

Directional linearly polarized terahertz emission from argon clusters irradiated by noncollinear double-pulse beams

Cite as: Appl. Phys. Lett. **111**, 241107 (2017); <https://doi.org/10.1063/1.4991736>

Submitted: 22 June 2017 • Accepted: 06 November 2017 • Published Online: 15 December 2017

 Kazuaki Mori, Masaki Hashida, Takeshi Nagashima, et al.



View Online



Export Citation



CrossMark

ARTICLES YOU MAY BE INTERESTED IN

Increased energy of THz waves from a cluster plasma by optimizing laser pulse duration
AIP Advances **9**, 015134 (2019); <https://doi.org/10.1063/1.5075712>

Intense terahertz emission from atomic cluster plasma produced by intense femtosecond laser pulses
Applied Physics Letters **99**, 261503 (2011); <https://doi.org/10.1063/1.3672814>

Observation of broadband terahertz wave generation from liquid water
Applied Physics Letters **111**, 071103 (2017); <https://doi.org/10.1063/1.4990824>

 QBLOX



1 qubit

Shorten Setup Time

Auto-Calibration

More Qubits

Fully-integrated

Quantum Control Stacks

Ultrastable DC to 18.5 GHz

Synchronized <<1 ns

Ultralow noise



100s qubits

visit our website >

Directional linearly polarized terahertz emission from argon clusters irradiated by noncollinear double-pulse beams

Kazuaki Mori,^{1,2} Masaki Hashida,^{1,2} Takeshi Nagashima,³ Dazhi Li,⁴ Kensuke Teramoto,^{1,2} Yoshihide Nakamiya,¹ Shunsuke Inoue,^{1,2} and Shuji Sakabe^{1,2}

¹Advanced Research Center for Beam Science, Institute for Chemical Research, Kyoto University, Gokasho, Uji, Kyoto 611-0011, Japan

²Department of Physics, Graduate School of Science, Kyoto University, Kitashirakawa, Sakyo, Kyoto 606-8502, Japan

³Faculty of Science and Engineering, Setsunan University, 17-8 Ikeda-Nakamachi, Neyagawa, Osaka 572-8508, Japan

⁴Institute for Laser Technology, 1-8-4 Utsubo-honmachi, Nishi-ku, Osaka 550-0004, Japan

(Received 22 June 2017; accepted 6 November 2017; published online 15 December 2017)

It has been demonstrated that the interaction between argon clusters and intense femtosecond double laser pulses with appropriate intervals in time and space provides important properties for terahertz electromagnetic wave generation, namely, high forward directivity, power enhancement, and linear polarization with a variable direction. Irradiating argon clusters with double pulses (1 and 3 mJ, 40 fs, 810 nm) in 133-ps and 40- μ m intervals results in terahertz wave emission in the forward direction that is 10 times greater than that for a single pulse. The polarization direction of terahertz electromagnetic waves can be varied by changing the relative focal positions of the first and second pulses. *Published by AIP Publishing.* <https://doi.org/10.1063/1.4991736>

In the past few decades, the scientific and engineering applications of terahertz (THz) waves have advanced remarkably with the rapid progress of laser technology. Intense THz waves have attracted attention for various applications.^{1–5} Currently, 0.9-mJ THz wave pulses are generated by optical rectification in an organic nonlinear crystal.⁶ However, it is difficult to increase the THz wave energy with this method because the damage threshold of the nonlinear crystal limits the incident laser energy. Using a plasma as the THz-wave-source element would reduce this limitation considerably and is a promising method for generating intense THz waves. Laser plasmas produced from a solid target generate 0.7-mJ THz waves by sheath proton acceleration.⁷ However, this method is not suitable for repeated generation of THz waves because of debris from the solid target and the difficulty in replenishing the target continuously. A gas target can overcome these problems but provides THz waves with much lower energies than those from a solid target.^{8,9} To have the benefits of both solid and gas targets, clusters consisting of atoms have been proposed as a target,¹⁰ and the generation of THz waves from cluster plasmas has been studied experimentally and theoretically.^{11–14} Jahangiri *et al.* reported that an atomic-cluster plasma generates 600 times more THz waves than a gas plasma does,^{11,12} thereby showing the potential of an atomic-cluster plasma as a THz wave source.

However, in conventional methods, THz waves from cluster plasmas produced by irradiation with a single-pulse laser beam generally have radial polarization and conical angular distribution. This is inconvenient for applications because such characteristics make handling the waves difficult and less efficient. Therefore, it is important to develop intense THz wave sources with good characteristics, in particular directional linearly polarized emission. For a gas target, the use of multiple incident laser pulses has been proposed to improve the radiation properties. Generating THz

waves via two-color laser focusing in a gas has overcome the problematic radiation properties of THz waves from a gas plasma.^{15–17} In this scheme, linearly polarized THz waves are generated in the forward direction. This requires precise adjustment of the temporal deviation between the two pulses to match their phases.¹⁸ An additional challenge is the damage threshold of the second-harmonic-generation crystal, limiting the THz output. As a method using two single-color pulses, a bi-filamentation scheme has been proposed.¹⁹ Although THz wave generation in the forward direction has been measured, the power is much lower than that for solid or cluster targets.

In this study, we use a cluster target to generate the plasma and the THz waves separately with double-pulse beams to search for appropriate intervals in both time and space for the pulses to be focused into the plasma. THz waves are emitted with linear polarization, and their power is greatly enhanced in the direction of laser propagation.

A schematic diagram of the experimental setup is shown in Fig. 1. We use a Ti:sapphire chirped pulse amplification laser system operating at a central wavelength of 810 nm. The laser pulse has a duration of 40 fs (FWHM). The laser beam is separated into two beams: the first pulse with an energy of 1 mJ and the second pulse with an energy of 3 mJ. The pulse energies of the two pulses are measured separately in front of a glass chamber with a pyroelectric energy meter. The two pulses are linearly polarized with a polarization parallel to the optical table. The time delay between the two pulses is varied from 0 to 500 ps with 0.03-fs resolution by a direct-drive translation stage (ANT95-100-L, Aerotech Inc.). The two beams are recombined and focused onto the argon clusters by a plano-convex lens with a focal length of 200 mm ($F = 10$). The first and second laser beams can be focused in Gaussian spots with diameters of 12 and 15 μ m ($F\text{We}^{-1}\text{M}$), respectively. The focal position of the second

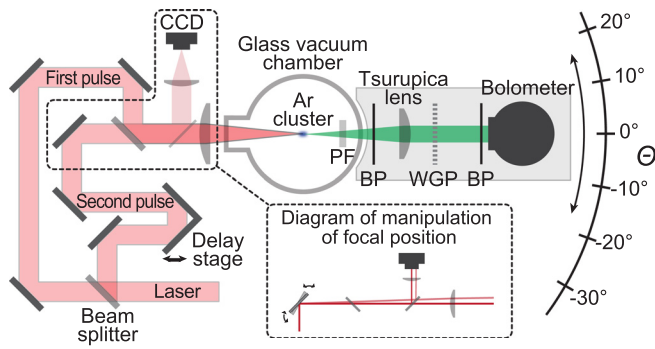


FIG. 1. Schematic diagram of the experimental setup. PF: polystyrene foam; BP: black polypropylene filter; WGP: wire-grid polarizer; and Θ : detection angle of THz waves.

pulse can be varied from 0 to $100\ \mu\text{m}$ (corresponding to 0 – $500\ \mu\text{rad}$) by a mirror just before the beam combiner. The focal position (pointing) of the two beams is monitored using a CCD camera during the experiment.

Argon clusters are produced in the center of a glass vacuum chamber by injecting argon gas with a backing pressure of 7 MPa through a conical nozzle whose length is 20 mm and whose input and output diameters are 0.5 and 5 mm, respectively. The distance of the nozzle from the focal point is 10 mm. The typical radius of clusters is measured as $R_0 \approx 5\ \text{nm}$ by the Rayleigh scattering method,²⁰ and the atomic density in the interaction region is estimated to be approximately $6 \times 10^{17}\ \text{cm}^{-3}$. The THz waves generated by irradiating the argon clusters with intense laser pulses are collimated by a Tsurupica lens (Pax Co.) and detected using a helium-cooled InSb bolometer (QFI-2BI, QMC Instruments Ltd.) whose detectable frequency range is up to approximately 2 THz. Five-millimeter-thick polystyrene foam and 60- μm -thick black polypropylene filters are placed in front of the bolometer to block residual laser pulses and undesirable emissions from the plasma. The detection system consists of the bolometer, a wire grid, the Tsurupica lens, and the black polypropylene filters on an optical plate (Fig. 1). To measure the angular distribution of the THz wave emission, the detection system is moved in an arc from $\Theta = 80^\circ$ to -60° with a resolution of 4° , and a chamber made of fused silica glass with 90% transparency at 0.5 THz is used.²¹ The polarization properties and spectrum of the THz waves are measured with wire-grid polarizers (extinction ratio approximately 10^{-5} at 0.5 THz) and a Martin-Puplett interferometer, respectively. Each data point is obtained by averaging 16 pulses.

First, the THz power dependence on the focal position of the second pulse relative to that of the first pulse is measured. Figure 2(a) shows the definition of the relative focal position of the double-pulse beams irradiating the argon clusters. The focal position of the second pulse relative to that of the first pulse is shown in Cartesian coordinates (x, y) and cylindrical coordinates (r, θ) . Figure 2(b) shows the THz power dependence on the focal position of the second pulse relative to that of the first pulse at a delay time of 133 ps. Although almost no THz waves are generated when the first and second pulses are irradiated collinearly [i.e., $(x, y) = (0, 0)$], THz wave generation is enhanced when the two pulses are irradiated noncollinearly [i.e., $(x, y) \neq (0, 0)$] and is

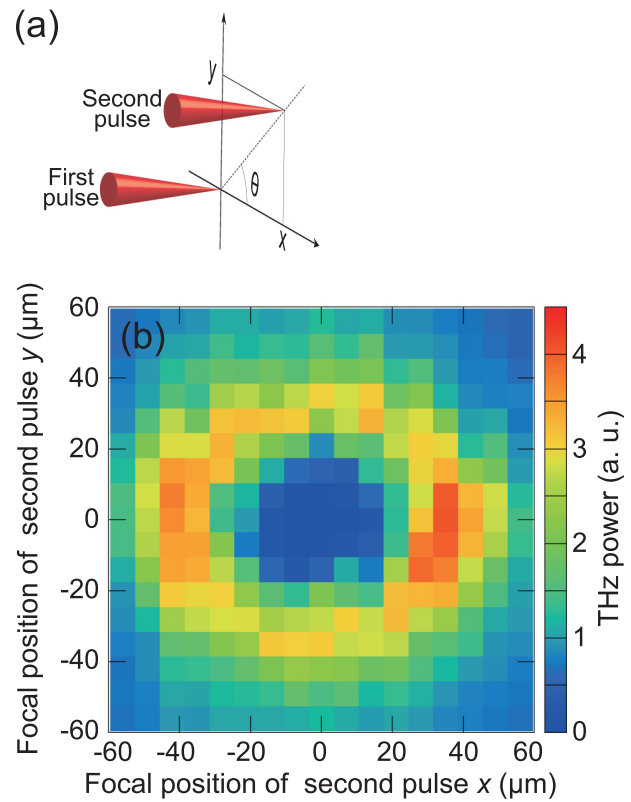


FIG. 2. (a) Schematic of x , y , and θ defined as the focal position of the second pulse relative to the focal position of the first pulse. (b) Dependence of THz power on the focal position (x, y) of the second pulse. The delay time between the two pulses is 133 ps. THz power is detected in the laser-propagation direction ($\Theta = 0^\circ$ in Fig. 1).

maximized by shifting the second pulse by approximately $40\ \mu\text{m}$ regardless of θ . In this experiment, THz waves are detected in the direction of $\Theta = 0^\circ$. The THz waves are detected at a solid angle of $7.9 \times 10^{-3}\ \text{sr}$, which is much larger than the maximum shift of the second pulse ($500\ \mu\text{rad}$). Therefore, the shift of the second pulse is small enough not to influence THz detection. Under conditions in which the delay time is 133 ps and the focal position of the second pulse is $(x, y) = (0, 40\ \mu\text{m})$, the THz-power stability (standard deviation/average value) is 11%.

Figure 3(a) shows the dependence of THz power on the delay time of the second pulse. The focal position of the second pulse is shifted to $(x, 0)$ ($\theta = 0$). As the delay time is increased from 1 to 500 ps, position x of the second pulse that gives the maximum THz wave power increases, which is shown in Fig. 3(b) more clearly. Figure 3(b) shows the second-pulse positions that give each maximum THz power in Fig. 3(a) as a function of their delay time. The second-pulse position for maximum THz wave emission moves at a speed of $8 \times 10^4\ \text{m/s}$.

Figure 4 shows the angular distribution of THz waves from a cluster plasma produced by double-pulse laser beams compared with that produced by a single-pulse beam. The angular distribution of the THz emission is then measured under conditions of a delay time of 133 ps and $(x, y) = (0, 40\ \mu\text{m})$. The propagation direction of the first laser pulse is denoted by $\Theta = 0^\circ$. For a single-pulse beam, THz waves are emitted at $\pm 30^\circ$, which is consistent with previous results.¹¹ For double-pulse laser beams, the THz power is increased by

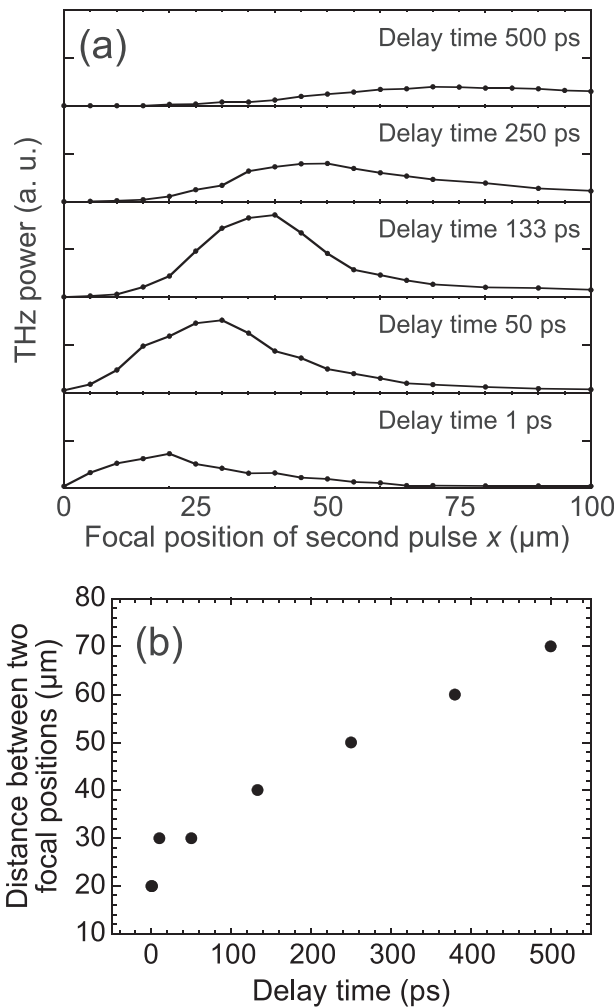


FIG. 3. (a) Temporal development of the THz power dependence with the distance between focal positions of the first and second pulses in the horizontal plane. (b) Distance between the two focal positions for various delay times at which the THz power reaches its peak value.

more than 10 times that for the single-pulse beam in the laser-propagation direction. The THz emission is directed much more in the laser-propagation direction when the first pulse forms the plasma and the second pulse generates the THz waves, whereas THz waves are emitted conically in the single-pulse case. The frequency of the THz waves is measured separately by the Martin–Puplett interferometer in the range of 0.03–2 THz. The observed spectrum of THz waves covers 0.1–1 THz and peaks around 0.2–0.5 THz with a frequency width (FWHM) of approximately 0.2–0.3 THz.

The polarization of the THz waves is measured using a wire grid. The experiment is performed at a laser-propagation direction of $\Theta = 0^\circ$. The inset in Fig. 5(b) shows the transmitted THz wave power as a function of the wire-grid angle Φ for $\theta = -1^\circ$ and 57° . The THz waves are linearly polarized at every value of θ . The polarization angle φ of the THz waves is obtained by analyzing the wire-grid angle Φ for maximum transmission; the transmitted THz power is well fitted by the function $I = \cos^2(\Phi - \varphi)$. The polarization angle of the THz waves is 4° for $\theta = -1^\circ$ and 59° for $\theta = 57^\circ$ and increases with the second-pulse angle θ . Figure 5(b) summarizes the polarization angle φ of the THz waves for various angles θ of the second pulse, which shows

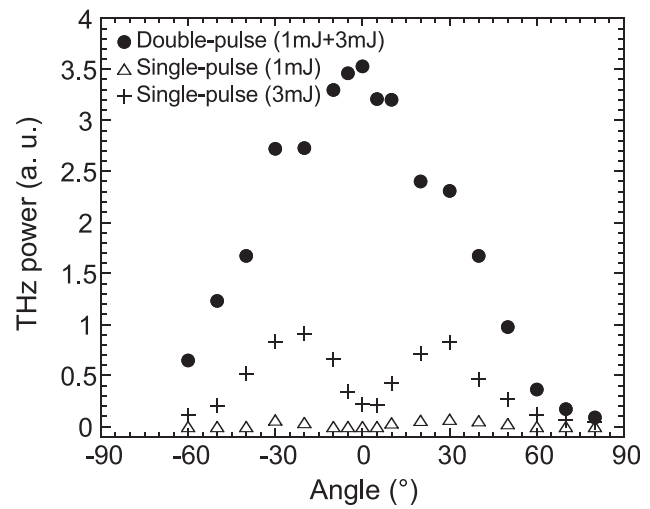


FIG. 4. Angular distributions of THz waves for a double-pulse beam (pulse energies of 1 and 3 mJ) and a single-pulse beam (pulse energies of 1 or 3 mJ); 0° is the laser-propagation direction. THz emission is measured under conditions of a delay time of 133 ps and $(x, y) = (0, 40 \mu\text{m})$.

that $\varphi = \theta$. Thus, the polarization angle of linearly polarized THz waves can be rotated by the relative position of the first and second pulses without changing the incident laser polarization.

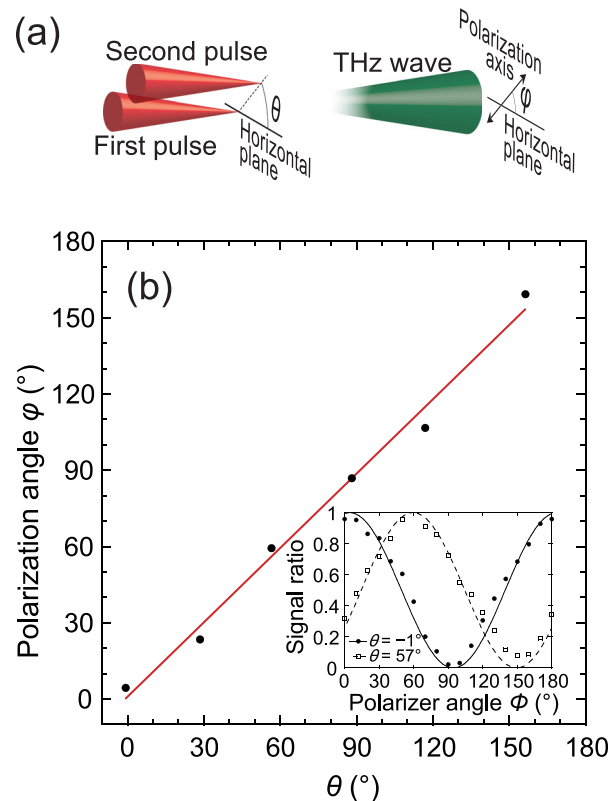


FIG. 5. (a) Schematic diagram of the definition of θ and φ . θ is defined as the angle between the line connecting the two focal positions and the horizontal plane. φ is the angle between the direction of the vibration of the THz electric field and horizontal plane. (b) Polarization angle φ of the THz wave for various θ . The red line is the least-squares regression line $\varphi = 0.97\theta + 1.0$, $R = 0.99$. The delay time between the two pulses is 133 ps. THz power is detected in the laser-propagation direction ($\Theta = 0^\circ$ in Fig. 1). Inset: THz transmission of the wire-grid polarizer for various wire-grid angles. Horizontal linearly polarized THz waves are transmitted by the wire-grid polarizer when the wire-grid angle is 0° .

The characteristics of THz waves emitted from clusters irradiated by intense femtosecond double laser pulses can be summarized as follows: (1) higher directivity than for a single pulse, (2) high intensity enhancement with appropriate intervals between the double pulses in space and time, and (3) linear polarization in the direction of the line connecting the focal points of the double pulses. These characteristics are useful for applications because they result in (1) highly efficient use, (2) higher power, and (3) variable polarization, respectively.

Linearly polarized directional THz emission from plasmas cannot be explained by conventional mechanisms such as transition radiation²² or transition-Cherenkov radiation²³ which have been studied as THz generation mechanisms for laser plasmas produced from gas targets. In the present work, because the laser intensity is in the range of 10^{16} – 10^{17} W/cm², it is expected that THz waves should be emitted in the direction $\Theta = 80^\circ$ – 90° .¹² Although directional THz emission via four-wave optical rectification or second-order optical rectification inside an asymmetric filament^{24,25} and parametric decay of laser light²⁶ have been reported in other works on THz generation using a plasma filament, the polarization of the THz waves was elliptical, not linear.

The generation of directional linearly polarized THz waves from clusters under double-pulse irradiation can be understood as the generation of THz waves from a plasma channel near Coulomb-exploded ions, which work as a bias. First, argon clusters are Coulomb-exploded by the first pulse in the focal space, and ions are emitted from the focal space to form expanding (positively charged) ion clouds. Then, the second pulse is irradiated from near the explosion front and provides electrons along the ion-cloud front as the second pulse travels. A transverse current²⁷ or dipole moment²⁸ would be excited in the plasma channel produced by the second pulse, resulting in emission of electromagnetic waves in the laser propagation direction just like THz wave generation in a DC field.^{27,29} The electromagnetic wave is emitted most efficiently at the ion-cloud front, which moves as fast as the ions that are emitted from the Coulomb-exploded clusters. Because the laser intensity of the first pulse is in the range of 10^{16} – 10^{17} W/cm², the cluster survives as a skinned cluster. Based on the scaling law of the ionized radius with laser intensity, the remaining core part of the cluster is estimated to have a radius of 2.5 nm. The maximum and probable energy of argon ions can be estimated as $E_{\max} = 4\pi Z^2 e^2 n_i (R_0^3 - R^3) / (3R_0)$, where Z is the ion charge state, e is the electron charge, n_i is the ion density, and R_0 and R are the radii of an initial argon cluster and remaining core argon cluster, respectively. For $R_0 = 5$ nm, $R = 2.5$ nm, and $Z = 7$, we have $E_{\max} = 1.6 \times 10^2$ keV. However, both the laser intensity and cluster radius are distributed spatially, which contributes to the observed distributions in the actual experiments. We have reported previously that the most probable energy of argon ions is distributed in the energy range of 1–14 keV at a laser intensity of 10^{16} W/cm².³⁰ These energies correspond to speeds of 7×10^4 to 2.6×10^5 m/s, which agree well with the temporal development speed of the optimal distance of the focal positions in Fig. 3(b). In this mechanism, the THz wave is polarized in the direction along a line connecting the focal points of the double pulses and is consistent with the

experimental results. The characteristics of the measured THz wave show that the combination of a cluster target and double-pulse irradiation is effective for generating THz waves. Future work will involve systematic studies of the mechanism of THz wave generation from cluster targets irradiated with double-pulse beams.

In summary, distinguishing properties have been observed in the emission of THz waves from argon clusters irradiated by intense femtosecond double laser pulses with appropriate intervals in time and space. These properties are high directivity, power enhancement, and linear polarization, which can be varied by changing the pointing of the second pulse relative to that of the first pulse. All these properties are useful for applications.

This work was financially supported by the Collaborative Research Program of the Institute for Chemical Research, Kyoto University (Grant Nos. 2016-4, 2017-2), and a Grant-in-Aid for Scientific Research (C) (JP16K06745) from the Ministry of Education, Culture, Sports, Science and Technology (MEXT), Japan.

¹X. C. Zhang, A. Shkurinov, and Y. Zhang, *Nat. Photonics* **11**, 16 (2017).

²C. Kealhofer, W. Schneider, D. Ehberger, A. Ryabov, F. Krausz, and P. Baum, *Science* **352**, 429 (2016).

³W. R. Huang, E. A. Nanni, K. Ravi, K. H. Hong, A. Fallahi, L. J. Wong, P. D. Keathley, L. E. Zapata, and F. X. Kärtner, *Sci. Rep.* **5**, 14899 (2015).

⁴K. Kovács, E. Balogh, J. Hebling, V. Tosa, and K. Varjú, *Phys. Rev. Lett.* **108**, 193903 (2012).

⁵M. Shalaby, C. Vicario, and C. P. Hauri, *Appl. Phys. Lett.* **108**, 182903 (2016).

⁶C. Vicario, A. V. Ovchinnikov, S. I. Ashtikov, M. B. Agranat, V. E. Fortov, and C. P. Hauri, *Opt. Lett.* **39**, 6632 (2014).

⁷A. Gopal, P. Singh, S. Herzer, A. Reinhard, A. Schmidt, U. Dillner, T. May, H.-G. Meyer, W. Ziegler, and G. G. Paulus, *Opt. Lett.* **38**, 4705 (2013).

⁸T. Löffler, F. Jacob, and H. G. Roskos, *Appl. Phys. Lett.* **77**, 453 (2000).

⁹C. D'Amico, A. Houard, M. Franco, B. Prade, A. Mysyrowicz, A. Couairon, and V. T. Tikhonchuk, *Phys. Rev. Lett.* **98**, 235002 (2007).

¹⁰T. Nagashima, H. Hirayama, K. Shibuya, M. Hangyo, M. Hashida, S. Tokita, and S. Sakabe, *Opt. Express* **17**, 8907 (2009).

¹¹F. Jahangiri, M. Hashida, T. Nagashima, S. Tokita, M. Hangyo, and S. Sakabe, *Appl. Phys. Lett.* **99**, 261503 (2011).

¹²F. Jahangiri, M. Hashida, S. Tokita, T. Nagashima, M. Hangyo, and S. Sakabe, *Appl. Phys. Lett.* **102**, 191106 (2013).

¹³A. A. Frolov, *Plasma Phys. Rep.* **42**, 637 (2016).

¹⁴A. V. Balakin, M. S. Dzhidzhoev, V. M. Gordienko, M. N. Esaulkov, I. A. Zhvaniya, K. A. Ivanov, I. A. Kotelnikov, N. A. Kuzechkin, I. A. Ozheredov, V. Y. Panchenko, A. B. Savel'ev, M. B. Smirnov, P. M. Solyankin, and A. P. Shkurinov, *IEEE Trans. Terahertz Sci. Technol.* **7**, 70 (2017).

¹⁵D. J. Cook and R. M. Hochstrasser, *Opt. Lett.* **25**, 1210 (2000).

¹⁶J. Das and M. Yamaguchi, *J. Opt. Soc. Am. B* **30**, 1595 (2013).

¹⁷R. G. Ungureanu, O. V. Grigore, M. P. Dinca, G. V. Cojocaru, D. Ursescu, and T. Dascalu, *Laser Phys. Lett.* **12**, 045301 (2015).

¹⁸X. Xie, J. Dai, and X.-C. Zhang, *Phys. Rev. Lett.* **96**, 075005 (2006).

¹⁹Y. Liu, A. Houard, B. Prade, S. Akturk, A. Mysyrowicz, and V. T. Tikhonchuk, *Phys. Rev. Lett.* **99**, 135002 (2007).

²⁰M. Hirokane, S. Shimizu, M. Hashida, S. Okada, S. Okihara, F. Sato, T. Iida, and S. Sakabe, *Phys. Rev. A* **69**, 063201 (2004).

²¹M. Naftaly and R. E. Miles, *J. Non-Cryst. Solids* **351**, 3341 (2005).

²²U. Happek, A. J. Sievers, and E. B. Blum, *Phys. Rev. Lett.* **67**, 2962 (1991).

²³C. D'Amico, A. Houard, S. Akturk, Y. Liu, J. L. Bloas, M. Franco, B. Prade, A. Couairon, V. T. Tikhonchuk, and A. Mysyrowicz, *New J. Phys.* **10**, 013015 (2008).

²⁴Y. Chen, C. Marceau, W. Liu, Z.-D. Sun, Y. Zhang, F. Théberge, M. Châteauneuf, J. Dubois, and S. L. Chin, *Appl. Phys. Lett.* **93**, 231116 (2008).

- ²⁵Y. Zhang, Y. Chen, C. Marceau, W. Liu, Z.-D. Sun, S. Xu, F. Théberge, M. Châteauneuf, J. Dubois, and S. L. Chin, [Opt. Express](#) **16**, 15483 (2008).
- ²⁶F. Jahangiri, M. Hashida, S. Tokita, T. Nagashima, K. Otani, M. Hangyo, and S. Sakabe, [Appl. Phys. Express](#) **5**, 026201 (2012).
- ²⁷A. Houard, Y. Liu, B. Prade, V. T. Tikhonchuk, and A. Mysyrowicz, [Phys. Rev. Lett.](#) **100**, 255006 (2008).
- ²⁸N. A. Panov, O. G. Kosareva, V. A. Andreeva, A. B. Savel'ev, D. S. Uryupina, R. V. Volkov, V. A. Makarov, and A. P. Shkurinov, [JETP Lett.](#) **93**, 638 (2011).
- ²⁹T. Löffler, M. Kress, M. Thomson, and H. G. Roskos, [Acta Phys. Pol. A](#) **107**, 99 (2005).
- ³⁰S. Sakabe, K. Shirai, M. Hashida, S. Shimizu, and S. Masuno, [Phys. Rev. A](#) **74**, 043205 (2006).