

Influence of Magnetic Field and Magnetic Anisotropy on the Quench Characteristics of Bi-2223/Ag Multifilamentary Tapes

Taketsune Nakamura, Shingo Tsuchiya, Akihisa Fujio, Tsutomu Hoshino and Itsuya Muta

Abstract— We examined the normal propagation characteristics of Bi-2223/Ag multifilamentary tapes as a function of magnetic field direction. Anisotropic current transport characteristics were estimated by taking into account the effect of the misorientation of the grains. Anisotropic normal zone propagation properties were examined based on the numerical simulation, and showed that the thermal properties were influenced by the anisotropic current transport characteristics.

Keywords— Bi-2223, Anisotropy, Normal Zone Propagation Properties

I. INTRODUCTION

TREMENDOUS efforts have been made for the development of the power application systems using high- T_c superconducting (HTS) tapes. Silver sheathed $\text{Bi}_2\text{Sr}_2\text{Ca}_2\text{Cu}_3\text{O}_{10+x}$ (Bi-2223/Ag) tape is the most promising candidate for such applications. For the design of the HTS systems, the estimation of the stability is crucial. However, HTS is thermally stable at elevated temperatures. Furthermore, current transport curves of HTS are complicated depending on temperature and/or transport current, and are also anisotropic with respect to the applied field direction. This means that the Joule heating process is complicated. Therefore, taking account of such effects is important to the quantitative understanding of the thermal properties. Shimizu et al.[1] showed that the power-law expression for the $E-J$ curves in zero magnetic field can be used for the calculation of the normal zone propagation properties. On the other hand, Kiss et al.[2,3] proposed the analytical expression for the properties in high magnetic fields, and Okamoto et al.[4] showed that the $E-J$ curves of Bi-2223/Ag multifilamentary tapes can be described quantitatively by that expression. Therefore, appropriate expression for the $E-J$ characteristics has to be selected depending upon the application, i.e., zero-field or non-zero field application. On the other hand, we showed that the angular dependence of the current transport characteristics in Bi-2223/Ag multifilamentary tape is well described by the model which takes into account the misorientation of the grains[5].

In this paper, we present the normal propagation properties of the Bi-2223/Ag multifilamentary tape as a function of temperature and magnetic field direction studied by computer simulation.

T. Nakamura, S. Tsuchiya, A. Fujio, T. Hoshino and I. Muta are with the Graduate School of Engineering, Kyoto University, Kyoto 606-8501, Japan. E-mail: tk.naka@kuee.kyoto-u.ac.jp

II. EXPRESSION FOR $E-J$ CURVES

A. $E-J$ curves in zero external field

$E-J$ characteristics in zero external field are expressed by the power-law approximation with index, n , as follows:

$$E = E_{c0} \left(\frac{J}{J_{c0}} \right)^n \quad (1)$$

where E_{c0} is electric field criterion to define J_{c0} . As shown in Fig.2, this expression is valid in a wide range of temperatures.

B. $E-J$ curves in high external field

$E-J$ properties in high magnetic field are very complicated, depending upon the temperature and/or current. We use the expression proposed by Kiss et al.[2,3]. In this expression, so-called glass-liquid transition temperature, T_g , is introduced, and $E-J$ properties are expressed as follows:

for $T \leq T_g$,

$$E(J, T) \cong \frac{\rho_{ff}}{m+1} J \left[\frac{J}{J_0(T)} \right]^m \left[1 - \frac{J_{cm}(T)}{J} \right]^{m+1} \quad (2)$$

for $T > T_g$,

$$E(J, T) \cong \frac{\rho_{ff}}{m+1} |J_{cm}(T)| \left[\frac{|J_{cm}(T)|}{J_0(T)} \right]^m \cdot \left[\left(1 + \frac{J}{|J_{cm}(T)|} \right)^{m+1} - 1 \right] \quad (3)$$

where ρ_{ff} and m are, respectively, the uniform flux flow resistivity and the parameter describing the shape of the local critical current distribution. J_{cm} is a minimum local critical current density, and expressed as follows:

$$J_{cm} = J_T \left| 1 - \frac{T}{T_g} \right|^{2\nu} \quad (4)$$

where J_T and ν are constant. $J_0(T)$ is roughly proportional to T^{-1} .

C. Anisotropy of $E-J$ curves

We have taken account of the misorientation of the grains in Bi-2223/Ag tape, and shown that the glass-liquid transition flux density, B_g , as a function of the angle between the external field direction and the axis perpendicular to

the tape surface, φ , is well described by the following expression[5]:

$$B_g(\varphi, T) = B_{g,\perp}(T) (\cos^2\varphi + \gamma^{-2}\sin^2\varphi)^{-\frac{1}{2}} \quad (5)$$

where γ is the effective anisotropy factor and defined as follows:

$$\gamma = \frac{B_{g,\parallel}}{B_{g,\perp}} \quad (6)$$

$B_{g,\parallel}$ and $B_{g,\perp}$ are B_g for, respectively, parallel and perpendicular field direction with respect to the tape face. We assume that (5) and (6) are valid for arbitrary temperature considered in this paper.

III. SIMULATION DETAIL

Fig.1 shows the schematics of the simulation model. The total length of 100mm is considered, and the 50mm of it is simulated as both sides are symmetrical. For simplicity, Bi-2223/Ag multifilamentary tape is treated as two layers of superconducting Bi layer and silver sheath. The following heat balance equation is solved by finite difference method:

$$c(T) \frac{\partial \theta}{\partial t} = k \left(\frac{\partial^2 \theta}{\partial x^2} + \frac{\partial^2 \theta}{\partial y^2} \right) + G + Q_d \quad (7)$$

where c and k are, respectively, the heat capacity and the thermal conductivity. θ is temperature difference and defined as $T - T_0$ with bias temperature (T_0). $G(= EJ)$ is Joule heating and calculated with one of (1), (2) and (3). Q_d is disturbance which is applied by the heater, considered on the silver sheath as shown in Fig.1. We take into account the temperature dependence of c . Parameter k is assumed to be constant. Conduction cooling is considered, and boundary condition for the surface of the tape is adiabatic. The temperature of the end of the tape is fixed at T_0 . Furthermore, we assume that the equipotential surface through any points is at right angles to the longitudinal direction of the tape at that points. Then, current density flowing in the superconducting Bi layer, J_{sc} , is determined with the ratio of silver, η , as follows:

$$\frac{\eta}{\eta + 1} \frac{E(J_{sc})}{\rho_{ag}} = J_{total} - \frac{1}{\eta + 1} J_{sc} \quad (8)$$

where J_{total} and ρ_{ag} are the total transport current and the resistivity of the silver sheath, respectively. Equation (8) is calculated numerically using Newton-Raphson method.

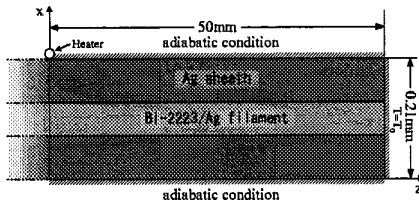


Fig. 1. Simulation Model.

IV. RESULTS AND DISCUSSION

A. Normal Propagation Properties in Zero Magnetic Field

Firstly, we simulate the normal propagation properties in zero magnetic field by using the measured $V - I$ curves. Bi-2223/Ag tape, which parameters are listed in TABLE I, was cooled by a GM cryocooler. $V - I$ measurements were performed by means of the DC four points method.

Fig.2 shows the temperature dependence of $V - I_{sc}$ characteristics in zero magnetic field. Effect of the sharing current flowing in the silver sheath was eliminated with the measured voltage, V , by $I_{sc} = I_{total} - V/R_{ag}$. R_{ag} is the resistance of the silver sheath. As can be seen in Fig.2, we can roughly describe the $V - I_{sc}$ curves with (1) in a wide range of temperature. Fig.3 shows the temperature dependence of the critical currents and the n -values which are defined by the voltage criterion of $0.6\mu V$, corresponding to the electric field criterion of $1\mu Vcm^{-1}$. We can describe I_c with the following expression:

$$I_c(T) = I_{c0} \left(1 - \frac{T}{T_c} \right)^\alpha \quad (9)$$

Critical temperature, $T_c(=107.6K)$, was obtained from the Resistance vs. Temperature relation measured with small bias current ($=20mA$). α ($=1.39$) is a fitting parameter. We can see that the n -value is temperature dependent weakly at $T \leq 95K$, and decreases drastically above the temperature. Then this dependency is also taken into account for the simulation.

TABLE I
THE PARAMETERS OF THE SAMPLE USED FOR THE $V - I$
MEASUREMENTS IN ZERO MAGNETIC FIELD

Cross sectional area	$3.2 \text{ mm}^w \times 0.21 \text{ mm}^t$
Number of filaments	55
Ratio of silver	2.8
T_c	107.6 K

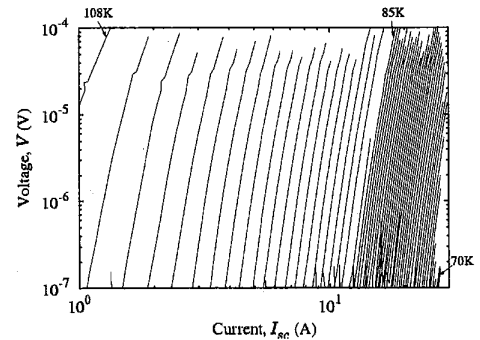


Fig. 2. Temperature dependence of $V - I_{sc}$ curves in zero magnetic field. The sharing current flowing in the silver sheath (V/R_{ag}) is eliminated by $I_{sc} = I_{total} - V/R_{ag}$. Temperature intervals are 0.4K from 70.0K to 85.0K, and 1.0K from 85.0K to 108.0K, respectively.

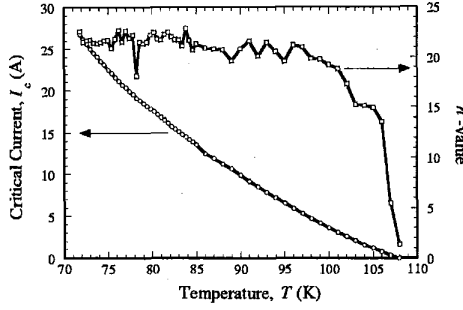


Fig. 3. Temperature dependence of I_c and n -value in Bi-2223/Ag multifilamentary tape in zero external field. I_c and n -value are defined with electric field criterion of $1\mu\text{Vcm}^{-1}$.

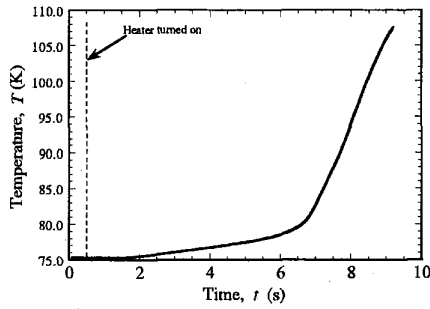


Fig. 4. Temperature trace with a transport current of 19.0A, $I_{\text{total}}/I_c=0.9$. Initial temperature is 75.3K. Distance from the heater is 6.0mm. Heater (heat flux of $1.6 \times 10^5 \text{Wm}^{-2}$) is turned on at $t=1.0\text{s}$.

Combining (1), (7), (8) and (9), one can simulate the normal zone propagation properties. Fig.4 shows temperature traces for different heater powers. Bias current, 19.0A, is applied to the sample. After the temperature is stabilized, the heater (heat flux of $1.6 \times 10^5 \text{Wm}^{-2}$) is turned on from $t=1.0$ to 15.0s. By turning on the heater, temperature begins to increase gradually due to the Joule heating in the Bi filaments. Then, current sharing is increased due to the rising temperature, and temperature increases rapidly at $t=15\text{s}$. This behavior is qualitatively in agreement with the experimental data in [6].

B. Normal Propagation Properties in high magnetic field

Next, we simulate the normal propagation properties in high magnetic field with the same simulation code in subsection IV-A. In high magnetic field, $V - I_{\text{sc}}$ characteristics vary complicatedly depending upon the temperature and/or the transport current, and (1) is no more applicable. Then, we adopt (2) and (3) for the calculation of the Joule heating. The experimental data in high magnetic field is not obtained yet, hence we use the data measured in a field of 1T by Okamoto et al.[4]. Fig.5 shows the $E - J$ curves reproduced by (2) and (3), and the parameters obtained in [4] are listed in TABLE II. Furthermore, in order to study the influence of the anisotropic current transport

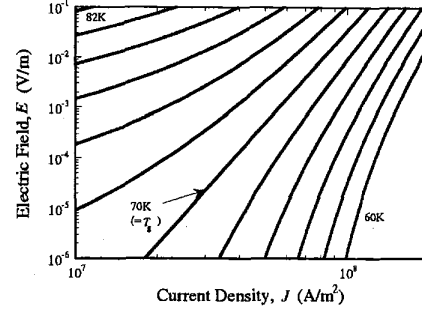


Fig. 5. Temperature dependence of $E - J$ characteristics reproduced by (2) and (3) with the parameters listed in TABLE II. External magnetic field (1T) is applied perpendicularly to the sample face [4].

TABLE II
THE PARAMETERS OF THE EXPERIMENTAL DATA USED FOR THE SIMULATION IN HIGH MAGNETIC FIELD[4]

m	5.1
ν	0.5
ρ_{ff}	$1 \times 10^{-7} \Omega\text{m}$
J_T	$5.15 \times 10^8 \text{ A/m}^2$
T_g at 1T	70.0K
$B_{g,\parallel}/B_{g,\perp}$	5.0
T_c	101 K

characteristics, (5) and (6) are assumed to be valid for (2), (3) and (4). Temperature dependence of glass-liquid transition flux density for perpendicular field direction, $B_{g,\perp}$, is expressed from [4] as follows:

$$B_{g,\perp} = B_{g,\perp}(0) \left(1 - \frac{T}{T_c}\right)^\beta \quad (10)$$

where $B_{g,\perp}(0)=5.9\text{T}$, $T_c=101\text{K}$, and $\beta=1.5$, respectively. Combining (5) and (10), and replacing parameters from B_g to B and T to T_g , one can estimate the angular dependent T_g as follows:

$$T_g(\varphi, B) = \left[1 - \left(\frac{B}{B_{g,\perp}(0)} \right)^{\frac{1}{\beta}} (\cos^2 \varphi + \gamma^{-2} \sin^2 \varphi)^{\frac{1}{2\beta}} \right] T_c \quad (11)$$

Fig.6 shows T_g vs. φ curves calculated by (11). $E - J_{\text{sc}}$ curves at temperature below and above $T_g(\varphi)$ are determined by (2) and (3), respectively. Angular dependence of the normal propagation properties are considered by varying T_g with (11). Fig.7 and 8 show the time evolutions of the voltages and temperatures for different field directions. Thermal disturbance (heat flux of $1.8 \times 10^5 \text{Wm}^{-2}$) is applied by the heater from $t=1.0$ to 15s with the bias current of 15.0A. Initial temperature is 70.0K which corresponds to $T_{g,\perp} (= T_g(\varphi=0^\circ))$ in this field (1.0T). That is, the properties except at $\varphi=0^\circ$ are at temperatures below $T_g(\varphi)$ before turning on the heater, and $E - J_{\text{sc}}$ properties are

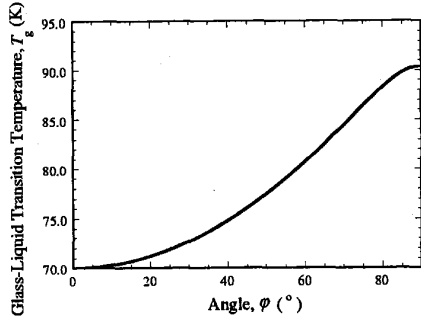


Fig. 6. T_g vs. ϕ characteristics at $B=1T$ obtained from (11). $B_{g,\perp}(0)=5.9T$, $T_c=101K$, and $\beta=1.5$, respectively.

determined by (2). As can be seen, voltage traces vary drastically depending upon the field direction after turning on the heater. Namely, we can't see any dissipations for the properties at $\phi = 30, 45$ and 60° within the voltage window of Fig.7. Temperature increase is also small correspondingly. However, voltage and temperature increase rapidly for $\phi = 0, 10$ and 20° . Fig.9 shows the current sharing rate and the normalized temperature (T/T_g) at $t=16.0s$, plotted as a function of angle. As can be seen, current sharing rate increases steeply when T is higher than T_g . Because

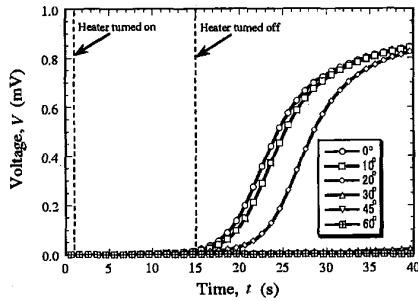


Fig. 7. Voltage traces for different angles with a transport current of $15.0A$. Heater(heat flux of $1.8 \times 10^5 Wm^{-2}$) is turned on from $t=1.0$ to $15s$. Initial temperature is $70.0K$.

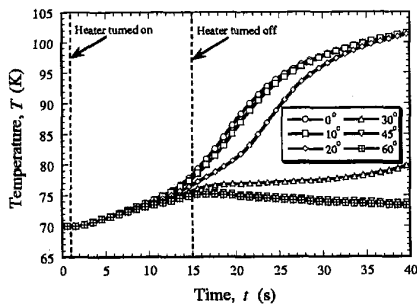


Fig. 8. Temperature traces for different angles. Heater(heat flux of $1.8 \times 10^5 Wm^{-2}$) is turned on from $t=1.0$ to $15s$. Distance from the heater is $6.0mm$. Initial temperature is $70.0K$.

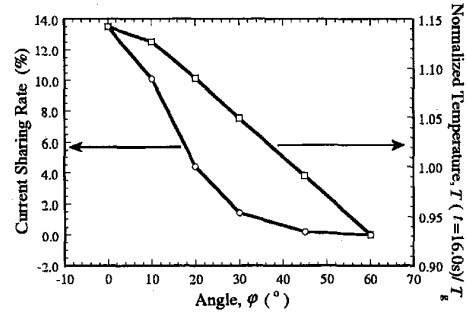


Fig. 9. Current sharing rate and normalized temperature (T/T_g) vs. angle (ϕ) characteristics at $t=16.0s$.

$E - J_{sc}$ curves are determined by (3) at $T > T_g$ and the dissipation is increased. That is to say, normal zone propagation properties depend on the field direction and stability have to be discussed by taking into account the field direction. Experimental study has been planned.

V. CONCLUSION

Normal propagation properties were studied in zero and high magnetic field based on the computer simulation. In zero magnetic field, the power-law expression for the $E - J$ curves was used to calculate the Joule heating. In high magnetic field, however, $E - J$ curves vary complicatedly, and the analytical expressions proposed by Kiss et al. was adopted. Anisotropy of the current transport characteristics were also considered. It was shown that the normal propagation property was influenced by the complicated $E - J$ characteristics and the anisotropy.

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