



A New Cation-Ordered Structure Type with Multiple Thermal Redistributions in $\text{Co}_2\text{InSbO}_6$

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Abstract: Cation ordering in solids is important for controlling physical properties and leads to ilmenite (FeTiO_3) and LiNbO_3 type derivatives of the corundum structure, with ferroelectricity resulting from breaking of inversion symmetry in the latter. However, a hypothetical third ABO_3 derivative with $R32$ symmetry has never been observed. Here we show that $\text{Co}_2\text{InSbO}_6$ recovered from high pressure has a new, ordered- $R32$ A_2BCO_6 variant of the corundum structure. $\text{Co}_2\text{InSbO}_6$ is also remarkable for showing two cation redistributions, to $(\text{Co}_{0.5}\text{In}_{0.5})_2\text{CoSbO}_6$ and then $\text{Co}_2\text{InSbO}_6$ variants of the ordered- LiNbO_3 A_2BCO_6 structure on heating. The cation distributions change magnetic properties as the final ordered- LiNbO_3 product has a sharp ferrimagnetic transition unlike the initial ordered- $R32$ phase. Future syntheses of metastable corundum derivatives at pressure are likely to reveal other cation-redistribution pathways, and may enable ABO_3 materials with the $R32$ structure to be discovered.

Cation ordering within extended oxide structures is an important way to control physical properties such as introduction of ferroelectricity and multiferroism from arrangements that break inversion symmetry.^[1,2] This is notably illustrated by the corundum ($\alpha\text{-Al}_2\text{O}_3$) type A_2O_3 structure which has a simple centrosymmetric arrangement with rhombohedral space group symmetry. A_2O_9 dimer units of two octahedra sharing a common face are separated by single vacant octahedra to form AA_AA chains (Figure 1).

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Two cation-ordered ABO_3 derivatives are known—the ilmenite (FeTiO_3) and LiNbO_3 types with centric and acentric $R3c$ symmetry respectively. Both have AB cation pairs in the dimer units with antiparallel AB_BA alignment in ilmenite but parallel AB_AB order leading to polarity in the LiNbO_3 type. It is intriguing to note that a third ABO_3 cation ordering type is also possible within the corundum unit cell as shown in Figure 1. This structure has $R32$ symmetry with AA_BB chains of dimer pairs, and no examples have been reported. Further cation ordering within the ABO_3 structures leads to A_2BCO_6 derivatives, referred to in the literature as “ordered- ABO_3 ” types. These were first found in Li_2MTeO_6 phases which adopt the

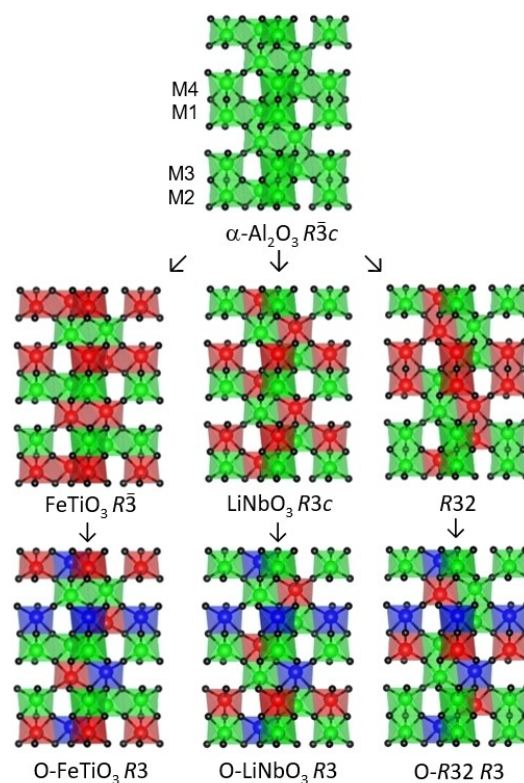


Figure 1. Crystal structures of the corundum type A_2O_3 structure (top), and ABO_3 (middle row) and A_2BCO_6 (“ordered- ABO_3 ”, bottom row) derivatives obtained through cation ordering. Symmetry descents are indicated by the arrows and space groups are shown. Colours indicate the cation occupancies of octahedra in each structure (A/B/C = green/red/blue). The four site labels shown by the A_2O_3 structure are used throughout this paper, where M4 = Sb in $\text{Co}_2\text{InSbO}_6$.

ordered-LiNbO₃ structure for M=Zr and Hf,^[3] and the ordered-ilmenite type for M=Ge.^[4] The “ordered-*R32*” A₂BCO₆ derivative of the *R32*-type has not been reported. These three A₂BCO₆ arrangements all have *R3* symmetry with four symmetry-independent octahedral cation sites. This structure is known as the Ni₃TeO₆ (NTO) type and represents a special case of all three A₂BCO₆ types (Figure 1) where A=B=Ni and C=Te.

The non-centrosymmetric cation ordered corundum derivatives offer a rich variety of properties^[5–7] according to their point group symmetry.^[8] *R3c* LiNbO₃-types (*C*_{3v} point group) have allowed polar and piezoelectric activity, and LiNbO₃ itself is an important ferroelectric, piezoelectric and non-linear optical material.^[9] The *R32* ABO₃ structure (*D*₃) is notable as belonging to the class of space groups that are non-polar but allow enantiomorphic and piezoelectric activity, so discovered examples would be of great interest. The ordered A₂BCO₆ (NTO-type) structures (*C*₃) have allowed polar, chiral and piezoelectric properties arising from their *R3* symmetry. Further coupling of these structural orders to magnetism (multiferroism) can be introduced by use of magnetic cations that adopt long-range spin orders at low temperatures.^[10] High pressure is often used to stabilise these acentric cation-ordered corundum derivatives, for example, MnTiO₃ changes from a centric ilmenite type at ambient pressure to an acentric LiNbO₃-type high-pressure polymorph where weak ferromagnetism offers a mechanism for multiferroic switching.^[11] Mn(Fe_{0.5}M_{0.5})O₃ (M=Nb, Ta) are further examples of ABO₃ LiNbO₃-types, with Fe/M disorder.^[12] Within the *R3* A₂BCO₆ types, β-Mn₂InSbO₆ has the ordered-ilmenite arrangement^[13] while M₂ScSbO₆ (M=Mn,^[14] Co,^[15] Ni^[16]) and Mn₂FeWO₆^[17] are ordered-LiNbO₃ types. Mn₂FeMoO₆ recovered from high pressure synthesis has an ordered-LiNbO₃ structure but this changes to an ordered-ilmenite type on heating and the stabilisation of these two types was rationalised from band structure calculations.^[18] Magnetoelectric effects are reported in Ni₃TeO₆^[19–21] and ternary NTO-type analogues have recently been discovered for A₃TeO₆ (A=Mn, Co)^[22] and Mn₃WO₆^[23] at high pressure. In this communication, we report the synthesis of the new double-corundum material Co₂InSbO₆ and thermal cation redistributions of unprecedented complexity from the previously unobserved ordered-*R32* type to two different ordered-LiNbO₃ types.

A mixture of CoO, In₂O₃ and Sb₂O₃ in stoichiometric proportions for product Co₂InSbO₆ was treated under high pressure and temperature conditions using a multi-anvil apparatus. Further details are in Supporting Information. A sample recovered from 6 GPa and 1373 K was found to contain a CaCl₂-type product with an orthorhombic structure that is unrelated to the corundum types and characterisation of this phase is described in Supporting Information. Synthesis under 8 GPa and 1373 K led to a recovered Co₂InSbO₆ product with *R3* symmetry (lattice parameters *a*=5.2882(3) Å and *c*=14.029(1) Å) consistent with A₂BCO₆ structures shown in Figure 1.

Synchrotron powder X-ray diffraction data from the Co₂InSbO₆ sample were collected in situ while heating from 300 to 1073 K to determine the structure and any thermal

changes. Refinement of the recovered Co₂InSbO₆ product structure at 300 K (fit and results in Supporting Information) gave cation site occupancies M1=Co_{0.3}In_{0.7}, M2=Co_{0.7}In_{0.3}, M3=Co and M4=Sb. Remarkably, the two Co-rich sites M2 and M3 are present in the same dimer units so that the cation distribution is close to the ordered-*R32* type rather than the ordered-ilmenite or LiNbO₃ types. This is an important structural discovery given that no ordered-*R32* A₂BCO₆ or *R32* ABO₃ structures have been reported amongst many known corundum-derived phases. The present compound thus represents a new structural type within the corundum group (Figure 1).

Comparison of the variable temperature patterns in Figure 2 shows that Co₂InSbO₆ persists as a *R3* corundum-derived material up to 1073 K but changes in peak positions and intensities near 900 K reveal structural rearrangement. Initial unconstrained fits (summarised in Supporting Information) demonstrated that while one cation site (M4) remains occupied by Sb throughout, Co/In occupancies at

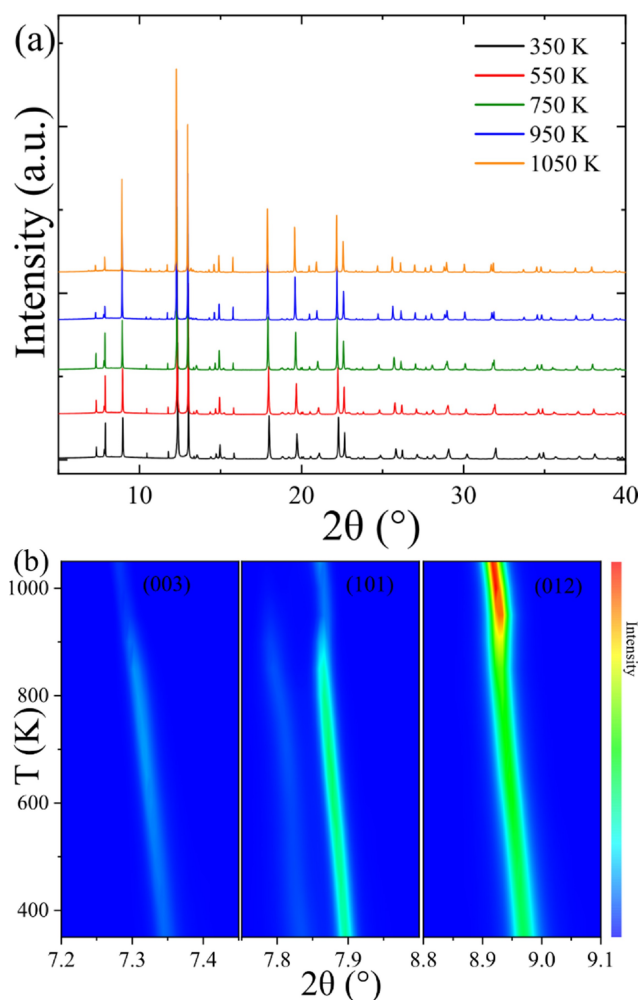
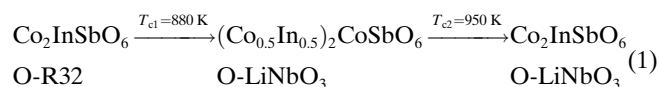


Figure 2. a) Selected powder X-ray diffraction data from the high pressure Co₂InSbO₆ sample collected in situ while heating from 300 to 1073 K. b) Diffraction intensity map for low-angle (003), (101) and (012) peaks. Changes between 850 and 950 K reflect the evolving cation distributions.

the other three sites change with temperature. Thermal variations of the cell parameters, M1–M3 site occupancies and M–O bond lengths from final refinements constrained to the overall $\text{Co}_2\text{InSbO}_6$ stoichiometry are shown in Figure 3. Further results are tabulated in Supporting Information.

Refined cell parameters for $\text{Co}_2\text{InSbO}_6$ in Figure 3a show a change in curvature on heating above 650 K and a dramatic lattice anomaly between 850 and 950 K. A negative expansion in the a parameter, a small excess positive expansion in c , and almost zero volume expansion over this interval are observed. Corresponding changes in the Co/In occupancies at the M1–M3 sites in Figure 3b reveal that two successive cation rearrangements occur on heating. Cation populations remain constant from 300 up to 650 K, reflecting the kinetic sluggishness of migration on the timescale of the X-ray diffraction experiment. Above 650 K the M3 site remains fully occupied by Co, but the Co/In occupancies of the M1 and M2 sites gradually converge and are estimated to become equal at $T_{c1} = 880$ K from a mean field fit to the occupancy difference (shown in Supporting Information). This transition corresponds to a change between different A_2BCO_6 types in $\text{Co}_2\text{InSbO}_6$: from

the ordered $R32$ -type in the recovered sample to an ordered- LiNbO_3 arrangement upon heating. In the latter structure where the two A sites (M1 and M2) have identical disordered $\text{Co}_{0.5}\text{In}_{0.5}$ compositions at T_{c1} . This transition highlights the instability of the $R32$ cation distribution and ordered- $R32$ derivative at ambient conditions. Furthermore, the ordered- LiNbO_3 phase ($\text{Co}_{0.5}\text{In}_{0.5}$) $_2\text{CoSbO}_6$ shows a thermal instability immediately above T_{c1} as In from the M1 and M2 sites rapidly exchanges with Co from the M3 site between T_{c1} and $T_{c2} \approx 950$ K. Above T_{c2} , the cation distribution is close to another A_2BCO_6 ordered- LiNbO_3 type, with A sites having M1 $\approx 80\%$ and M2 = 100% Co. Hence the discovered sequence of structural changes (showing ideal A_2BCO_6 cation site occupancies as displayed in the Table of Contents graphic) is;



These observations demonstrate that the overall transformation of $\text{Co}_2\text{InSbO}_6$ from an ordered- $R32$ to an ordered- LiNbO_3 polymorph occurs via a cation-disordered ($\text{Co}_{0.5}\text{In}_{0.5}$) $_2\text{CoSbO}_6$ ordered- LiNbO_3 intermediate phase.

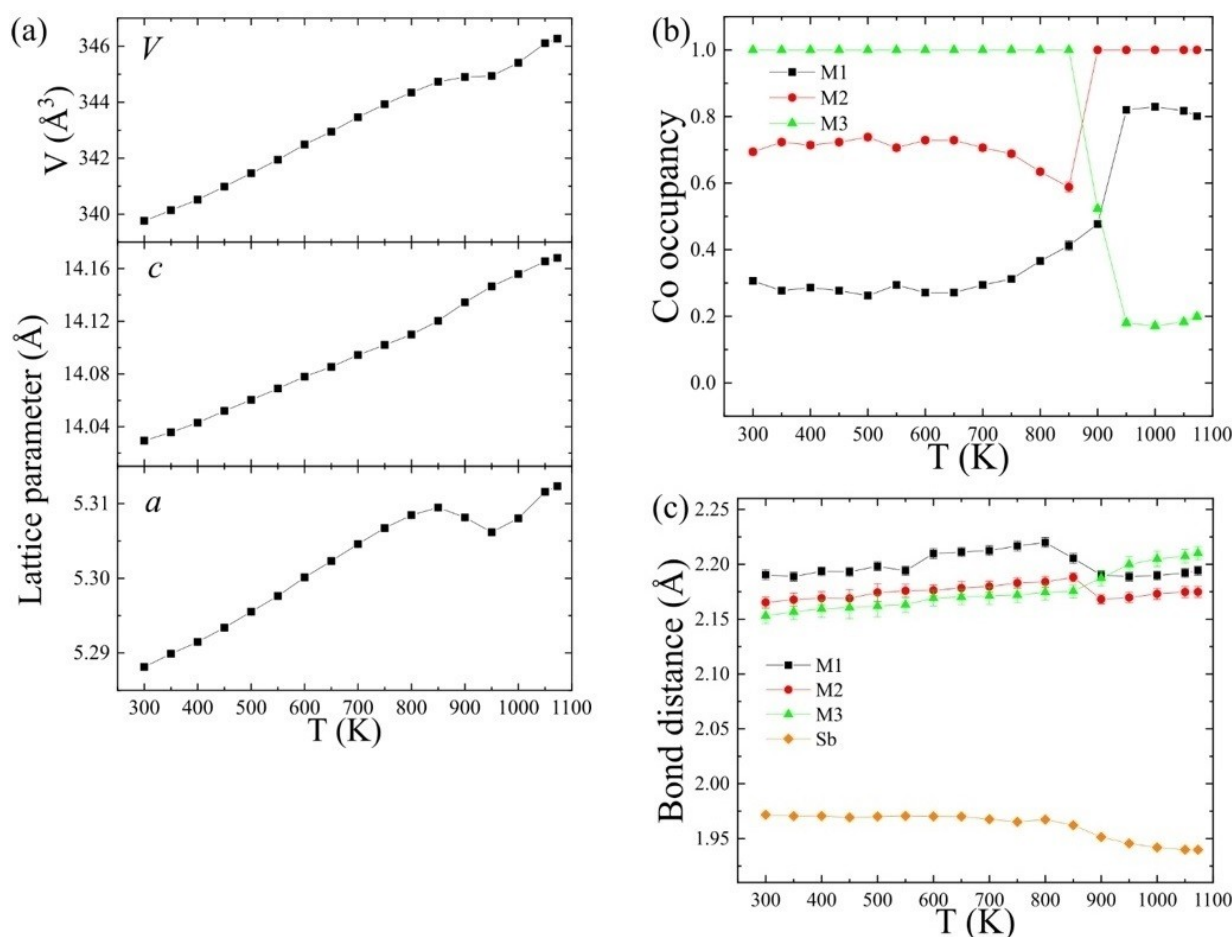


Figure 3. Refined X-ray structure parameters from $\text{Co}_2\text{InSbO}_6$ while heating from 300 to 1073 K. a) Lattice parameters and cell volume showing the structural anomaly between 850 and 950 K. b) Co occupancies at M1, M2 and M3 sites revealing Co/In intersite rearrangements. c) Average M–O bond lengths for each MO_6 octahedron where M1–M3 are occupied by Co/In and Sb is at the M4 site.

The second transition, from a cation-disordered to a cation-ordered structure on heating, is unusual given the loss of configurational entropy and evidences likely metastability of the intermediate $(\text{Co}_{0.5}\text{In}_{0.5})_2\text{CoSbO}_6$ phase.

Average metal–oxygen bond lengths (Figure 3c) are consistent with the Co/In occupancy rearrangements, given the cation sizes (6-coordinate ionic radii are $\text{Co}^{2+}=0.745$, $\text{In}^{3+}=0.80$, and $\text{Sb}^{5+}=0.60$ Å).^[24] M1–O and M2–O distances both decrease on heating from 850 to 950 K as their Co-populations increase, while an increase in the M3–O distance reflects the almost complete replacement of Co^{2+} by In^{3+} . The similar sizes and charges of Co^{2+} and In^{3+} cations allow the changing cation distributions at M1–M3 sites while the smaller and more highly charged Sb^{5+} occupies only the M4 site throughout.

Interplay between cation sizes and charges provides a likely explanation for the observed sequence of structures for $\text{Co}_2\text{InSbO}_6$. Efficient cation packing is favoured under the high pressure (8 GPa) at which the initial sample was synthesised. The ordered-*R*32 structure with $\text{Co}^{2+}\text{Co}^{2+}\text{In}^{3+}\text{Sb}^{5+}$ chains of cation pairs is thus stabilised as Co^{2+} (0.745 Å) is similar in size to the average (0.70 Å) of the larger In^{3+} and smaller Sb^{5+} cations. Thermal relaxation at ambient pressure leads to $(\text{Co}_{0.5}\text{In}_{0.5})^{2.5+}\text{Co}^{2+}(\text{Co}_{0.5}\text{In}_{0.5})^{2.5+}\text{Sb}^{5+}$ and then to $\text{Co}^{2+}\text{In}^{3+}\text{Co}^{2+}\text{Sb}^{5+}$ sequences of cation

pairs in the successive ordered- LiNbO_3 type products. This reduces electrostatic repulsions between cations in the dimer pairs which becomes more significant at ambient pressure where packing constraints are less important. Repulsions between cation charges q_i in the dimer pairs can be quantified in a simple nearest-neighbour approximation as $E=q_Aq_B+q_Cq_D$ for AB_2CD cation order in the chains assuming fixed cation-cation separations. The sequence of structures shown above as (1) have $E=19\rightarrow17.5\rightarrow16$ and the $\text{Co}^{2+}\text{In}^{3+}\text{Co}^{2+}\text{Sb}^{5+}$ sequence in the final ordered- LiNbO_3 product has the lowest possible electrostatic repulsion energy within the family of A_2BCO_6 structures (Figure 1), as the ordered-ilmenite alternative would have greater repulsion across $\text{Co}^{2+}\text{In}^{3+}\text{Sb}^{5+}\text{Co}^{2+}$ pairs. This lowering of cation-cation repulsion is consistent with the decrease in thermal expansion of a and V cell parameters on heating across the two transitions seen in Figure 3a.

The effects of the cation rearrangement on the magnetic properties of $\text{Co}_2\text{InSbO}_6$ were explored by comparing the original sample recovered from high pressure with ordered-*R*32 structure type (OR32 sample), with a sample subsequently heated to 1073 K having the final ordered- LiNbO_3 cation arrangement (OLN sample). Magnetic susceptibility measurements show that both samples are Curie–Weiss paramagnets at high temperatures (Figure 4), with effective

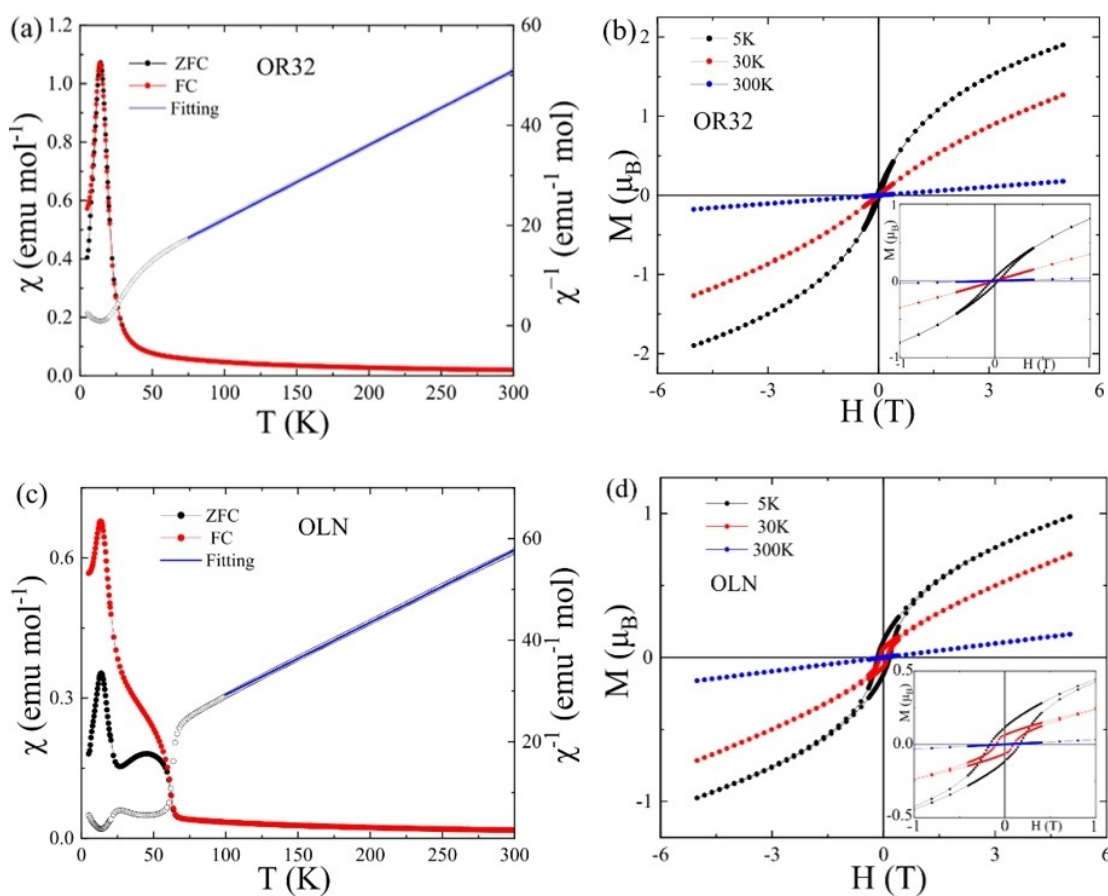


Figure 4. Magnetic measurements for a) and b) the OR32, and c) and d) the OLN, samples of $\text{Co}_2\text{InSbO}_6$. a) and c) ZFC and FC susceptibilities and inverse ZFC susceptibilities with high-temperature Curie–Weiss fits. b) and d) Magnetisation-field loops with insets showing low field regions.

paramagnetic moment $\mu_{\text{eff}} = 5.19 \mu_{\text{B}}$ per Co^{2+} and Weiss temperature $\theta = -43$ K for the OR32 sample, and $\mu_{\text{eff}} = 5.30 \mu_{\text{B}}$ and $\theta = -106$ K for OLN. The moments are in excess of spin-only values showing that strong orbital contributions are present, and similar values up to $\approx 5.20 \mu_{\text{B}}$ have been reported for other Co^{2+} oxides such as $\text{Co}_2\text{ScSbO}_6$.^[15] Negative values of θ indicate that dominant spin-spin interactions are antiferromagnetic. Both samples show deviation of the susceptibility above the Curie–Weiss variation at temperatures below ≈ 65 K suggesting antiparallel but ferrimagnetic spin alignments, given the negative values of θ . The OR32 sample shows no discontinuity or divergence of zero-field cooled and field cooled (ZFC and FC) susceptibilities, which suggest short range ferrimagnetism. However, the OLN sample has a sharp Curie transition at $T_{\text{C}} = 65$ K, similar to $T_{\text{C}} = 59$ K for isostructural ferrimagnetic $\text{Co}_2\text{ScSbO}_6$.^[15] This contrasting behaviour reflects a key difference in cobalt spin distributions in the two structures. In the *R32* structure, the Co spins are located in dimers which results in frustration between successive dimer layers, but in the ordered- LiNbO_3 structure the spins are distributed in a less frustrated, three-dimensional network. Both samples have susceptibility peaks at 15 K indicative of a possible antiferromagnetic or a spin-glass transition. The latter could result from the 20–25 % Co/In disorder between two sites observed for both samples. Neutron diffraction will be needed to confirm the spin orders or their absence. Magnetization-field loops show substantial magnetization for the two samples at low temperatures (Figure 4). The moments at 5 K and 5 T approach 2 and $1 \mu_{\text{B}}$ per $\text{Co}_2\text{InSbO}_6$ formula unit for OR32 and OLN samples, respectively. The OR32 sample exhibits a small hysteresis at 5 K (remnant magnetization $M_{\text{r}} = 0.04 \mu_{\text{B}}$ and coercive field $B_{\text{c}} = 0.03$ T) but hysteresis for the OLN sample is more substantial ($M_{\text{r}} = 0.12 \mu_{\text{B}}$ and $B_{\text{c}} = 0.16$ T), consistent with the well-defined ferrimagnetic transition for this phase.

These results demonstrate that a new A_2BCO_6 ordered-*R32* variant of the corundum structure is stabilised in $\text{Co}_2\text{InSbO}_6$ prepared at high pressure. This discovery of dimer units containing the same cations is unprecedented in ABO_3 or A_2BCO_6 corundum derivatives and likely results from similar average cation sizes in the Co_2O_9 and InSbO_9 dimer units minimising volume at pressure. High pressure may thus enable discovery of ABO_3 phases with the as-yet unreported *R32* structure. These would be of interest for the stabilization of spin-dimer materials with A/B = magnetic/non-magnetic cation combinations. $\text{Co}_2\text{InSbO}_6$ is also remarkable for showing an unprecedented sequence of two cation rearrangements on heating: first to $(\text{Co}_{0.5}\text{In}_{0.5})_2\text{CoSbO}_6$ and then to $\text{Co}_2\text{InSbO}_6$ phases with the ordered- LiNbO_3 structure. These rearrangements reduce cation-cation repulsions and favour the ordered- LiNbO_3 structure that is often observed in corundum-derived A_2BCO_6 materials. The cation distributions change magnetic properties as the final ordered- LiNbO_3 product has a sharp ferrimagnetic transition at 65 K whereas the initial ordered-*R32* phase has a broader feature indicative of short-range spin ordering. Both samples have susceptibility peaks at 15 K indicative of an antiferromagnetic or spin-glass tran-

sition. All of the $\text{Co}_2\text{InSbO}_6$ phases have acentric *R3* symmetry, and so are of interest for future exploration of acentric and multiferroic properties.

In conclusion, $\text{Co}_2\text{InSbO}_6$ and previously-reported $\text{Mn}_2\text{FeMoO}_6$ ^[18] demonstrate that high pressure may be used to recover metastable cation arrangements within the corundum family that can be thermally relaxed to new structures. This enables dependence of physical properties on the cation-ordering patterns to be explored. Further syntheses of high-pressure corundum derivatives are likely to reveal other cation-redistribution pathways, and may enable ABO_3 materials with the as-yet unreported *R32* structure type to be discovered.

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Conflict of Interest

The authors declare no conflict of interest.

Data Availability Statement

The data that support the findings of this study are openly available in Edinburgh DataShare at <https://datashare.ed.ac.uk/handle/10283/838>.

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- [1] J. M. Rondinelli, C. J. Fennie, *Adv. Mater.* **2012**, *24*, 1961–1968.
- [2] R. Shankar P N, F. Orlandi, P. Manuel, W. Zhang, P. S. Halasyamani, A. Sundaresan, *Chem. Mater.* **2020**, *32*, 5641–5649.
- [3] J. Choisnet, A. Rulmont, P. Tarte, *J. Solid State Chem.* **1988**, *75*, 124–135.
- [4] P. M. Woodward, A. W. Sleight, L. S. Du, C. P. Grey, *J. Solid State Chem.* **1999**, *147*, 99–116.
- [5] R. Shankar P N, S. Mishra, S. Athinarayanan, *APL Mater.* **2020**, *8*, 040906.
- [6] M. Ye, D. Vanderbilt, *Phys. Rev. B* **2016**, *93*, 134303.
- [7] H. Niu, M. J. Pitcher, A. J. Corkett, S. Ling, P. Mandal, M. Zanella, K. Dawson, P. Stamenov, D. Batuk, A. M. Abakumov, C. L. Bull, R. I. Smith, C. A. Murray, S. J. Day, B. Slater, F. Cora, J. B. Claridge, M. J. Rosseinsky, *J. Am. Chem. Soc.* **2017**, *139*, 1520–1531.
- [8] P. S. Halasyamani, K. R. Poeppelmeier, *Chem. Mater.* **1998**, *10*, 2753–2769.

- [9] K. K. Wong, *Properties of Lithium Niobate*, INSPEC, UK, 2002.
- [10] G. H. Cai, M. Greenblatt, M. R. Li, *Chem. Mater.* **2017**, 29, 5447–5457.
- [11] A. M. Arévalo-López, J. P. Attfield, *Phys. Rev. B* **2013**, 88, 104416.
- [12] M.-R. Li, D. Walker, M. Retuerto, T. Sarkar, J. Hadermann, P. W. Stephens, M. Croft, A. Ignatov, C. P. Grams, J. Hemberger, I. Nowik, P. S. Halasyamani, T. T. Tran, S. Mukherjee, T. S. Dasgupta, M. Greenblatt, *Angew. Chem. Int. Ed.* **2013**, 52, 8406–8410; *Angew. Chem.* **2013**, 125, 8564–8568.
- [13] A. M. Arévalo-López, E. Solana-Madruga, E. P. Arévalo-López, D. Khalyavin, M. Kepa, A. J. Dos Santos-García, R. Sáez-Puche, J. P. Attfield, *Phys. Rev. B Phys. Rev. B* **2018**, 98, 214403.
- [14] E. Solana-Madruga, A. Dos Santos-García, A. Arévalo-López, D. Ávila-Brandé, C. Ritter, J. P. Attfield, R. Sáez-Puche, *Dalton Trans.* **2015**, 44, 20441–20448.
- [15] K. Ji, E. Solana-Madruga, A. M. Arévalo-López, P. Manuel, C. Ritter, A. Senyshyn, J. P. Attfield, *Chem. Commun.* **2018**, 54, 12523–12526.
- [16] S. A. Ivanov, R. Mathieu, P. Nordblad, R. Tellgren, C. Ritter, E. Politova, G. Kaleva, A. Mosunov, S. Stefanovich, M. Weil, *Chem. Mater.* **2013**, 25, 935–945.
- [17] M.-R. Li, M. Croft, P. W. Stephens, M. Ye, D. Vanderbilt, M. Retuerto, Z. Deng, C. P. Grams, J. Hemberger, J. Hadermann, W.-M. Li, C.-Q. Jin, F. O. Saouma, J. I. Jang, H. Akamatsu, V. Gopalan, D. Walker, M. Greenblatt, *Adv. Mater.* **2015**, 27, 2177–2181.
- [18] M.-R. Li, M. Retuerto, P. W. Stephens, M. Croft, D. Sheptyakov, V. Pomjakushin, Z. Deng, H. Akamatsu, V. Gopalan, J. Sánchez-Benítez, F. O. Saouma, J. I. Jang, D. Walker, M. Greenblatt, *Angew. Chem. Int. Ed.* **2016**, 55, 9862–9867; *Angew. Chem.* **2016**, 128, 10016–10021.
- [19] Y. S. Oh, S. Artyukhin, J. J. Yang, V. Zapf, J. W. Kim, D. Vanderbilt, S. W. Cheong, *Nat. Commun.* **2014**, 5, 3201.
- [20] J. W. Kim, S. Artyukhin, E. D. Mun, M. Jaime, N. Harrison, A. Hansen, J. J. Yang, Y. S. Oh, D. Vanderbilt, V. S. Zapf, S.-W. Cheong, *Phys. Rev. Lett.* **2015**, 115, 137201.
- [21] M. O. Yokosuk, A. al-Wahish, S. Artyukhin, K. R. O'Neal, D. Mazumdar, P. Chen, J. J. Yang, Y. S. Oh, S. A. McGill, K. Haule, S.-W. Cheong, D. Vanderbilt, J. L. Musfeldt, *Phys. Rev. Lett.* **2016**, 117, 147402.
- [22] E. Solana-Madruga, C. Aguilar-Maldonado, C. Ritter, M. Huvé, O. Mentré, J. P. Attfield, A. M. Arévalo-López, *Chem. Commun.* **2021**, 57, 2521–2514.
- [23] M. R. Li, E. E. McCabe, P. W. Stephens, M. Croft, L. Collins, S. V. Kalinin, Z. Deng, M. Retuerto, A. Sen Gupta, H. Padmanabhan, V. Gopalan, C. P. Grams, J. Hemberger, F. Orlandi, P. Manuel, W.-M. Li, C.-Q. Jin, D. Walker, M. Greenblatt, *Nat. Commun.* **2017**, 8, 2037.
- [24] R. D. Shannon, *Acta Crystallogr. Sect. A* **1976**, 32, 751–767.

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