ after using Poincaré inequality so that the first term of the right hand side is absorbed in the left hand side,

$$\frac{1}{m} \frac{d}{dt} |u_N(t)|_{L^2}^{2m} + \frac{\gamma_1}{4C_m^{2m}} |(-H)^{\frac{1}{2m}} u_N(t)|_{L^2}^{2m} \lesssim C_0 C_\varepsilon \sum_{k+l=1}^3 | : S_N Z_{N,R}^l S_N Z_{N,I}^k : |_{\mathcal{W}^{-s,q}}^{m'_{k,l}}.$$

Integrating in time, taking expectations on both sides and using the stationarity of u_N and of the Wick products, together with (2.10), yields

$$\mathbb{E}(|(-H)^{\frac{1}{2m}} u_N|_{L^2}^{2m}) \lesssim_m \left(\frac{\gamma_1}{4(\gamma_1 + |\gamma_2|)} \right)^{-\kappa'} \sum_{k+l=1}^3 M_{s,q,k,l,m'_{k,l}},$$

for some $\kappa' = \kappa'_\alpha > 0$ and the conclusion. $\square$

§ 3. Proof of Theorem 2.4

We are in position to construct a stationary solution of (2.11) for any values of $\gamma_1 > 0$ and $\gamma_2 \in \mathbb{R}$. We have seen in the previous section that the system (2.13) has a stationary solution (u_N, Z_N) where $Z_N = \Pi_N Z_\infty^{\gamma_1, \gamma_2}$. Moreover, it is clear that $(u_N, Z_\infty^{\gamma_1, \gamma_2})$ is then a stationary solution of

$$(3.1) \quad \begin{cases} \frac{du}{dt} = (\gamma_1 + i\gamma_2) [Hu - S_N(|S_N(u+Z)|^2 S_N(u+Z))] \\ dZ = (\gamma_1 + i\gamma_2) HZ dt + \sqrt{2\gamma_1} dW. \end{cases}$$

In this section we will denote $Z_\infty^{\gamma_1, \gamma_2}$ by Z for the sake of simplicity. Using Proposition 2.5, we first prove that the law of this sequence $\{(u_N, Z)\}_{N \in \mathbb{N}}$ is tight in an appropriate space to construct a martingale solution.

We will use the following lemma (see [3] for the proof).

Lemma 3.1. *Let $1 < p < q < +\infty$, $0 < s < \beta < 2/p$, and $m \in \mathbb{N} \setminus \{0\}$. Suppose $\beta - s - (m-1)(\frac{2}{p} - \beta) > 0$, and $s + m(\frac{2}{p} - \beta) < 2(1 - \frac{1}{q})$. Then, there is a constant $C > 0$ such that*

$$|hf^m|_{\mathcal{W}^{-(s+m(\frac{2}{p}-\beta)),q}} \leq C|h|_{\mathcal{W}^{-s,q}}|f|_{\mathcal{W}^{\beta,p}}^m.$$

Lemma 3.2. *Let $\gamma_1 > 0$, $\gamma_2 \in \mathbb{R}$, $0 < s < 1$, and $q > 2$ such that $qs > 2$. Let also $0 < \delta < \frac{1}{6}$, and let $p > \max\{q, 24, \frac{2}{\delta}\}$. The sequence $(u_N, Z)_{N \in \mathbb{N}}$ is bounded in*

$$L^{2m}(\Omega, L^{2m}(0, T, H^{\frac{1}{m}})) \cap L^{\frac{4}{3}}(\Omega, \mathcal{W}^{1, \frac{4}{3}}(0, T, \mathcal{W}^{-2,p})) \times C^\alpha([0, T], \mathcal{W}^{-s,q} \cap \mathcal{W}^{-\delta,p})$$

for any $m > 0$, and $\alpha > 0$ satisfying $\alpha < \min(\frac{s}{2} - \frac{1}{q}, \frac{\delta}{2} - \frac{1}{p}, 1)$.

Proof. It suffices to check the bound in $L^{\frac{4}{3}}(\Omega, \mathcal{W}^{1, \frac{4}{3}}(0, T, \mathcal{W}^{-2,p}))$, since the other bounds follow from Proposition 2.5, the stationarity of u_N , and Lemma 2.1 for Z . We write the equation for u_N ;

$$\begin{aligned} u_N(t) &= u_N(0) + (\gamma_1 + i\gamma_2) \int_0^t H u_N(\sigma) d\sigma \\ &\quad - (\gamma_1 + i\gamma_2) \int_0^t S_N(|S_N(u_N + Z_N)|^2 S_N(u_N + Z_N))(\sigma) d\sigma. \end{aligned}$$

In the right hand side, the first term is constant and clearly bounded in $L^{\frac{4}{3}}(\Omega, \mathcal{W}^{1, \frac{4}{3}}(0, T, \mathcal{W}^{-2,p}))$ by Proposition 2.5 and Hölder inequality. For the second term, we have

$$\begin{aligned} \mathbb{E} \left(\left| \int_0^\cdot H u_N(\sigma) d\sigma \right|_{\mathcal{W}^{1, \frac{4}{3}}(0, T, \mathcal{W}^{-2,p})}^{\frac{4}{3}} \right) &\leq C_T \mathbb{E} \left(|H u_N|_{L^{\frac{4}{3}}(0, T, \mathcal{W}^{-2,p})}^{\frac{4}{3}} \right) \\ &\leq C_T \mathbb{E} \left(|u_N|_{L^2(0, T, \mathcal{W}^{1,2})}^2 \right), \end{aligned}$$

where we have used Sobolev embedding and Hölder inequality in the last inequality. To estimate the nonlinear terms, we decompose as in (2.7),

$$\begin{aligned} S_N(\cdot |S_N(u_N + Z)|^2 S_N(u_N + Z) \cdot) &= S_N(F_0(S_N u_N) + F_1(S_N u_N, S_N Z) \\ &\quad + F_2(S_N u_N, S_N Z) + F_3(S_N Z)). \end{aligned}$$

The terms in F_3 are simply estimated thanks to Proposition 2.2 as follows.

$$\mathbb{E} \left(\left| \int_0^\cdot S_N F_3(s) ds \right|_{\mathcal{W}^{1, \frac{4}{3}}(0, T, \mathcal{W}^{-2, p})}^{\frac{4}{3}} \right) \leq C_T \mathbb{E}(|S_N F_3|_{L^{\frac{4}{3}}(0, T, \mathcal{W}^{-s, p})}^{\frac{4}{3}}) \leq C_T M_{s, p}^{\frac{4}{3}}.$$

For the term F_0 , using Sobolev embeddings $L^{\frac{4}{3}} \subset \mathcal{W}^{-2, p}$ and $\mathcal{W}^{\frac{1}{2}, 2} \subset L^4$, we obtain

$$\begin{aligned} \mathbb{E}(|S_N F_0|_{L^{\frac{4}{3}}(0, T, \mathcal{W}^{-2, p})}^{\frac{4}{3}}) &= \int_0^T \mathbb{E}(|S_N u_N|^2 S_N u_N|_{\mathcal{W}^{-2, p}}^{\frac{4}{3}}) ds \\ &\leq C_T \mathbb{E}(|S_N u_N|^2 S_N u_N|_{L^{\frac{4}{3}}}^{\frac{4}{3}}) = C_T \mathbb{E}(|S_N u_N|_{L^4}^4) \leq C_T \mathbb{E}(|u_N|_{\mathcal{W}^{\frac{1}{2}, 2}}^4), \end{aligned}$$

which is bounded independently of N by Proposition 2.5. To estimate the F_1 -terms, we fix $s' > 0$ such that $s' < \frac{1}{12}$ and $s'p > 2$, and apply Lemma 3.1 to get

$$\mathbb{E} \left(|S_N F_1|_{L^{\frac{4}{3}}(0, T, \mathcal{W}^{-2, p})}^{\frac{4}{3}} \right) \leq C_T \mathbb{E} \left(|F_1|_{\mathcal{W}^{-(s' + \frac{1}{24}), p}}^{\frac{4}{3}} \right) \leq C_T \mathbb{E} \left(|S_N Z|_{\mathcal{W}^{-s', p}}^{\frac{4}{3}} |S_N u_N|_{\mathcal{W}^{\frac{3}{8}, 3}}^{\frac{8}{3}} \right).$$

Using then the Sobolev embedding $\mathcal{W}^{\frac{17}{24}, 2} \subset \mathcal{W}^{\frac{3}{8}, 3}$ and Hölder inequality, the right hand side is majorized by

$$C_T \mathbb{E}(|S_N Z|_{\mathcal{W}^{-s', p}}^{24})^{\frac{1}{18}} \mathbb{E}(|S_N u_N|_{\mathcal{W}^{\frac{17}{24}, 2}}^{\frac{48}{17}})^{\frac{17}{18}},$$

and is thus bounded independently of N thanks to Propositions 2.2 and 2.5. Finally, for the terms in F_2 , we apply again Lemma 3.1:

$$\begin{aligned} \mathbb{E} \left(|S_N F_2|_{L^{\frac{4}{3}}(0, T, \mathcal{W}^{-2, p})}^{\frac{4}{3}} \right) &\leq T \mathbb{E} \left(|F_2|_{\mathcal{W}^{-(s' + \frac{1}{6}), p}}^{\frac{4}{3}} \right) \\ &\leq C_T \mathbb{E} \left(\sum_{k+l=2} | : (S_N Z_R)^k (S_N Z_I)^l : |_{\mathcal{W}^{-s', p}}^{\frac{4}{3}} |S_N u_N|_{\mathcal{W}^{\frac{1}{3}, 4}}^{\frac{4}{3}} \right) \\ &\leq C_T \sum_{k+l=2} \mathbb{E}(| : (S_N Z_R)^k (S_N Z_I)^l : |_{\mathcal{W}^{-s', p}}^3)^{\frac{4}{9}} \mathbb{E}(|u_N|_{\mathcal{W}^{\frac{5}{8}, 2}}^{\frac{12}{5}})^{\frac{5}{9}}, \end{aligned}$$

which is bounded again by Propositions 2.2 and 2.5. Note that in the last inequality we have used the Sobolev embedding $\mathcal{W}^{\frac{5}{8}, 2} \subset \mathcal{W}^{\frac{1}{3}, 4}$ and Hölder inequality. $\square$

Remark. We note that by Lemma 3.2, $(u_N)_{N \in \mathbb{N}}$ is bounded in

$$\begin{aligned} L^3(\Omega, L^3(0, T, \mathcal{W}^{\frac{2}{3}, 2})) &\cap L^{\frac{4}{3}}(\Omega, \mathcal{W}^{1, \frac{4}{3}}(0, T, \mathcal{W}^{-2, p})) \\ &\subset L^{\frac{4}{3}}(\Omega, L^3(0, T, \mathcal{W}^{\frac{2}{3}, 2})) \cap L^{\frac{4}{3}}(\Omega, \mathcal{W}^{\frac{1}{12}, 3}(0, T, \mathcal{W}^{-2, p})), \end{aligned}$$

and $L^3(0, T, \mathcal{W}^{\frac{2}{3}, 2}) \cap \mathcal{W}^{\frac{1}{12}, 3}(0, T, \mathcal{W}^{-2, p})$ is compactly embedded in $L^3(0, T, L_x^4)$. On the other hand, $(u_N)_{N \in \mathbb{N}}$ is also bounded in

$$\begin{aligned} L^2(\Omega, L^2(0, T, \mathcal{W}^{1, 2})) \cap L^{\frac{4}{3}}(\Omega, \mathcal{W}^{1, \frac{4}{3}}(0, T, \mathcal{W}^{-2, p})) \\ \subset L^{\frac{4}{3}}(\Omega, L^2(0, T, \mathcal{W}^{1, 2}) \cap \mathcal{W}^{\frac{1}{12}, 2}(0, T, \mathcal{W}^{-2, p})), \end{aligned}$$

and $L^2(0, T, \mathcal{W}^{1, 2}) \cap \mathcal{W}^{\frac{1}{12}, 2}(0, T, \mathcal{W}^{-2, p})$ is compactly embedded in $L^2(0, T, \mathcal{W}^{s, 2})$ for any s with $0 \leq s < 1$. In particular, since $\mathcal{W}^{\frac{5}{6}, 2} \subset \mathcal{W}^{\frac{1}{3}, 4}$, the embedding is compact in $L^2(0, T, \mathcal{W}^{\frac{1}{3}, 4})$. Finally, we note that $\mathcal{W}^{1, \frac{4}{3}}(0, T, \mathcal{W}^{-2, p})$ is compactly embedded in $C([0, T], \mathcal{W}^{-3, p})$.

Proof of Theorem 2.4. Let $\alpha > 0$ satisfy the condition in Lemma 3.2. We deduce from Lemma 3.2, Remark 3 and Markov inequality that the sequence $\{(u_N, Z)_{N \in \mathbb{N}}\}$ is tight in

$$(3.2) \quad L^3(0, T, L^4) \cap L^2(0, T, \mathcal{W}^{\frac{5}{6}, 2}) \cap C([0, T], \mathcal{W}^{-3, p}) \times C^\beta([0, T], \mathcal{W}^{-s', q} \cap \mathcal{W}^{-\delta, p})$$

for any $\beta < \alpha$ and $s' > s$. Fix such β and s' . By Prokhorov Theorem, there exists a subsequence, still denoted $\{(u_N, Z)_{N \in \mathbb{N}}\}$ which converges in law to a measure ν on the space (3.2). By Skorokhod Theorem, there exist $(\tilde{\Omega}, \tilde{\mathcal{F}}, \tilde{\mathbb{P}})$, $(\tilde{u}_N, \tilde{Z}_N)_{N \in \mathbb{N}}$ and $(\tilde{u}, \tilde{Z})$ taking values in the same space (3.2), satisfying $\mathcal{L}((u_N, Z)) = \mathcal{L}((\tilde{u}_N, \tilde{Z}_N))$ for any $N \in \mathbb{N}$, $\mathcal{L}((\tilde{u}, \tilde{Z})) = \nu$, and $\tilde{u}_N$ converges to $\tilde{u}$, $\tilde{\mathbb{P}}$ -a.s. in $L^3(0, T, L^4) \cap L^2(0, T, \mathcal{W}^{\frac{5}{6}, 2}) \cap C([0, T], \mathcal{W}^{-3, p})$, $\tilde{Z}_N$ converges to $\tilde{Z}$, $\tilde{\mathbb{P}}$ -a.s. in $C^\beta([0, T], \mathcal{W}^{-s', q} \cap \mathcal{W}^{-\delta, p})$. Moreover, by diagonal extraction, it can be assumed that this holds for any $T > 0$. It is easily seen that $(\tilde{u}, \tilde{Z})$ is a stationary process thanks to the convergence of $(\tilde{u}_N, \tilde{Z}_N)$ to $(\tilde{u}, \tilde{Z})$ in $C([0, T], \mathcal{W}^{-3, p}) \times C^\beta([0, T], \mathcal{W}^{-s', q})$. This convergence also implies $\mathcal{L}(Z) = \mathcal{L}(\tilde{Z})$. Note also that if we extract from the subsequence we took for the weak convergence $\tilde{\rho}_N$ to ρ , we have that for each $t \in \mathbb{R}$, $\mathcal{L}(\tilde{X}(t)) = \mathcal{L}(\tilde{u}(t) + \tilde{Z}(t)) = \rho$.

Write then,

$$\begin{aligned} \tilde{u}_N(t) - \tilde{u}_N(0) &= (\gamma_1 + i\gamma_2) \int_0^t H \tilde{u}_N(\sigma) d\sigma \\ (3.3) \quad &- (\gamma_1 + i\gamma_2) \int_0^t S_N(|S_N(\tilde{u}_N + \tilde{Z}_N)|^2 S_N(\tilde{u}_N + \tilde{Z}_N))(\sigma) d\sigma. \end{aligned}$$

It remains us to show that the right hand side of (3.3) converges, up to a subsequence, to

$$(\gamma_1 + i\gamma_2) \int_0^t H \tilde{u}(\sigma) d\sigma - (\gamma_1 + i\gamma_2) \int_0^t : |(\tilde{u} + \tilde{Z})|^2 (\tilde{u} + \tilde{Z}) : (\sigma) d\sigma,$$

$\tilde{\mathbb{P}}$ -a.s. in $C([0, T], \mathcal{W}^{-2, p})$. This can be checked as follows. First, the convergence of the linear term follows from the convergence of $\tilde{u}_N$ to $\tilde{u}$ in $L^2(0, T, \mathcal{W}^{\frac{5}{6}, 2}) \subset L^1(0, T, \mathcal{W}^{-2, p})$.

In order to prove the convergence of nonlinear terms, we again decompose the nonlinear terms into the four terms $F_0, \dots, F_3$ as in (2.7) and estimate them separately, i.e.,

$$\begin{aligned} & \sup_{t \in [0, T]} \left| \int_0^t : |\tilde{u} + \tilde{Z}|^2 (\tilde{u} + \tilde{Z}) : (\sigma) d\sigma - \int_0^t S_N (: |S_N(\tilde{u}_N + \tilde{Z}_N)|^2 S_N(\tilde{u}_N + \tilde{Z}_N) :)(\sigma) d\sigma \right|_{\mathcal{W}^{-2,p}} \\ & \leq \sum_{k=0}^3 \int_0^T |F_k(\tilde{u}, \tilde{Z}) - S_N F_k(S_N \tilde{u}_N, S_N \tilde{Z}_N)|_{\mathcal{W}^{-2,p}} d\sigma \\ & = I_0^N + I_1^N + I_2^N + I_3^N. \end{aligned}$$

We begin with the convergence of I_0^N .

$$\begin{aligned} I_0^N &= \int_0^T |S_N(|S_N \tilde{u}_N|^2 S_N \tilde{u}_N) - |\tilde{u}|^2 \tilde{u}|_{\mathcal{W}^{-2,p}} d\sigma \\ &\leq \int_0^T \| |S_N \tilde{u}_N|^2 S_N \tilde{u}_N - |\tilde{u}|^2 \tilde{u} \|_{L^{\frac{4}{3}}} d\sigma + \int_0^T |(S_N - I)|\tilde{u}|^2 \tilde{u}|_{L^{\frac{4}{3}}} d\sigma. \end{aligned}$$

Since $\tilde{u}_N$ converges to $\tilde{u}$ a.s. in $L^3(0, T, L_x^4)$, by dominated convergence, the same holds for $S_N \tilde{u}_N$, thus $|S_N \tilde{u}_N|^2 S_N \tilde{u}_N$ converges to $|\tilde{u}|^2 \tilde{u}$ a.s. in $L^{\frac{4}{3}}(0, T, L_x^{\frac{4}{3}})$. Similarly, the second term converges to zero by dominated convergence, therefore, I_0^N converges to 0. Next, we use Lemma 3.1 to obtain

$$\begin{aligned} I_1^N &\leq \int_0^T |F_1(\tilde{u}, \tilde{Z}) - S_N F_1(S_N \tilde{u}_N, S_N \tilde{Z}_N)|_{\mathcal{W}^{-(\delta+\frac{1}{3}),p}} d\sigma \\ &\lesssim |\tilde{Z} - S_N \tilde{Z}_N|_{C([0,T], \mathcal{W}^{-\delta,p})} \left(|S_N \tilde{u}_N|_{L^2(0,T, \mathcal{W}^{\frac{1}{3},4})}^2 + |\tilde{u}|_{L^2(0,T, \mathcal{W}^{\frac{1}{3},4})}^2 \right) \\ &\quad + \left(|S_N \tilde{Z}_N|_{C([0,T], \mathcal{W}^{-\delta,p})} + |\tilde{Z}|_{C([0,T], \mathcal{W}^{-\delta,p})} \right) \left(|S_N \tilde{u}_N|_{L^2(0,T, \mathcal{W}^{\frac{1}{3},4})} + |\tilde{u}|_{L^2(0,T, \mathcal{W}^{\frac{1}{3},4})} \right) \\ &\quad \times |S_N \tilde{u}_N - \tilde{u}|_{L^2(0,T, \mathcal{W}^{\frac{1}{3},4})}. \end{aligned}$$

Hence, I_1^N converges to 0 since $S_N \tilde{Z}_N$ converges to $\tilde{Z}$ in $C([0, T], \mathcal{W}^{-\delta,p})$ and, by Remark 3, $S_N \tilde{u}_N$ converges to $\tilde{u}$ in $L^2(0, T, \mathcal{W}^{\frac{1}{3},4})$. Concerning F_2 , we proceed as for F_1 , and we use Lemma 3.1 in the same way;

$$\begin{aligned} I_2^N &\lesssim \sum_{k+l=2} (| : \tilde{Z}_R^k \tilde{Z}_I^l : |_{L^2(0,T, \mathcal{W}^{-\delta,p})} + |S_N (: (S_N \tilde{Z}_{N,R})^k (S_N \tilde{Z}_{N,I})^l :)|_{L^2(0,T, \mathcal{W}^{-\delta,p})}) \\ &\quad \times |S_N \tilde{u}_N - \tilde{u}|_{L^2(0,T, \mathcal{W}^{\frac{1}{3},4})} \\ &\quad + \sum_{k+l=2} | : \tilde{Z}_R^k \tilde{Z}_I^l : - S_N (: (S_N \tilde{Z}_{N,R})^k (S_N \tilde{Z}_{N,I})^l :)|_{L^2(0,T, \mathcal{W}^{-\delta,p})} \\ &\quad \times \left(|S_N \tilde{u}_N|_{L^2(0,T, \mathcal{W}^{\frac{1}{3},4})} + |\tilde{u}|_{L^2(0,T, \mathcal{W}^{\frac{1}{3},4})} \right). \end{aligned}$$

Note that Proposition 2.2 implies the convergence to zero of the second term in $L^p(\Omega)$, thus extracting a subsequence, $\tilde{\mathbb{P}}$ -a.s. The first term goes to 0, too since, again, $S_N \tilde{u}_N$

converges to $\tilde{u}$ in $L^2(0, T, \mathcal{W}^{\frac{1}{3}, 4})$. The term I_3^N can be treated similarly. We deduce from this convergence result that $\tilde{u}$ satisfies

$$\tilde{u}(t) - \tilde{u}(0) = (\gamma_1 + i\gamma_2) \int_0^t H\tilde{u}(\sigma) d\sigma - (\gamma_1 + i\gamma_2) \int_0^t : |(\tilde{u} + \tilde{Z})|^2 (\tilde{u} + \tilde{Z}) : (\sigma) d\sigma.$$

Since it is clear that $\tilde{Z}$ satisfies the second equation in (3.1), we easily deduce that $\tilde{X} = \tilde{u} + \tilde{Z}$ is a stationary solution of (2.11) on $(\tilde{\Omega}, \tilde{F}, \tilde{\mathbb{P}})$. Moreover, it is not difficult to prove that $\tilde{u}$ is continuous with values in $\mathcal{W}^{-s, q}$, which ends the proof of Theorem 2.4. $\square$

References

- [1] P.B. Blakie, A.S. Bradley, M.J. Davis, R.J. Ballagh and C.W. Gardiner, “Dynamics and statistical mechanics of ultra-cold Bose gases using c-field techniques,” *Advances in Physics*. **57** (2008) no.5, 363-455.
- [2] A. de Bouard, A. Debussche and R. Fukuizumi, “Long time behavior of Gross Pitaevskii equation at positive temperature,” *SIAM. J. Math. Anal.* **50** (2018) 5887-5920.
- [3] A. de Bouard, A. Debussche and R. Fukuizumi, “Two dimensional Gross-Pitaevskii equation with space-time white noise,” *Int. Math. Res. Not.* (2022) rnac137
- [4] G. Da Prato and A. Debussche, “Strong solutions to the stochastic quantization equations,” *The Annals of Probability*. **31** (2003) no.4, 1900-1916.
- [5] G. Da Prato and A. Debussche, “Two-dimensional Navier-Stokes equations driven by a space-time white noise,” *J. Funct. Anal.* **196** (2002) 180-210..
- [6] R. Duine and H. Stoof, “Stochastic dynamics of a trapped Bose-Einstein condensate,” *Phys. Rev. A* **65** (2001) p. 013603.
- [7] C.W. Gardiner and M.J. Davis, “The stochastic Gross-Pitaevskii equation: II,” *J. Phys. B* **36** (2003) 4731-4753.
- [8] M. Hoshino, “Global well-posedness of complex Ginzburg-Landau equation with a space-time white noise,” *Ann. l’institut Henri Poincaré (B) Probability and Statistics*. **54** (2018) 1969-2001.
- [9] A. Jensen and S. Nakamura, “ L^p -mapping properties of functions of Schrödinger operators and their applications to scattering theory,” *J. Math. Soc. Japan*. **47** (1995) 253-273.
- [10] T. Matsuda, “Global well-posedness of the two-dimensional stochastic complex Ginzburg-Landau equation with cubic nonlinearity,” *arXiv: 2003.01569*.
- [11] M. E. Taylor, “Tools for PDEs,” *Pseudodifferential Operators, Paradifferential Operators, and Layer Potentials*, Math. Surveys Monogr., vol. 81, American Mathematical Society, Providence, RI (2000)
- [12] W. J. Trenberth “Global well-posedness for the two-dimensional stochastic complex Ginzburg-Landau equation” *arXiv:1911.09246v1*.
- [13] P. Tsatsoulis and H. Weber, “Spectral gap for the stochastic quantization equation on the 2-dimensional torus,” *Ann. l’IHP. Probabilités et statistiques* **54** (2018) 1204-1249.
- [14] C.N. Weiler et al. “Spontaneous vortices in the formation of Bose-Einstein condensates,” *Nature* **455** (2008) nature 07334.